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NATURAL, FORCED AND MIXED CONVECTION IN FIBROUS INSULATION

Aydin Konuk

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ABSTRACT

A numerical solution of flow and temperature distribution in fibrous insulation has been obtained. Rectangular and cylindrical geometries have been modeled. Boundary conditions included permeable hot wall and convective heat transfer at the walls. Good agreement has been obtained with published experimental and numerical values on rectangular cavities. The computed velocity and temperature profiles gave a better understanding of flow and heat transfer phenomena in fibrous insulation. Local cold wall and average Nusselt numbers presented provide useful information in the design of the fibrous insulation for concrete reactor vessel and primary coolant piping of the gas cooled nuclear power plants. Average Nusselt number has been correlated with Rayleigh number when only natural convection is present, and with Rayleigh and Reynolds numbers when a combination of natural and forced convection is present.

1 - INTRODUCTION

Extensive theoretical and experimental investigations have been conducted on fibrous type thermal insulation of nuclear reactor during the last decade^(5,3,4,11). These studies were concerned with rectangular geometries of interest in the nuclear reactors. Investigation of the fibrous insulation in cylindrical geometries however does not appear in the literature, except one experimental work⁽²⁾.

The cylindrical geometry, among others, applies to double-pipe systems, which are widely used to carry hot fluids under high pressure⁽⁶⁾. The inner pipe, subject to high temperature, does not carry the pressure load since it is either perforated or made of non-sealed pieces. The outer wall works at low temperatures because of the temperature drop in the insulation layer and carries the pressure load. Thus, the insulation layer is subjected to large temperature differences across its thickness and to pressure gradients along its hot face. As a result, natural convection and forced convection permeation flows can develop within the insulating layer. These flows can degrade its performance and create hot local spots. This report describes the development and application of a computer program to predict the steady-state performance of fibrous insulation enclosed between two concentric cylinders (horizontal or vertical) or in a rectangular cavity.

Experiments^(5,3) show that the flow in fibrous insulation obeys Darcy's Law. For such materials, classified as porous media, the momentum equation is simplified since the pressure gradient is a linear function of the velocity.

It is also assumed that solid and the gas have the same local temperatures. In the present development, except those specified above (and common to all previous work), no other assumption is made. The fluid properties are allowed to vary locally with temperature, using empirical equations to predict them. In particular the local derivatives of the density on the right hand side of the momentum equation are evaluated directly, rather than being converted into temperature gradients by taking the compressibility factor Z of the gas as $Z = 1$. Thus, no inaccuracies are introduced when the gas departs largely from ideal gas behaviour.

Various boundary conditions are included in the model. Except for the horizontal cylinder, all three geometries considered include, on both the inner and the outer walls, convective boundary conditions or specified wall temperatures, and external flow over a perforated inner surface with its own permeability. The last boundary condition for a horizontal cylinder requires a three dimensional solution, therefore it has been excluded. The derivation of the equations and boundary conditions is described in the second section of the report.

A numerical method has been developed for the solution of the two non-linear partial differential equations (momentum and energy equations) with the proper boundary conditions resulting from the present formulation. Except at high Rayleigh or Reynolds numbers, no convergence problems have been encountered. The numerical procedure is the subject of the third section.

The computed results are in the fourth section. Velocity and temperature fields as well as local and average Nusselt numbers are presented.

Conclusions and remarks concerning future analytical and experimental work are in the final section.

Appendix contains a listing of the computer program with input information.

2 – FORMULATION OF EQUATIONS AND BOUNDARY CONDITIONS

2.1 – Horizontal Cylinder

The equations and boundary conditions describing flow and heat transfer in the case of a horizontal cylinder are formulated considering the half-circular cross section illustrated in Figure 1.

Considerations of symmetry allow the modeling of one half of the cross section instead of the full cross section. This has the advantage of decreasing the computer execution time and memory requirements. Equations of conservation of mass, momentum and energy are written for the two dimensional problem, the independent variables being the angular direction θ and the radial direction r . The axial direction is not considered assuming that flow and temperature fields do not change in that direction. This is true there is no permeation flow through the inside wall. A three dimensional modeling of the horizontal cylinder has not been possible within the framework of the present numerical method of solution due to the large increase in the computer execution time and memory requirements.

Conservation of Mass

The equation of conservation of mass is:

$$\frac{\partial}{\partial r} (r \rho v_r) + \frac{\partial}{\partial \theta} (\rho v_\theta) = 0 \quad (1)$$

Conservation of Momentum

For porous media, equation of conservation of momentum reduces to Darcy's law. In two dimensional cylindrical coordinates with r and θ , Darcy's law is written as

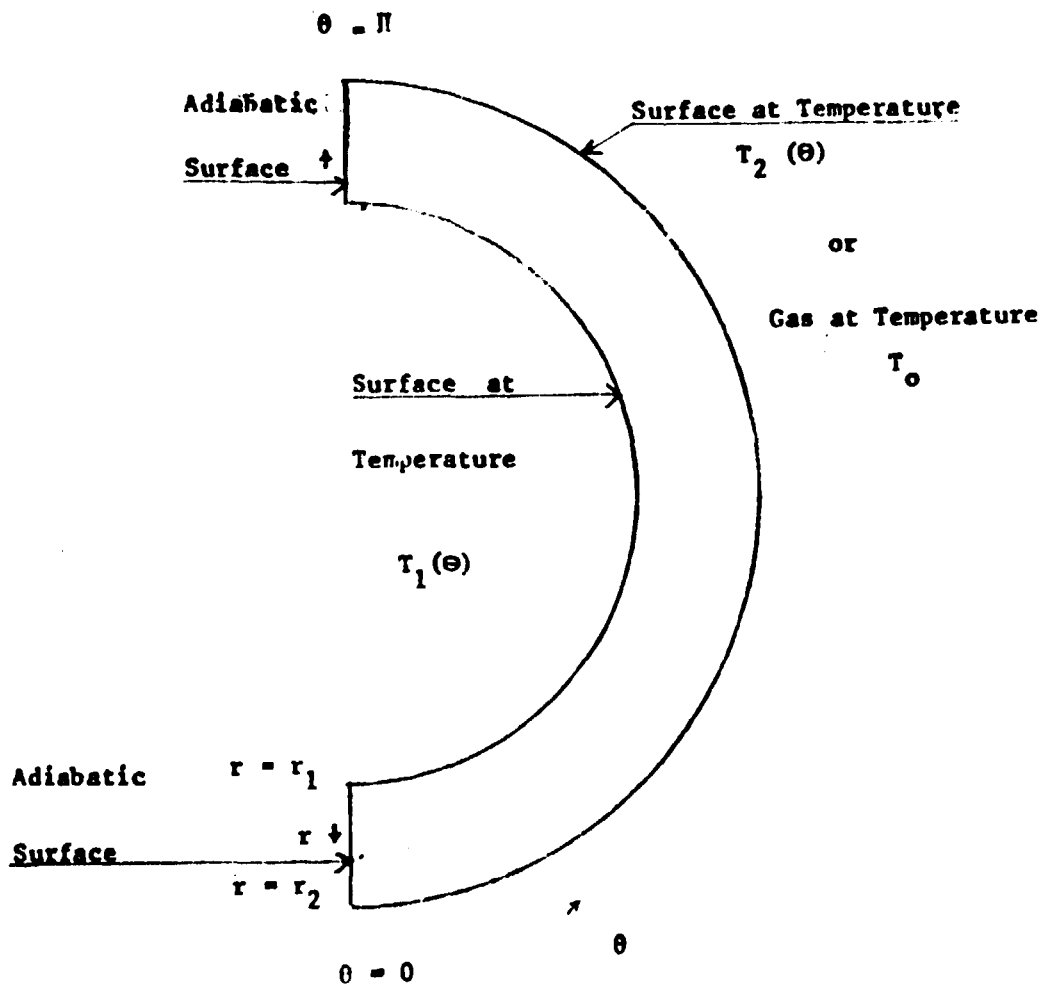


Figure 1 - Horizontal Cylinder Geometry for which the boundary conditions are described

$$v_r = \frac{K_r}{\mu} \left(\rho g_r - \frac{\partial \rho}{\partial r} \right) \quad (2)$$

for the r - momentum equation and as

$$v_\theta = \frac{K_\theta}{\mu} \left(\rho g_\theta - \frac{\partial \rho}{r \partial \theta} \right) \quad (3)$$

for the θ - momentum equation. The velocity v in these equations is the superficial velocity (volume flow rate of flow through a unit cross section area of the fluid plus solid) averaged over a small region of space-small with respect to the macroscopic dimensions of the flow system but large with respect to the pore size.

Equations (2) and (3) are combined to eliminate the pressure, which gives

$$\frac{1}{K_r} \frac{\partial}{\partial \theta} (\mu v_r) - \frac{1}{K_\theta} \frac{\partial}{\partial r} (\mu v_\theta) = \frac{\partial}{\partial \theta} (\rho g_r) - g_\theta \frac{\partial}{\partial r} (r \rho) \quad (4)$$

To reduce the 2 velocities into one variable and to eliminate the equation of conservation of mass, a stream function (which satisfies the equation of continuity) is defined by

$$v_r = \frac{1}{r \rho} \frac{\partial \psi}{\partial \theta} \quad (5)$$

$$v_\theta = - \frac{1}{\rho} \frac{\partial \psi}{\partial r} \quad (6)$$

With the above definitions of the stream function, and with $g_r = g \cos \theta$, $g_\theta = -g \sin \theta$, the equation of conservation of momentum becomes

$$\frac{1}{r K_r} \frac{\partial}{\partial \theta} \left(\nu \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{K_\theta} \frac{\partial}{\partial r} \left(r \nu \frac{\partial \psi}{\partial r} \right) = g (\cos \theta \frac{\partial \rho}{\partial \theta} + r \sin \theta \frac{\partial \rho}{\partial r}) \quad (7)$$

Conservation of Energy

The equation of conservation of energy is:

$$\frac{\partial}{\partial r} (r \rho c_p v_r T) + \frac{\partial}{\partial \theta} (\rho c_p v_\theta T) - \frac{\partial}{\partial r} (r \lambda \frac{\partial T}{\partial r}) - \frac{1}{r} \frac{\partial}{\partial \theta} (\lambda \frac{\partial T}{\partial \theta}) = 0 \quad (8)$$

Using the definitions of the stream function and using the equation of conservation of mass, one obtains:

$$\frac{\partial \psi}{\partial \theta} \frac{\partial}{\partial r} (c_p T) - \frac{\partial \psi}{\partial r} \frac{\partial}{\partial \theta} (c_p T) - \frac{\partial}{\partial r} \left(r \lambda \frac{\partial T}{\partial r} \right) - \frac{1}{r} \frac{\partial}{\partial \theta} \left(\lambda \frac{\partial T}{\partial \theta} \right) = 0 \quad (9)$$

Thus, the original group of governing equations has been reduced to two equations with two dependent variables ψ and T . The coefficients ρ , ν , c_p , and λ are variable coefficients which are functions of T and the system pressure.

Boundary Conditions

Since the governing differential equations (7) and (9) are each of second order in ψ and T , these equations are subject to two boundary conditions on the stream function and two on the temperature. For the problem shown in Figure 1, the boundary conditions are:

$$\text{At } \theta = 0 \text{ and } \theta = \Pi$$

$$\psi = 0 \quad (v_\theta = 0 \text{ because of the symmetry})$$

$$\frac{\partial T}{\partial \theta} = 0 \quad (\text{no heat flux because of symmetry})$$

$$\text{At } r = r_1 ; \text{ there are two possible boundary conditions:}$$

$$(a) \quad T = T_1(\theta)$$

$$(b) \quad h_1 (T_g - T) = -\lambda \frac{\partial T}{\partial r}$$

Condition (a) corresponds to a given inside wall temperature. Condition (b) is a convective boundary condition. T_g is the bulk temperature of the gas flowing through the inner pipe, and h_1 is the heat transfer coefficient. The temperature drop across the wall is neglected. For high values of h_1 , the temperature drop ($T_g - T$) will be small; in this case condition (b) is equivalent to condition (a), with $T = T_1(\theta) = T_g$.

At $r = r_2$, again two possible boundary conditions are considered:

$$(a) \quad T = T_2(\theta)$$

$$(b) \quad h_2 (T - T_0) = -\lambda \frac{\partial T}{\partial r}$$

Condition (a) corresponds to a given outside wall temperature. Condition (b) is again a convective boundary condition. This boundary condition is valid when there is no outside insulation. T_0 is the temperature of the gas outside (usually atmospheric air) and h_2 is the heat transfer coefficient (usually natural convection heat transfer coefficient for a horizontal cylinder). Condition (b) allows for the angular variation of the outside wall temperature.

2.2 - Vertical Cylinder

The equations and boundary conditions for flow and heat transfer for insulation pecked

between two vertical cylinders shown in Figure 2 are written in cylindrical coordinates r and x . For this geometry, velocity and temperature fields do not change with the angular coordinate θ , therefore a two dimensional modeling with the radial and axial coordinates is sufficient.

Conservation of Mass

$$\frac{1}{r} \frac{\partial}{\partial r} (\rho r v_r) + \frac{\partial}{\partial x} (\rho v_x) = 0 \quad (10)$$

Conservation of Momentum

Darcy's law, with velocity components v_r and v_x , is:

$$v_r = -\frac{K_r}{\mu} \left(\rho g_r - \frac{\partial p}{\partial r} \right) \quad (11)$$

$$v_x = -\frac{K_x}{\mu} \left(\rho g_x - \frac{\partial p}{\partial x} \right) \quad (12)$$

Eliminating the pressure, and noting that $g_r = 0$ and $g_x = -g$, one obtains

$$\frac{1}{K_r} \frac{\partial}{\partial x} (\mu v_r) - \frac{1}{K_x} \frac{\partial}{\partial r} (\mu v_x) = g \frac{\partial \rho}{\partial r} \quad (13)$$

A stream function is defined by

$$v_r = \frac{1}{\rho r} \frac{\partial \psi}{\partial x} \quad (14)$$

$$v_x = -\frac{1}{\rho r} \frac{\partial \psi}{\partial r} \quad (15)$$

Combining Equations (13), (14), and (15) one obtains

$$\frac{1}{K_r} \frac{1}{r} \frac{\partial}{\partial x} \left(\nu \frac{\partial \psi}{\partial x} \right) + \frac{1}{K_x} \frac{\partial}{\partial r} \left(\frac{\nu}{r} \frac{\partial \psi}{\partial r} \right) = g \frac{\partial \rho}{\partial r} \quad (16)$$

Conservation of Energy

The equation of conservation of energy, written in two dimensional cylindrical coordinates with coordinates r and x is:

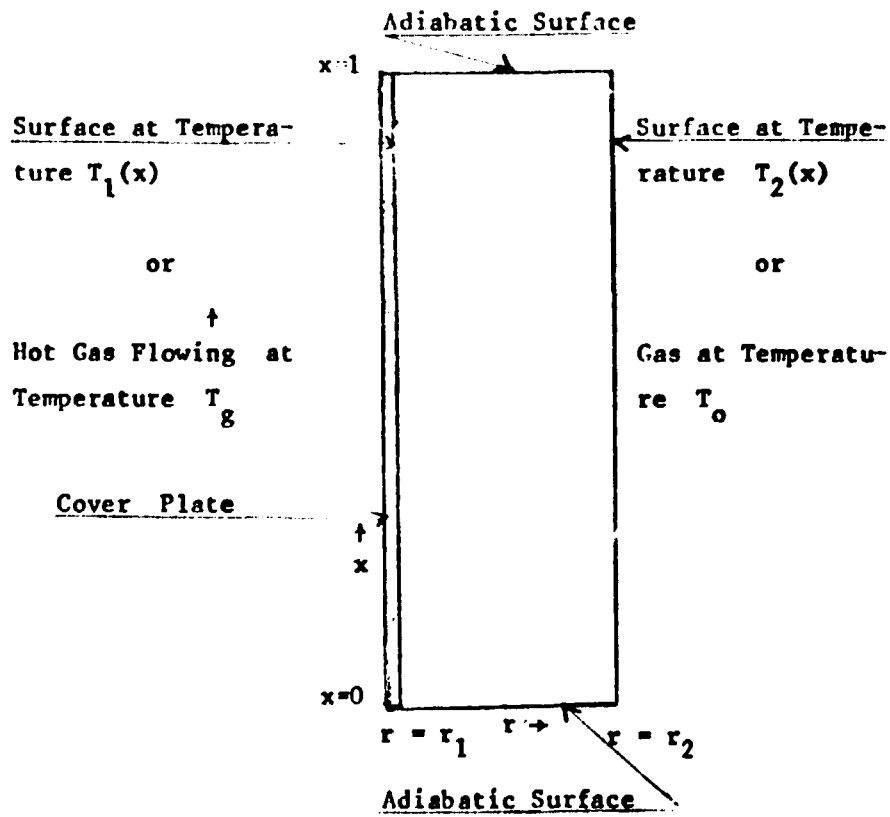


Figure 2 - Vertical Cylinder geometry for which the boundary conditions are described

$$\frac{\partial}{\partial r} (r v_r \rho c_D T) + \frac{\partial}{\partial x} (r v_x \rho c_D T) - \frac{\partial}{\partial r} (r \lambda \frac{\partial T}{\partial r}) - r \frac{\partial}{\partial x} (\lambda \frac{\partial T}{\partial x}) = 0 \quad (17)$$

Introducing the stream function and using the equation of conservation of mass, one obtains

$$\frac{\partial \psi}{\partial x} \frac{\partial}{\partial r} (C_p T) - \frac{\partial \psi}{\partial r} \frac{\partial}{\partial x} (C_p T) - \frac{\partial}{\partial r} (r \lambda \frac{\partial T}{\partial r}) - r \frac{\partial}{\partial x} (\lambda \frac{\partial T}{\partial x}) = 0 \quad (18)$$

Boundary Conditions

The boundary conditions for a vertical cylinder include the possibility of permeation flow through the perforated inner wall of a double-pipe system. The cover plate is considered as a porous medium with own permeability K_c , and thickness a . The boundary conditions are:

$$x = 0 \text{ and } x = L;$$

$$\psi = 0 \quad (v_x = 0, \text{ solid wall})$$

$$\frac{\partial T}{\partial x} = 0 \quad (\text{no heat flow through the solid wall})$$

$$r = r_1;$$

$$(a) \psi = 0 \quad (v_r = 0, \text{ solid inner wall without perforations})$$

or

$$(b) \frac{\partial \psi}{\partial r} + \frac{aK}{K_c} \frac{\partial^2 \psi}{\partial x^2} = \frac{rK}{\nu} (g\rho + \frac{\partial p_0}{\partial x}) \quad (19)$$

Boundary condition (b) corresponds to permeation flow through the perforated well, where $\frac{\partial p_0}{\partial x}$ is the pressure gradient in the inner pipe due to the flow of the gas.

There also two possible boundary conditions on T :

$$(a) T = T_1(x) \quad (\text{given inner well temperature})$$

or

$$(b) h_1 (T_g - T) = -\lambda \frac{\partial T}{\partial r} \quad (\text{convective boundary condition as in the horizontal cylinder})$$

$$r = r_2$$

$$\psi = 0 \quad (v_r = 0, \text{ solid outer wall})$$

$$(a) \quad T = T_2(x) \quad (\text{given outer wall temperature})$$

or

$$(b) \quad h_2 (T - T_0) = -\lambda \frac{\partial T}{\partial r} \quad (\text{convective boundary condition as in the horizontal cylinder})$$

Equation (19) is obtained as follows:

Darcy's law applied to the cover plate ($r = r_1$) gives

$$v_r = \frac{K_c}{\mu} \left(\rho g_r - \frac{\partial p}{\partial r} \right) \quad (20)$$

With $g_r = 0$ and $\frac{\partial p}{\partial r} = \frac{p - p_0}{a}$, Equation (20) gives

$$p = p_0 - \frac{a/\nu_r}{K_c} \quad (21)$$

Differentiating Equation (21) with respect to x and using definition of the stream function, one obtains (at $r = r_1$)

$$\frac{\partial p}{\partial x} = \frac{\partial p_0}{\partial x} - \frac{a}{K_c} \frac{\partial}{\partial x} \left(-\frac{\nu}{r} \frac{\partial \psi}{\partial x} \right) \quad (22)$$

Writing the x - momentum equation at $r = r_1$ gives

$$v_x = \frac{K}{\mu} \left(\rho g_x - \frac{\partial p}{\partial x} \right) \quad (23)$$

With $g_x = -g$ and using the definition of the stream function, one obtains for $\frac{\partial \psi}{\partial r}$

$$\frac{\partial \psi}{\partial x} = \frac{rK}{\nu} \left(g + \frac{p}{x} \right) \quad (24)$$

Substituting $\frac{\partial p}{\partial x}$ from Equation (22) into Equation (24), Equation (19) is derived.

2.3 - Rectangular Geometries

The equations and boundary conditions for the flow and heat transfer in insulation packed in a rectangular cavity are written considering the rectangular cross section shown in Figure 3.

Equations are formulated in two dimensional rectangular coordinates x and y , where x is along the cover plate. The equations are the following.

Conservation of Mass

$$\frac{\partial}{\partial x} (\rho v_x) + \frac{\partial}{\partial y} (\rho v_y) = 0 \quad (25)$$

Conservation of Momentum

$$U_x = \frac{K_x}{\mu} (\rho g_x - \frac{\partial p}{\partial x}) \quad (26)$$

$$U_y = \frac{K_y}{\mu} (\rho g_y - \frac{\partial p}{\partial y}) \quad (27)$$

Eliminating the pressure gives

$$\frac{1}{K_x} \frac{\partial}{\partial y} (\mu v_x) - \frac{1}{K_y} \frac{\partial}{\partial x} (\mu v_y) = g_x \frac{\partial \rho}{\partial y} - g_y \frac{\partial \rho}{\partial x} \quad (28)$$

The stream function is defined by

$$v_x = - \frac{1}{\rho} \frac{\partial \psi}{\partial y} \quad (29)$$

$$v_y = \frac{1}{\rho} \frac{\partial \psi}{\partial x} \quad (30)$$

Combining Equations (28), (29), and (30), one obtains

$$\frac{1}{K_x} \frac{\partial}{\partial y} (\nu \frac{\partial \psi}{\partial y}) + \frac{1}{K_y} \frac{\partial}{\partial x} (\nu \frac{\partial \psi}{\partial x}) = g_y \frac{\partial \rho}{\partial x} - g_x \frac{\partial \rho}{\partial y} \quad (31)$$

Conservation of Energy

$$\frac{\partial}{\partial x} (v_x \rho C_p T) + \frac{\partial}{\partial y} (v_y \rho C_p T) - \frac{\partial}{\partial x} (\lambda \frac{\partial T}{\partial x}) - \frac{\partial}{\partial y} (\lambda \frac{\partial T}{\partial y}) = 0 \quad (32)$$

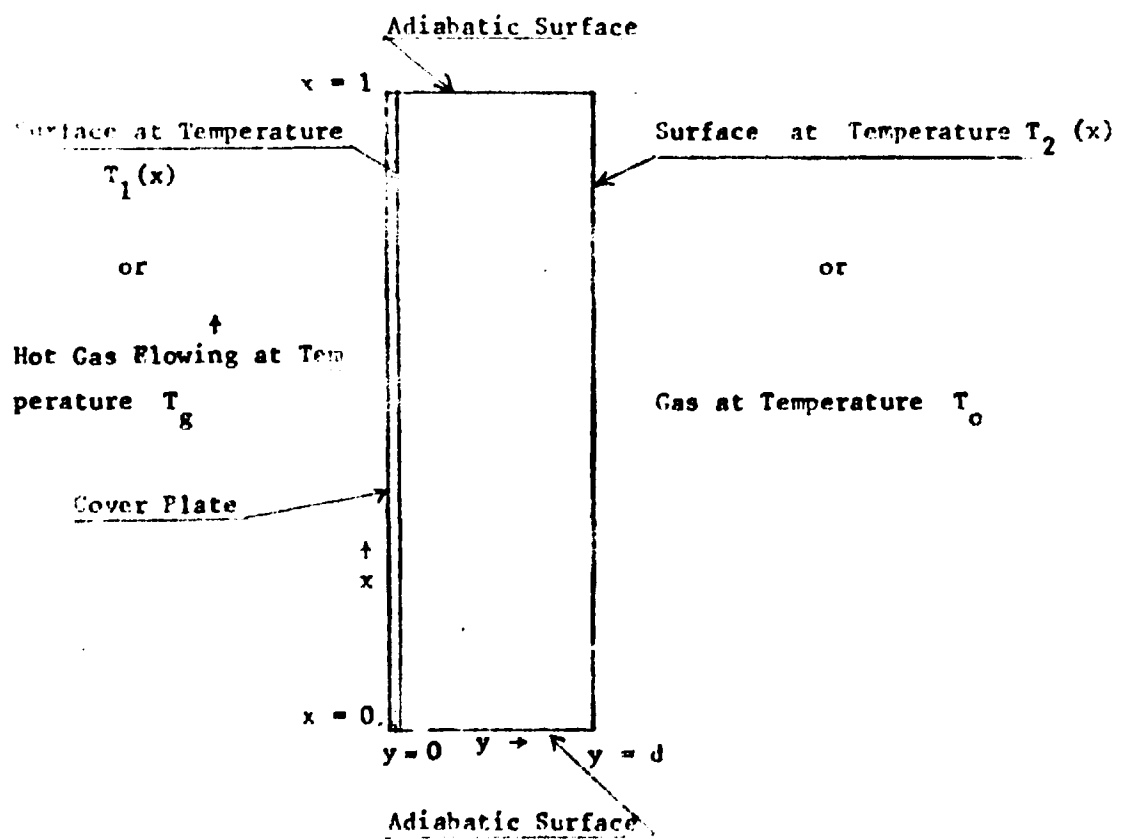


Figure 3 -- Rectangular Geometry for which the boundary conditions are described

Combining Equations (25), (26), (27), and (32), results in equation of conservation of energy

$$\frac{\partial \psi}{\partial y} \frac{\partial}{\partial x} (C_p T) - \frac{\partial \psi}{\partial x} \frac{\partial}{\partial y} (C_p T) = \frac{\partial}{\partial x} \left(\lambda \frac{\partial T}{\partial x} \right) - \frac{\partial}{\partial y} \left(\lambda \frac{\partial T}{\partial y} \right) = 0 \quad (33)$$

Boundary Conditions

Boundary conditions are similar to those written for the vertical cylinder.

At $x = 0$ and $x = 1$;

$$\psi = 0 \quad (v_x = 0, \text{ solid wall})$$

$$\frac{\partial T}{\partial x} = 0 \quad (\text{no heat flux through the solid wall})$$

At $y = 0$

$$\psi = 0 \quad (v_y = 0, \text{ cover plate without perforations})$$

or

$$\frac{\partial \psi}{\partial y} = -\frac{K}{\nu} \left(g_x - \frac{\partial \rho_o}{\partial x} \right) + \frac{K a g_y}{\nu} - \frac{a K}{K_c} \frac{\partial^2 \psi}{\partial x^2} \quad (34)$$

$$T = T_1(x) \quad (\text{given cover plate temperature})$$

or

$$h_1 (T_g - T) = -\frac{\partial T}{\partial y} \quad (\text{convective boundary condition})$$

At $y = 1$

$$\psi = 0 \quad (v_y = 0, \text{ solid wall})$$

$$T = T_2(x) \quad (\text{given outer wall temperature})$$

or

$$h_2 (T - T_0) = -k \frac{\partial T}{\partial y} \quad (\text{convective boundary condition})$$

Equation (34) is derived as follows:

Darcy's law across the cover plate gives

$$\frac{\mu v_y}{K_c} = \rho g_y - \frac{\partial p}{\partial y} \quad (35)$$

or

$$\frac{\mu v_y}{K_c} = \rho g_y + \frac{P_0 - P}{a} \quad (36)$$

Differentiating Equation (36) with respect to x and rearranging, one obtains

$$\frac{\partial p}{\partial x} = \frac{\partial p_0}{\partial x} + a \frac{\partial}{\partial x} \left(\rho g_y - \frac{\mu v_y}{K_c} \right) \quad (37)$$

Writing the x -momentum equation at $y = 0$

$$v_x = \frac{K}{\nu} \left(g_y - \frac{\partial p}{\partial x} \right) \quad (38)$$

Substituting $\frac{\partial p}{\partial x}$ from Equation (37) into Equation (38) gives the permeation flow boundary condition on ψ .

3 – NUMERICAL SOLUTION

3.1 – Horizontal Cylinder

The coupled Equations (7) and (8), and the boundary condition equations are solved numerically using a finite difference method implemented by a computer program called FINS (Fibrous Insulation). The numerical procedure is described as follows, considering the grid shown in Figure 4

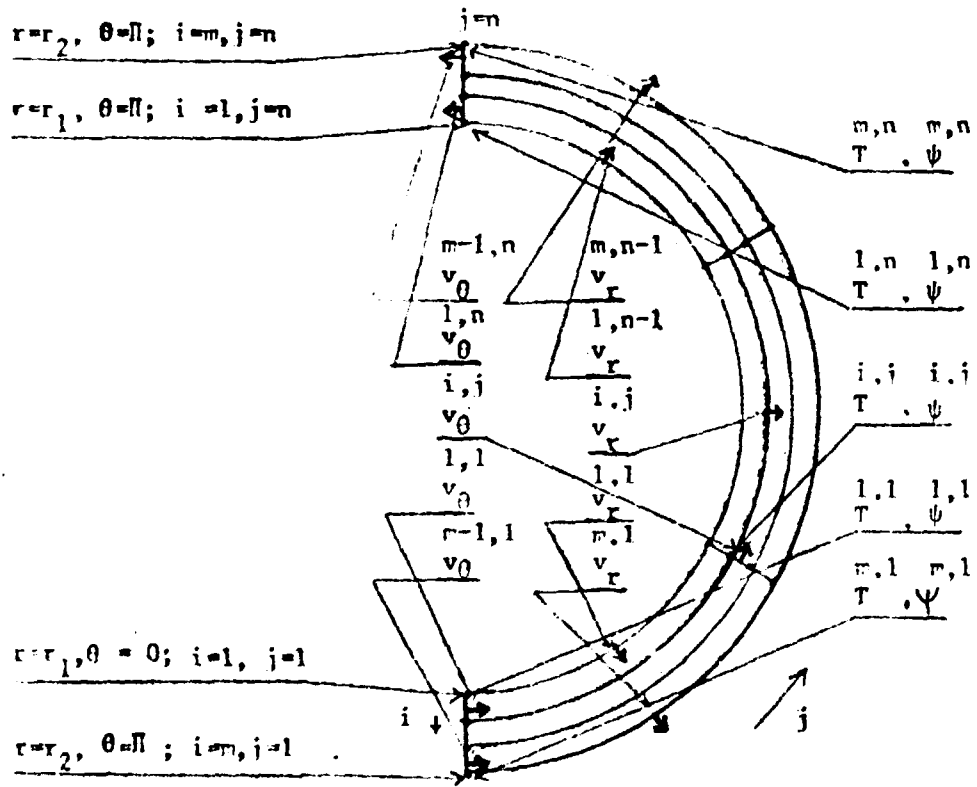


Figure 4 - Horizontal Cylinder, Finite Difference Grid

The derivatives terms in ψ and ρ in Equation (7) are expressed as central differences for all the interior points.

Thus, $(n-2)(m-2)$ algebraic equations are obtained. Boundary conditions expressing that $\psi^{i,j} = 0$ on the periferial points supply $2(n+m)-4$ algebraic equations, thus yielding a total of $n \times m$ equations, with $n \times m$ stream function $\psi^{i,j}$ variables. When the gas viscosities $\nu^{i,j}$ and densities $\rho^{i,j}$ are known at each grid point, Equation (7), which is a non-linear partial differential equation, is reduced to a system of linear algebraic equations which can be solved readily by numerical methods.

Similarly, Equation (9), which is also a non-linear partial differential equation, is put in a finite difference form using central differences for T and ψ derivatives at the interior points. For the boundary conditions, $\frac{\partial T}{\partial \theta}$ is written as a forward difference, and $\frac{\partial T}{\partial r}$ at $r=r_1$ and $r=r_m$ is obtained by differentiating the equation of the parabola passing through the 3 points next to the walls. Thus, when the gas properties $C_p^{i,j}$ and $\lambda^{i,j}$, and the stream function $\psi^{i,j}$ are known, Equation (9) is reduced to a system of $n \times m$ linear algebraic equations with $n \times m$ temperature variables $T^{i,j}$. The convergence is achieved when the temperature variations are within a set error criterion. The details of the procedure are written below.

Momentum Equation

The finite difference form to Equation (7) is:

$$\begin{aligned} & \frac{\nu^{i,j-1} + \nu^{i,j}}{2K_r r^i (\Delta\theta)^2} \psi^{i,j-1} + \frac{\nu^{i,j} + \nu^{i,j+1}}{2K_r r^i (\Delta\theta)^2} \psi^{i,j+1} \\ & + \frac{r^{i-1} \nu^{i-1,j} + r^i \nu^{i,j}}{2K_\theta (\Delta r)^2} \psi^{i-1,j} + \frac{r^i \nu^{i,j} + r^{i+1} \nu^{i+1,j}}{2K_\theta (\Delta r)^2} \psi^{i+1,j} \\ & - \left[\frac{\nu^{i,j-1} + 2\nu^{i,j} + \nu^{i,j+1}}{2K_r r^i (\Delta\theta)^2} + \frac{r^{i-1} \nu^{i-1,j} + r^i \nu^{i,j} + r^{i+1} \nu^{i+1,j}}{2K_\theta (\Delta r)^2} \right] \psi^{i,j} \\ & = g \frac{\cos \theta^j}{2\Delta\theta} (\rho^{i,j+1} - \rho^{i,j-1}) + \frac{r^i \sin \theta^j}{2\Delta r} (\rho^{i+1,j} - \rho^{i-1,j}) \end{aligned} \quad (39)$$

with $i = 2$ to $m-1$, and $j = 2$ to $n-1$.

The boundary condition equations are:

$$\psi^{1,j} = 0$$

$$j = 1, n$$

$$\psi^{m,j} = 0$$

$$\psi^{i,1} = 0$$

$$i = 2, m-1$$

$$\psi^{i,n} = 0$$

Energy Equation

The finite difference form of Equation (9) is:

$$\begin{aligned} & \left[\frac{(\psi^{i+1,j} - \psi^{i-1,j}) C_p^{i,j-1}}{4 \Delta \theta \Delta r} - \frac{\lambda^{i,j} + \lambda^{i,j-1}}{2r^i (\Delta \theta)^2} \right] T^{i,j-1} \\ & - \left[\frac{(\psi^{i+1,j} - \psi^{i-1,j}) C_p^{i,j+1}}{4 \Delta \theta \Delta r} + \frac{\kappa^{i,j+1} + \kappa^{i,j}}{2r^i (\Delta \theta)^2} \right] T^{i,j+1} \\ & - \left[\frac{(\psi^{i,j+1} - \psi^{i,j-1}) C_p^{i-1,j}}{4 \Delta \theta \Delta r} + \frac{r^i \lambda^{i,j} + r^{i-1} \lambda^{i-1,j}}{2 (\Delta r)^2} \right] T^{i-1,j} \\ & + \left[\frac{(\psi^{i,j+1} - \psi^{i,j-1}) C_p^{i+1,j}}{4 \Delta \theta \Delta r} - \frac{r^{i+1} \lambda^{i+1,j} + r^i \lambda^{i,j}}{2 (\Delta r)^2} \right] T^{i+1,j} \\ & + \left[\frac{r^{j+1} \lambda^{i+1,j} + 2r^j \lambda^{i,j} + r^{j-1} \lambda^{i-1,j}}{2 (\Delta r)} + \frac{\lambda^{i,j+1} + 2\lambda^{i,j} + \lambda^{i,j-1}}{2r^i (\Delta \theta)^2} \right] T^{i,j} = 0 \quad (42) \end{aligned}$$

The boundary condition equations are:

$$T^{1,j} = T_1^j \quad (\text{given inner wall temperature}) \quad (43-a)$$

or

$$\left(h_1 + \frac{3\lambda^{1,j}}{2\Delta r} \right) T^{1,j} - \frac{2\lambda^{1,j}}{\Delta r} T^{2,j} + \frac{\lambda^{1,j}}{2\Delta r} T^{3,j} = h_1 T_0^j \quad (\text{convective boundary condition}) \quad (43-b)$$

$$T^{m,j} = T_2^j \quad (\text{given outer wall temperature}) \quad (44-a)$$

or

$$\left(h_2 + \frac{3\lambda^{m,j}}{2\Delta r} \right) T^{m,j} - \frac{2\lambda^{m,j}}{\Delta r} T^{m-1,j} + \frac{\lambda^{m,j}}{2\Delta r} T^{m-2,j} = h_2 T_0^j \quad (\text{convective boundary condition}) \quad (44-b)$$

with $j = 1$ to m ,

$$\begin{aligned}
 T^{1,2} - T^{1,1} &= 0 \\
 T^{1,n} - T^{1,n-1} &= 0
 \end{aligned}
 \tag{45}$$

Solution of the set of linear algebraic equations

The solution starts with initial guesses of the temperatures. In the computer program, initial $T^{i,j}$'s are set equal to the mean temperature \bar{T} of the insulation, which is taken as the arithmetic average of the cold and hot wall temperatures. The gas properties $\nu^{i,j}$ and $\rho^{i,j}$ are evaluated at \bar{T} and at the system pressure, using empirical equations to predict the properties. With numerical values of $\nu^{i,j}$ and $\rho^{i,j}$, Equation (39) is linearized.

To solve the set of $n \times m$ linear algebraic equations made up by Equations (39), (40) and (41), the coefficient matrix is formed in such a way as to obtain a band structure. The $n \times m$ grid points $z^{i,j}$ are reindexed as a one dimensional array z^k , with $z^1 = z^{1,1}$, $z^2 = z^{2,1}$, or $z^{(j-1)m + i} = z^{i,j}$. The variables $T^{i,j}$ are reindexed into T^K in a similar manner. The equations are numbered so that the first equation is written at the point z^1 , second at z^2 , etc. In this way, the coefficient matrix has a band structure, with m upper and m lower codiagonals, and it is stored in a one dimensional array of size $s = m \times n (2m + 1) - \frac{m(m + 1)}{2}$. The size of the general coefficient matrix for the same system would be $(nm)^2$. The band structure thus allows the solution of the large sets of linear equations necessary to obtain accurate numerical results.

For the solution of the equations, subroutine GELB from IBM SSP (scientific Subroutine Package) is used. In GELB, solution is done by means of Gauss elimination with column pivoting only, in order to preserve band structure in remaining coefficient matrices.

The structure of the coefficient matrix of the momentum equation and boundary conditions for a grid of 4 x 4 nodes is shown in Figure 5. Energy equation has a similar matrix structure.

The first solution of the momentum equation, with $T^{i,j} = \bar{T}$, results in $\psi^{i,j} = 0$ since $\frac{\partial \rho^{i,j}}{\partial r} = \frac{\partial \rho^{i,j}}{\partial \theta} = 0$. Next, the energy equation is solved, and with $\psi^{i,j} = 0$ it yields the solution of the conduction problem, with the gas properties evaluated at \bar{T} . The properties $\nu^{i,j}$ and $\rho^{i,j}$ for the next solution of the momentum equation are evaluated at $T^{i,j}$'s obtained by a weighted average of the last two set of temperatures. For the next solution of the energy equation, a weighted average of the last two $\psi^{i,j}$'s and gas properties evaluated at $T^{i,j}$ are used. The procedure is continued until convergence of temperatures is obtained. For a grid at 20 x 20, at Ra numbers smaller than 1000, the temperatures converged within 0.1°C after about 15 iterations. At 3000 > Ra > 1000, the temperatures fluctuated around the solution by about 1°C. Approximately at Ra numbers greater than 3000, convergence difficulties were encountered. The convergence should improve at finer grid sizes.

The solution of a problem with 20 x 20 nodes required a storage at 200 K (about 130 K to store the coefficient matrix in double precision), and 5 minutes of CPU time for 15 iterations on the I.E.A.'s IBM 370/55.

The flow chart of the computer program FINS is shown in Figure 6.

	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16
1	X															
2		X														
3			X													
4				X												
5					X											
6		X			X	X	X			X						
7			X			X	X	X				X				
8								X								
9									X							
10						X			X	X	X				X	
11										X	X	X				
12												X				
13													X			
14														X		
15															X	
16																X

Figure 5 - Coefficient Matrix for a grid of 4 x 4 nodes for the Momentum Equation and the Boundary Conditions. x's are non-zero elements.

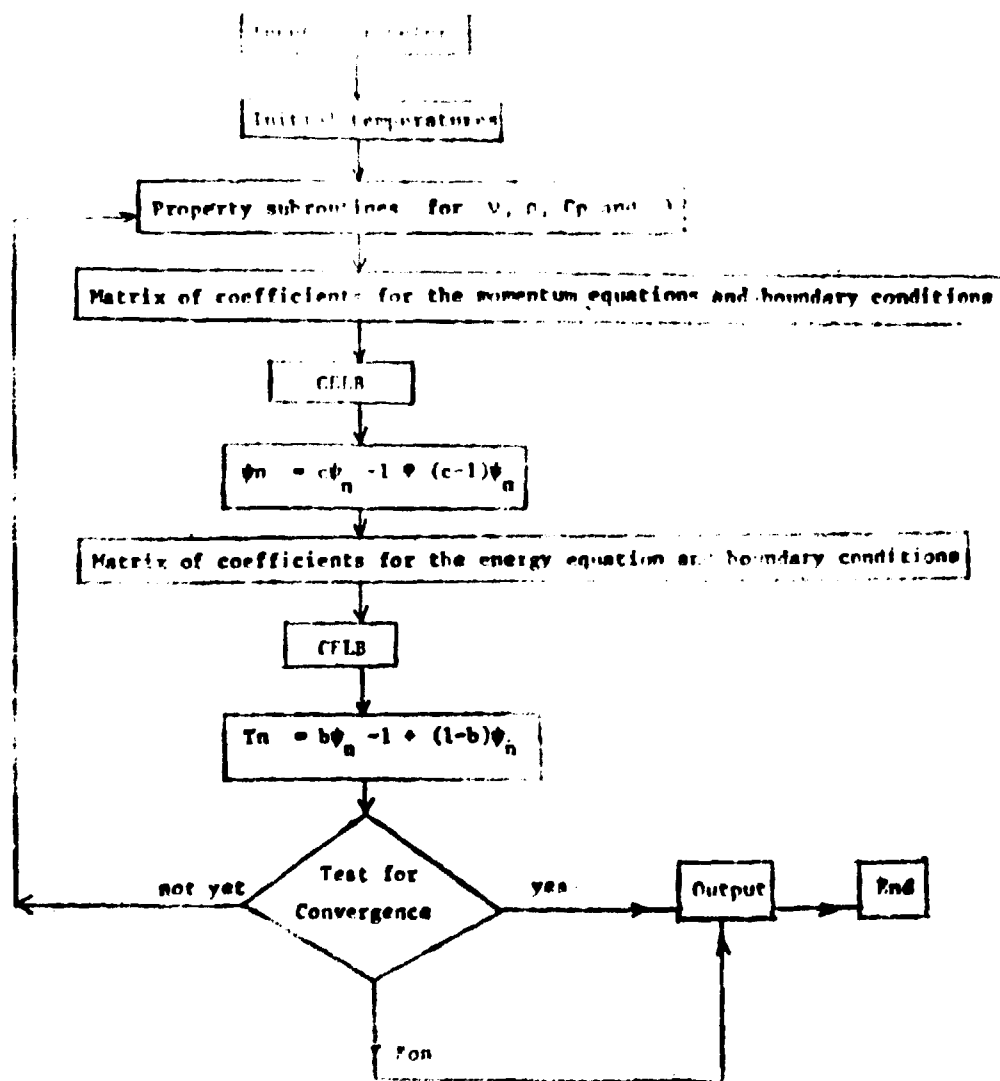


Figure 8. Flow Chart for FIMS.

For the weight factors a and b several combinations have been tried. For small Ra numbers, $a = b = 0$ gives rapid convergence. At high Ra numbers, $a \geq .5$, $b \geq 0$ is necessary. The final version of FINS is with $a = .5$, and $b = .0$.

3.2 – Vertical Cylinder

The finite difference equations for the vertical cylinder are derived similarly to horizontal cylinder (Figure 7).

Momentum Equation (Equation (16))

$$\begin{aligned}
 & \frac{\nu^{i,j} + \nu^{i,j-1}}{2K_r r^i (\Delta x)^2} \psi^{i,j-1} + \frac{\nu^{i,j+1} + \nu^{i,j}}{2K_r r^i (\Delta x)^2} \psi^{i,j+1} \\
 & + \frac{\frac{\nu^{j,j} + \nu^{j-1,j}}{r^j} + \frac{\nu^{j+1,j}}{r^{j+1}}}{2K_x (\Delta r)^2} \psi^{j-1,j} + \frac{\frac{\nu^{j+1,j} + \nu^{j,j}}{r^{j+1}} + \frac{\nu^{j,j}}{r^j}}{2K_x (\Delta r)^2} \psi^{j+1,j} \\
 & - \left[\frac{\nu^{i,j+1} + 2\nu^{i,j} + \nu^{i,j-1}}{2K_r r^i (\Delta x)^2} + \frac{\frac{\nu^{j+1,j}}{r^{j+1}} + 2\frac{\nu^{j,j}}{r^j} + \frac{\nu^{j-1,j}}{r^{j-1}}}{2K_x (\Delta r)^2} \right] \psi^{i,j} \\
 & = \frac{g}{2 \Delta r} (\rho^{i+1,j} - \rho^{i-1,j})
 \end{aligned} \tag{46}$$

with $i = 2$ to $m-1$ and $j = 2$ to $n-2$.

Boundary Conditions

Equations corresponding to $\psi = 0$ at $x = 0$ and $x = L$ are:

$$\psi^{1,1} = 0 \tag{47}$$

$$\psi^{1,n} = 0$$

with $i = 1$ to m , and at $r = r_1$

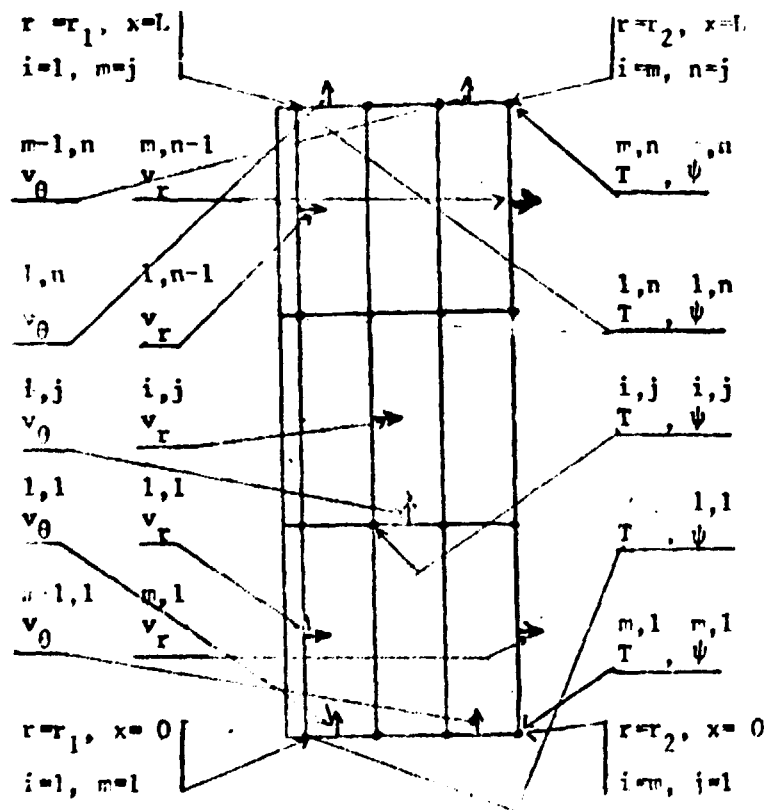


Figure 7 - Vertical Cylinder, Finite Difference Grid

$$\psi^{1,j} = 0 \quad (\text{solid wall}) \quad (48-a)$$

or

$$\begin{aligned} & - \left[\frac{1}{\Delta r} + \frac{2aK}{(\Delta x)^2 K_c} \right] \psi^{1,j} + \frac{1}{\Delta r} \psi^{2,j} + \frac{aK}{(\Delta x)^2 K_c} \psi^{1,j-1} + \frac{aK}{(\Delta x)^2 K_c} \psi^{1,j+1} \\ & - \frac{rK^1}{\rho^{1,j}} \left(\rho^{1,j} + \frac{\partial \rho_o}{\partial x} \right) \end{aligned} \quad (48-b)$$

with $j = 2, n-1$, and at $r = r_2$

$$\psi^{m,i} = 0 \quad (49)$$

with $i = 1, m$

Equation (48b) corresponds to Equation (19) which is used in the case permeation flow.

Energy Equation (Equation (18))

$$\begin{aligned} & \left[\frac{(\psi^{1,j+1} - \psi^{1,j}) C_p^{1,j+1}}{4 \Delta x \Delta r} - \frac{r^j (\lambda^{1,j} + \lambda^{1-1,j})}{2 (\Delta x)^2} \right] T^{1,j-1} \\ & - \left[\frac{(\psi^{1,j+1} - \psi^{1,j}) C_p^{1,j+1}}{4 \Delta x \Delta r} - \frac{r^{j+1} \lambda^{1+1,j} + r^j \lambda^{1,j}}{2 (\Delta r)^2} \right] T^{1,j+1} \\ & - \left[\frac{(\psi^{1,j+1} - \psi^{1,j-1}) C_p^{1-1,j}}{4 \Delta x \Delta r} + \frac{r^j \lambda^{1,j} + r^{j-1} \lambda^{1-1,j}}{2 (\Delta r)^2} \right] T^{1-1,j} \\ & + \left[\frac{(\psi^{1,j+1} - \psi^{1,j-1}) C_p^{1+1,j}}{4 \Delta x \Delta r} - \frac{r^{j+1} \lambda^{1+1,j} + r^j \lambda^{1,j}}{2 (\Delta r)^2} \right] T^{1+1,j} \\ & + \left[\frac{r^{j+1} \lambda^{1+1,j} + 2r^j \lambda^{1,j} + r^{j-1} \lambda^{1-1,j}}{2 (\Delta x)^2} + \frac{r^j (\lambda^{1,j+1} + 2\lambda^{1,j} + \lambda^{1,j-1})}{2 (\Delta x)^2} \right] T^{1,i} = 0 \quad (50) \end{aligned}$$

With $i = 2$ to $m, 1$ and $j = 2$ to $n-1$

Boundary Conditions

$$T^{1,i} = T_1^i \quad (\text{given hot wall temperature}) \quad (51-a)$$

$$T^{m,i} = T_2^i \quad (\text{given cold wall temperature})$$

or

$$\left(h_1 + \frac{3\lambda^{1,i}}{2\Delta r} T^{1,j} - \frac{2\lambda^{1,i}}{\Delta r} T^{2,j} + \frac{\lambda^{1,i}}{2\Delta r} T^{3,j} \right) = h_1 T_g \quad (51-b)$$

$$\left(h_2 + \frac{3\lambda^{m,i}}{2\Delta r} T^{m,j} - \frac{2\lambda^{m,i}}{\Delta r} T^{m-1,j} + \frac{\lambda^{m,i}}{2\Delta r} T^{m-2,j} \right) = h_2 T_o \quad (\text{convective boundary conditions})$$

with $j = 1$ to n , and to express $\frac{\partial T}{\partial x} = 0$ at $x = 0$ and $x = L$:

$$T^{i,2} - T^{i,1} = 0 \quad (52)$$

$$T^{i,n} - T^{i,n-1} = 0$$

with $i = 2, m-1$

The solution of the finite difference equations is similar to that for the horizontal cylinder.

3.3 – Rectangular Geometries**Momentum Equation**

Finite difference form of Equation (31) is (Figure 8):

$$\begin{aligned} & \frac{v^{i,j} + v^{i,j-1}}{2K_y (\Delta x)^2} \psi^{i,j-1} + \frac{v^{i,j+1} + v^{i,j}}{2K_y (\Delta x)^2} \psi^{i,j+1} \\ & + \frac{v^{i,j} + v^{i-1,j}}{2K_x (\Delta y)^2} \psi^{i-1,j} + \frac{v^{i+1,j} + v^{i,j}}{2K_x (\Delta y)^2} \psi^{i+1,j} \end{aligned}$$

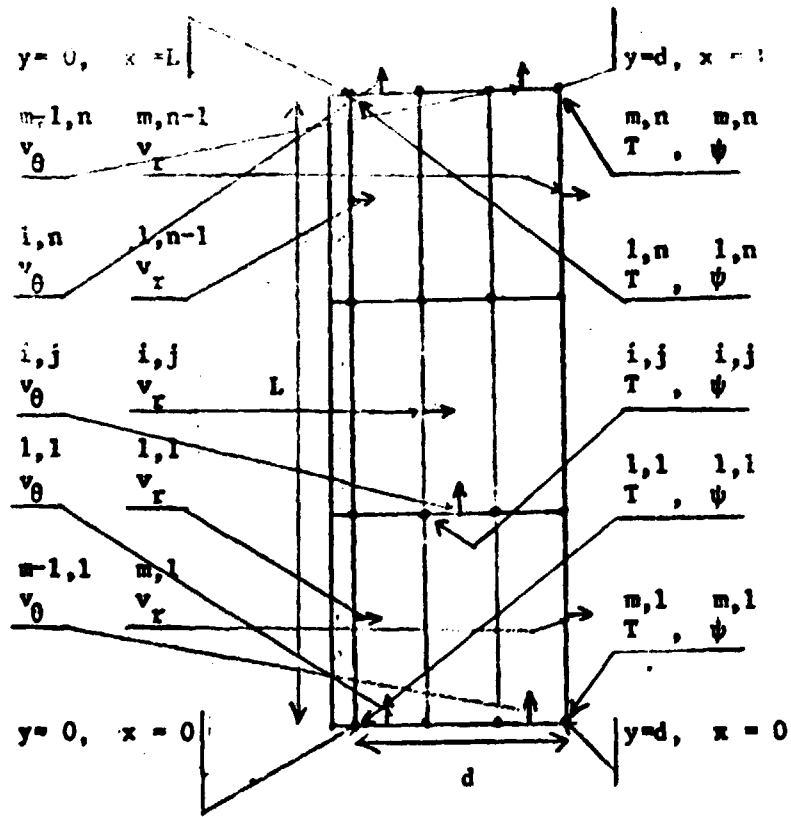


Figure 8 - Rectangular Geometry, Finite Difference Grid

$$- \left[\frac{\nu^{i,j+1} + 2\nu^{i,j} + \nu^{i,j-1}}{2K_y (\Delta x)^2} + \frac{\nu^{i+1,j} + 2\nu^{i,j} + \nu^{i-1,j}}{2K_x (\Delta y)^2} \right] \psi^{i,j} \quad (53)$$

Boundary Conditions

Equation corresponding to $\psi = 0$ at $x = 0$ and $x = L$ are:

$$\psi^{1,1} = 0 \quad (54)$$

$$\psi^{i,n} = 0$$

with $i = 1$ to m , and $r = r_1$

$$\psi^{1,j} = 0 \quad (\text{solid wall}) \quad (55-a)$$

or

$$\begin{aligned} & - \left(\frac{1}{\Delta y} + \frac{2aK}{K_c (\Delta x)^2} \right) \psi^{1,j} + \frac{1}{\Delta y} \psi^{2,j} + \frac{aK}{K_c (\Delta x)^2} \psi^{1,j-1} + \frac{aK}{K_c (\Delta x)^2} \psi^{1,j+1} \\ & = \frac{K}{\nu^{1,j}} \left(-\rho^{1,j} g_x + \frac{\partial p_o}{\partial x} + a g_y \right) \end{aligned} \quad (55-b)$$

with $j = 2$ to $n-1$, and at $r = r_2$

$$\psi^{m,i} = 0 \quad (56)$$

with $i = 1$ to m .

Equation (55-b) corresponds to Equation (34) which is used in the case of permeation flow through the cover plate.

Energy Equation

Finite difference form of Equation (33) is:

$$\left[\frac{(\psi^{i+1,j} - \psi^{i-1,j}) C_p^{i,j-1}}{4 \Delta x \Delta y} - \frac{\lambda^{i,j} + \lambda^{i-1,j}}{2 (\Delta x)^2} \right] T^{i,j-1}$$

$$\begin{aligned}
& - \left[\frac{(\psi^{i,j+1} - \psi^{i,j}) C_p^{i,j+1}}{4 \Delta x \Delta y} + \frac{\lambda^{i,j+1} + \lambda^{i,j}}{2 (\Delta x)^2} \right] T^{i,j+1} \\
& - \left[\frac{(\psi^{i,j+1} - \psi^{i,j-1}) C_p^{i-1,j}}{4 \Delta x \Delta y} + \frac{\lambda^{i,j} + \lambda^{i-1,j}}{2 (\Delta y)^2} \right] T^{i-1,j} \\
& + \left[\frac{(\psi^{i,j+1} - \psi^{i,j-1}) C_p^{i+1,j}}{4 \Delta x \Delta y} - \frac{\lambda^{i+1,j} + \lambda^{i,j}}{2 (\Delta y)^2} \right] T^{i+1,j} \\
& + \left[\frac{\lambda^{i+1,j} + 2\lambda^{i,j} + \lambda^{i-1,j}}{2 (\Delta y)^2} + \frac{\lambda^{i,j+1} + 2\lambda^{i,j} + \lambda^{i,j-1}}{2 (\Delta x)^2} \right] T^{i,j} = 0 \quad (57)
\end{aligned}$$

Boundary Conditions

$$T^{1,j} = T_1^j \quad (\text{given hot wall temperature})$$

(58-a)

$$T^{m,j} = T_2^j \quad (\text{given cold wall temperature})$$

or

$$(h_1 + \frac{3\lambda^{1,j}}{2 \Delta y}) T^{1,j} - \frac{2\lambda^{1,j}}{\Delta y} T^{2,j} + \frac{\lambda^{1,j}}{2 \Delta y} T^{3,j} = h_1 T_0$$

(58-b)

$$(h_2 + \frac{3\lambda^{m,j}}{2 \Delta y}) T^{m,j} - \frac{2\lambda^{m,j}}{\Delta y} T^{m-1,j} + \frac{\lambda^{m,j}}{2 \Delta y} T^{m-2,j} = h_2 T_0 \quad (\text{convective boundary condition})$$

with $j = 1$ to n , and to express $\frac{\partial T}{\partial x} = 0$ at $x = 0$ and $x = L$:

$$T^{i,2} - T^{i,1} = 0$$

(59)

$$T^{i,n} - T^{i,n-1} = 0$$

with $i = 2$ to $m-1$.

The solution of the finite difference equations is similar to that of the horizontal cylinder.

4 – RESULTS

The program FINS has been run for the cylindrical and rectangular geometries with the various boundary conditions discussed earlier. The fluid has been taken as carbon dioxide, with pressures from 1 to 100 bar and mean temperatures from 350°K to 550°K an effective thermal conductivity has been calculated from the velocity and temperature distribution results. The effective Nusselt number has been correlated with the Rayleigh and Peclet number in the case of a permeable hot wall, and with Rayleigh number only in the case of a solid hot wall. The results for the rectangular geometries showed close agreement with a previously published correlation⁽⁵⁾.

4.1 – Horizontal Cylinder

The parameters covered the following range:

Wall temperature T_2 :	316 – 400° K
Hot wall temperature T_1 :	365 – 700° K
Temperature difference $T_1 - T_2$:	50 – 300° K
Pressure:	1.3 – 100 bar
Permeability ($K_\theta = K_r$):	1.E-8 – 1.E-11 m^2
Heat transfer coefficient h_1 :	5 – 1000 $\frac{w}{m^2 \cdot ^\circ K}$
Heat transfer coefficient h_2 :	5 – 1000 $\frac{w}{m^2 K}$
Inside radius r_1 :	.05 – .09 m
Outside radius r_2 :	.075 – .01 m
Insulation thickness $r_2 - r_1$:	.01 – .05
Aspect ratio A:	4.71 – 30.0

Velocity and Temperature Distributions

At values of parameter M ($M = \frac{Ra}{A}$) less than 1, the computations showed that the gas does not circulate. Therefore, the isotherms are concentric cylinders, corresponding to the case of heat transfer by conduction only. At $M > 1$, the gas flows upwards (Figure 9-a) along the hot wall and downward along the cold wall. The corresponding temperature distribution deviates considerably from the conduction problem (Figure 9-b). The mass velocity components m_r and m_θ are shown on Figures 9-c and 9-d.

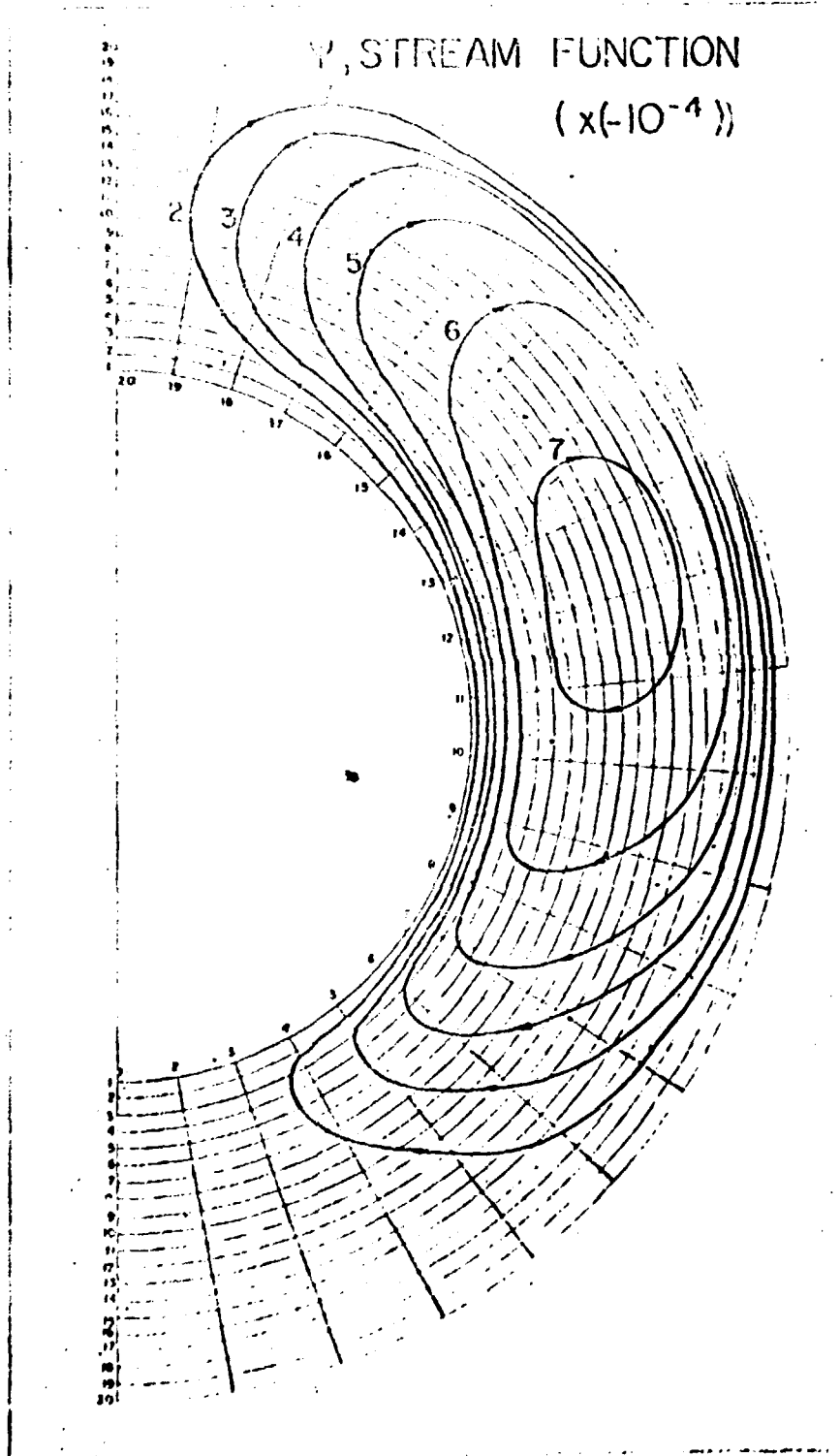


Figure 9-a - Horizontal Cylinder, Stream Lines $Re = 425$, $Nu = 4.83$

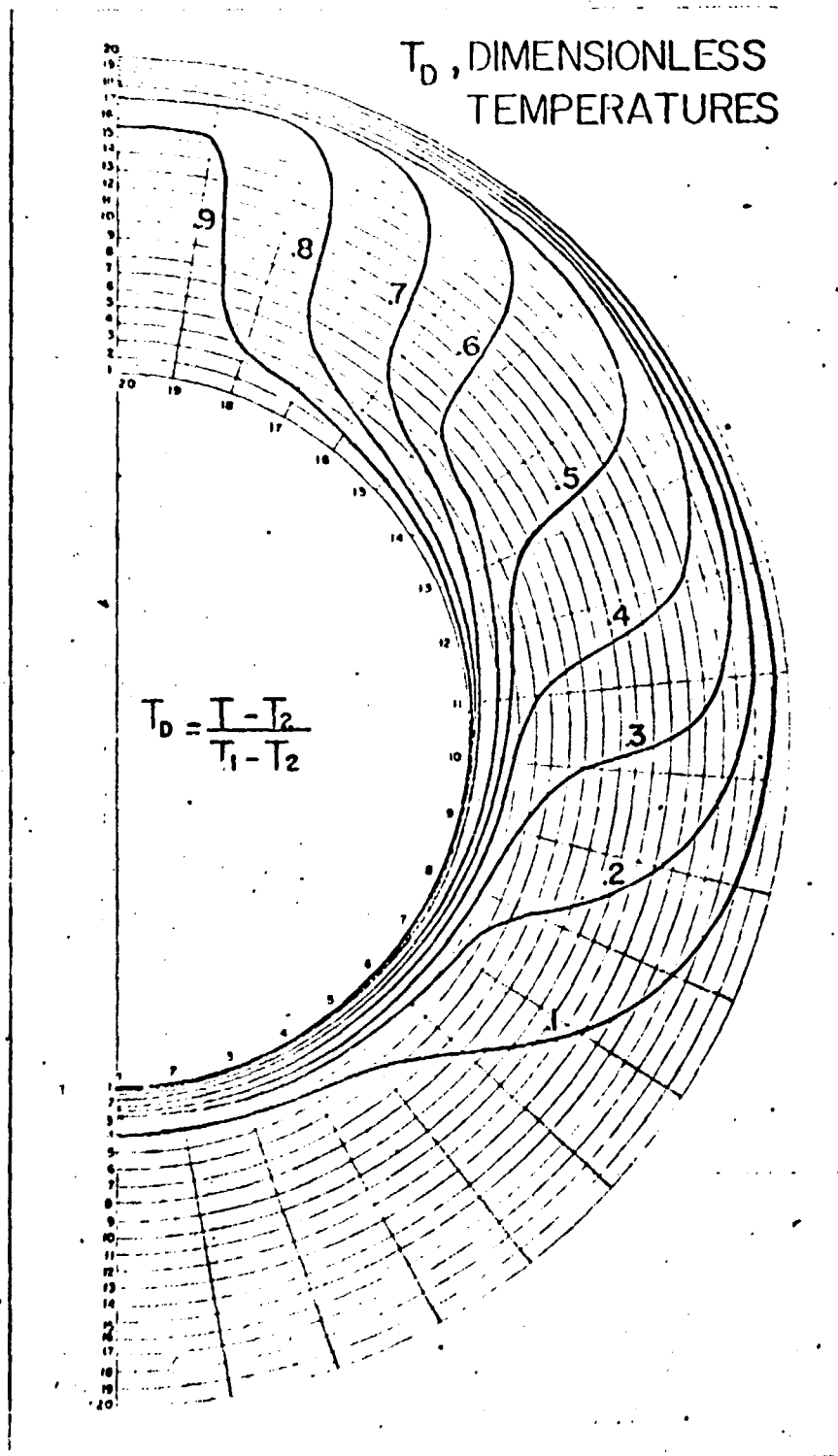
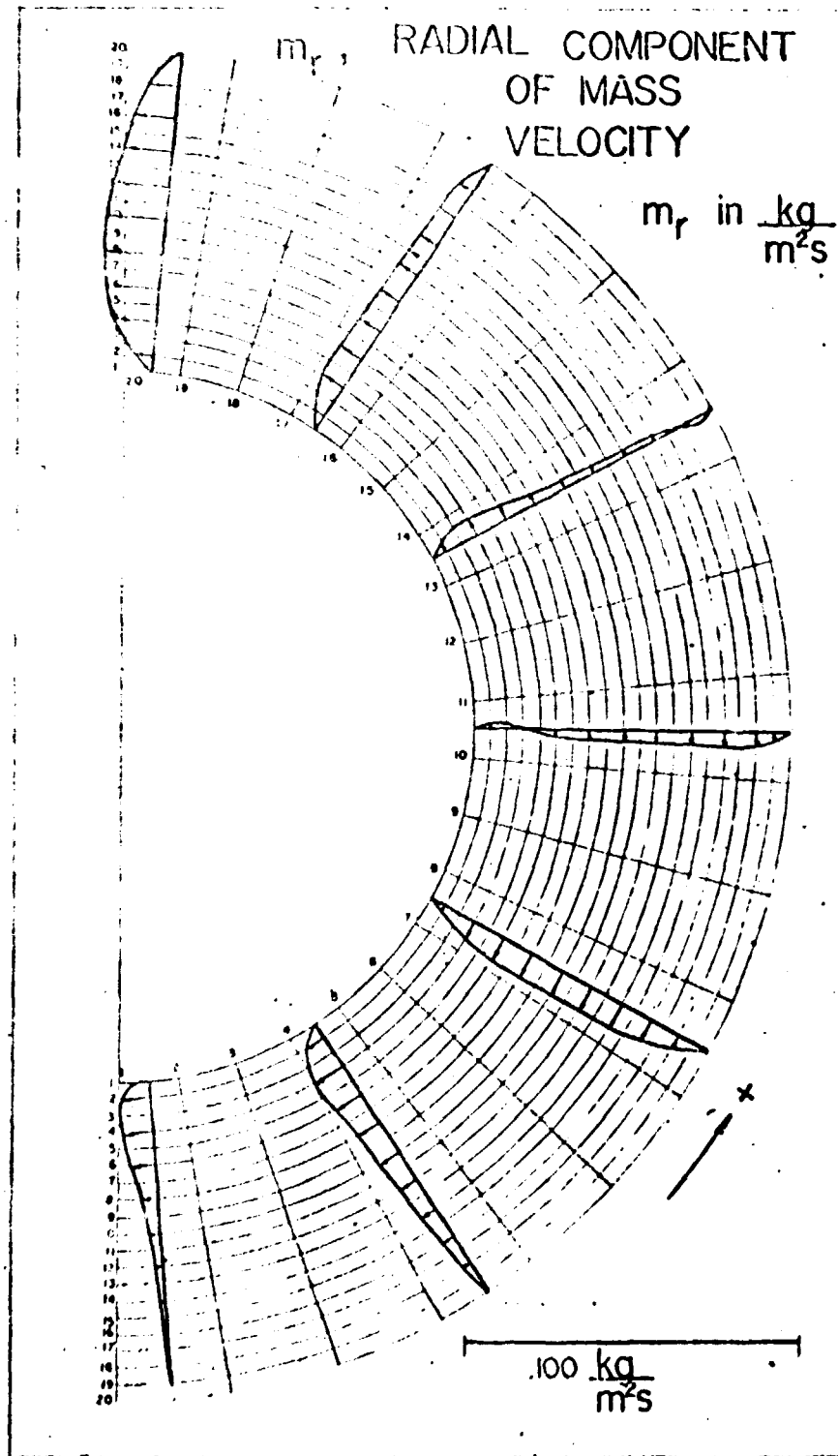


Figure 9 b - Horizontal Cylinder, Dimensionless Temperatures $Ra = 425$ $Nu = 4.83$



—Figure 9-c — Horizontal Cylinder, Radial Component of Mass Velocity, $Re = 425$, $Nu = 4.83$

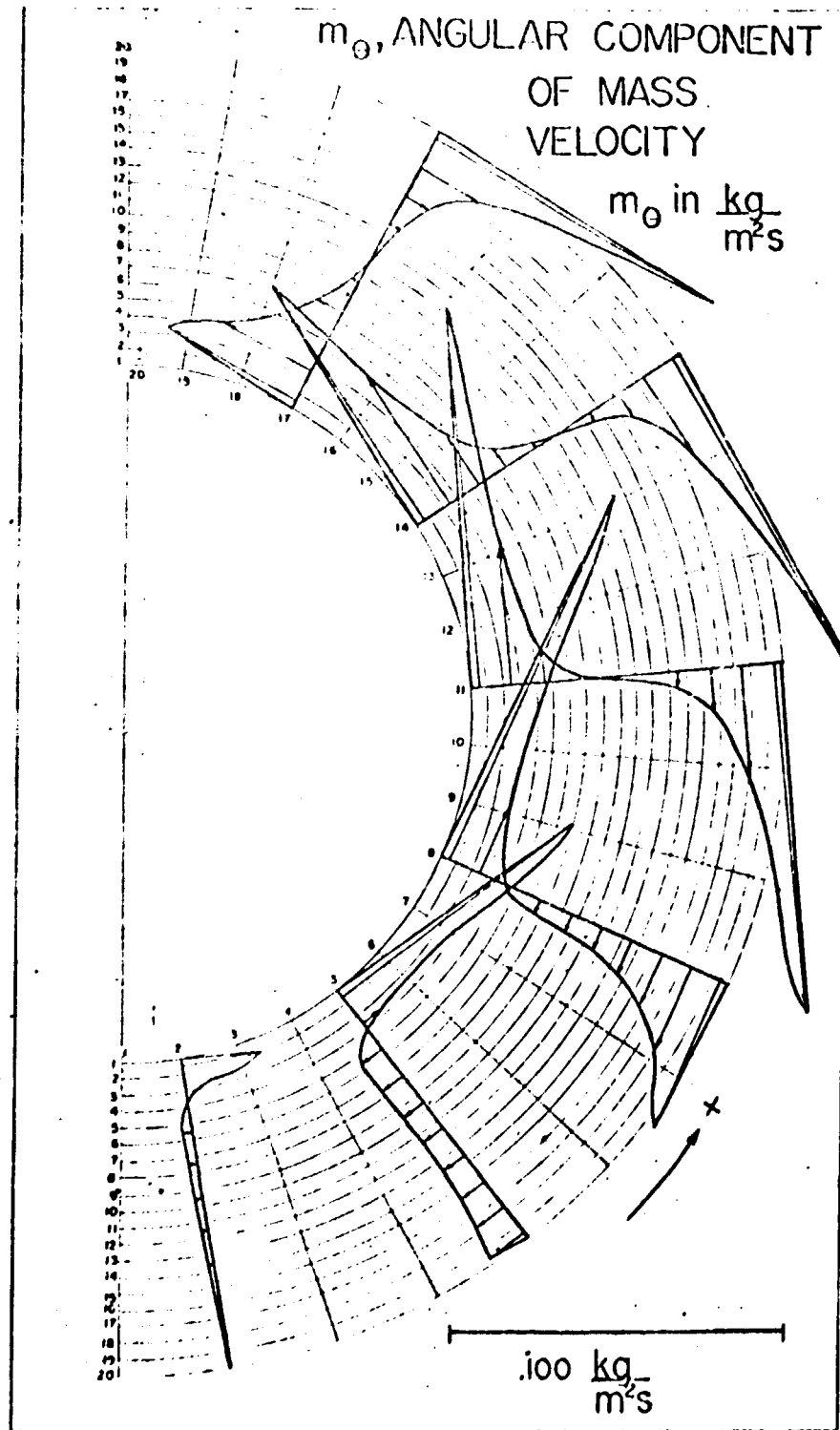


Figure 9-d – Horizontal Cylinder, Angular Component of Mass Velocity, $Re = 425$, $Nu = 4.83$

Effective Thermal Conductivity

The effective thermal conductivity λ^e is defined by

$$\lambda^e = \frac{Q \ln(r_2/r_1)}{(T_1 - T_2) \Pi} \quad (60)$$

where Q is the heat flow through the insulation. Q is calculated at each radial position r from

$$Q_r = \sum_{j=1}^{n-1} v^j \Delta \theta (\rho^{i,j} v_r^i) = \lambda^{i,j} \frac{\partial T}{\partial r} \quad (61)$$

where i and j are the node indices in the radial and angular directions respectively. At the walls ($i=1$ and $i=n$), the temperature gradient $\frac{\partial T}{\partial r}$ is evaluated from a second degree approximation of the temperature profile near the wall. The convective terms v_r are zero in this case. At other values of r , $\frac{\partial T}{\partial r}$ is approximated by central differences. The computed Q_r 's differed by about $\pm 5\%$, and the heat flow was taken as the arithmetic average of the hot and cold wall values. When $\frac{\partial T}{\partial r}$ was evaluated with a linear approximation of the temperature profile near the walls, lower heat flows (up to 30%) were obtained. This is explained by the temperature gradients at the walls which are not well approximated by a linear temperature profile.

The overall Nusselt number Nu is defined as

$$Nu = \frac{\lambda^e}{k_f}$$

where λ is the thermal conductivity of the insulation (last = first).

The overall Nusselt number correlated with an equation of the form $Nu = c \left(\frac{Ra}{A} \right)^n$. The result was (Figure 10)

$$\begin{aligned} Nu &= .48 \left(\frac{Ra}{A} \right)^{.5} & Ra > 1 \\ Nu &= 1 & Ra < 1 \end{aligned} \quad (62)$$

Where the Rayleigh number Ra is defined by

$$Ra = \frac{g(T_H - T_C) d K \beta}{\nu_r \alpha_r}$$

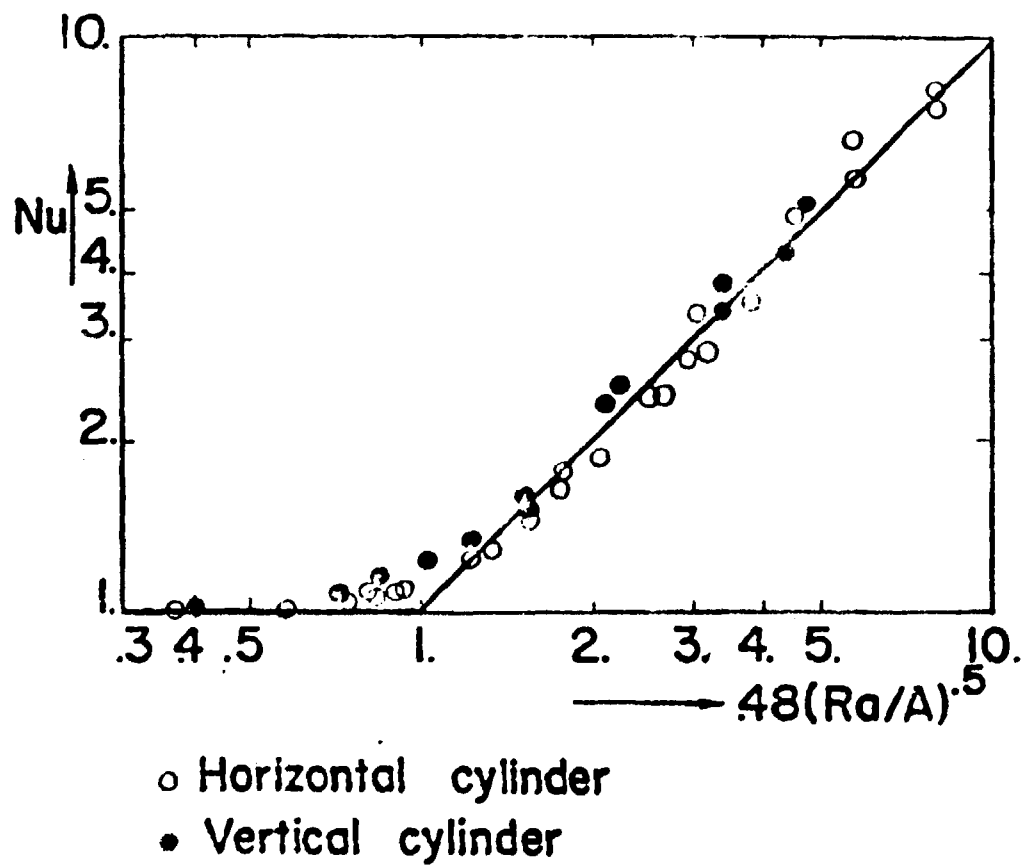


Figure 10 - Correlation of Overall Nusselt Number for Horizontal Cylinder and Vertical Cylinder Without Permeation Flow.

Where $K = K_r$ for cylindrical and $K = K_v$ for rectangular geometries.

The computed values of Nu showed a spread of $\pm 10\%$ around the curve $Nu = .48 \left(\frac{Ra}{A} \right)^{.5}$. The lower values are from runs with 13 and 25 bars and higher ones from runs at 80 and 100 bars.

The variation of the local Nusselt number Nu_j with θ at the cold wall is shown in Figure 11.

$$\frac{Nu_j}{Nu} \text{ is obtained from}$$

$$\frac{Nu_j}{Nu} = \frac{\lambda^{m,j} \frac{\partial T^{m,j}}{\partial r}}{Q / (n-1)} \quad (63)$$

4.2 - Vertical Cylinder

The range of input parameters was similar to that for the horizontal cylinder. In addition the following, permeation flow parameters have been used:

Cover plate permeability K_c :	1.E-8 - 1.E-12 m ²
Pressure gradient in the gas stream $\frac{\partial p_o}{\partial x}$:	- 1.5 to + 1.5 $\frac{\text{bar}}{\text{m}}$
Cover plate thickness:	.002 m
Prandtl Number:	.40 - .90

Closed Hot Wall

The computed velocity and temperature distributions (Figures 12-a and 12-b) found to be similar to those discussed earlier for the horizontal cylinder. λ^e was calculated from

$$\lambda^e = \frac{Q \ln(r_2/r_1)}{2(T_1 - T_2) L H} \quad (64)$$

Where L is the height of the cylinder. Q , the heat flux through the isolation has been again computed at each radial position r^i from

$$Q_i = \sum_{j=1}^{n-1} 2 \Pi \Delta x r^i \left(\rho \cdot v_r^{i,j} - \lambda^{i,j} \frac{\partial T^{i,j}}{\partial r} \right) \quad (65)$$

Where i and j are the node indices in the radial and vertical directions respectively. $\frac{\partial T}{\partial r}$ evaluation is similar to that for the horizontal cylinder.

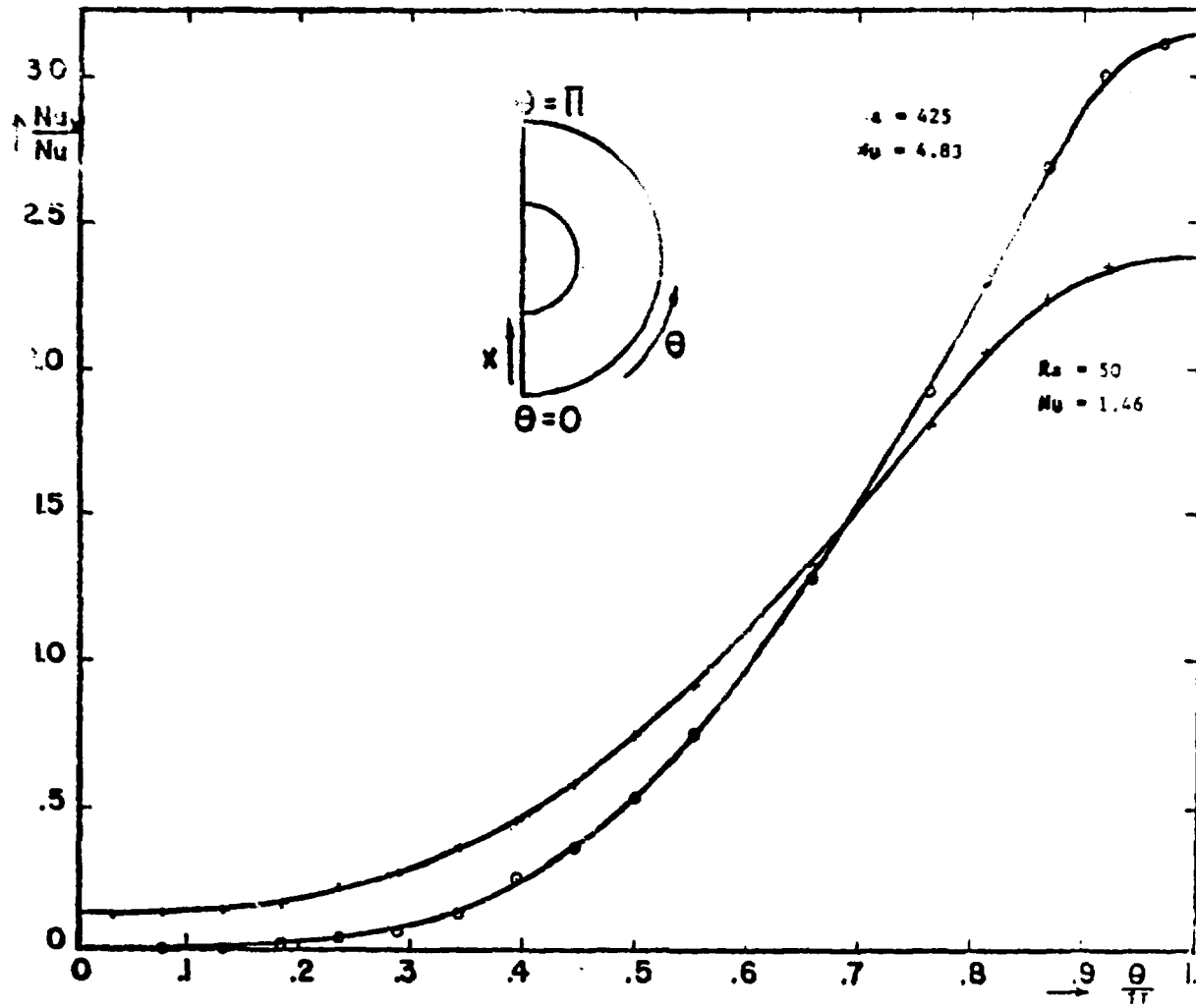


Figure 11 - Horizontal Cylinder, Variation of the Local Nusselt Number Along the Cold Wall

STREAM FUNCTION

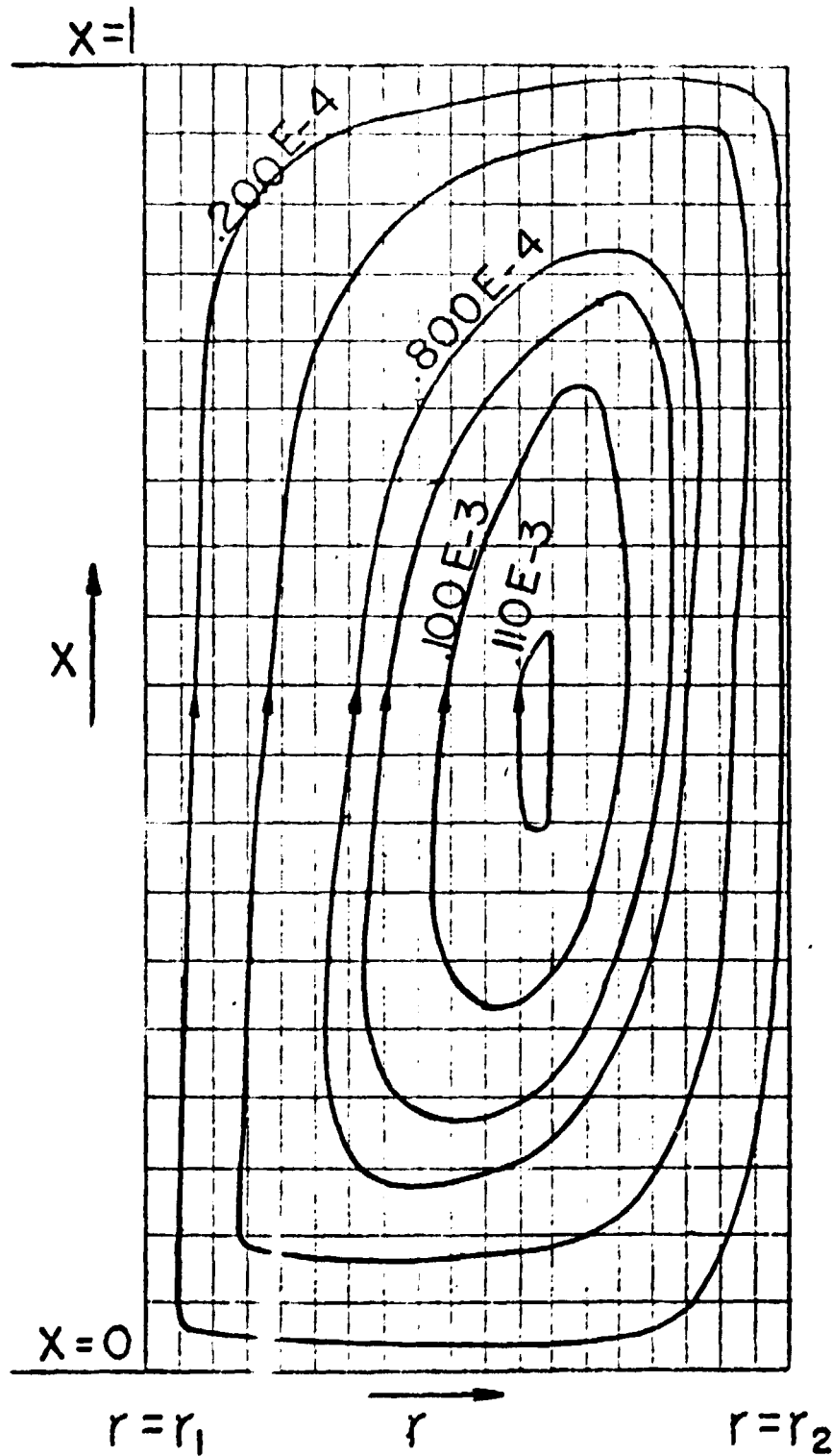


Figure 12-a - Vertical Cylinder, Stream Lines $Ra = 204$, $Nu = 2.32$. The scale in the radial direction is 5 times the scale in the vertical direction.

DIMENSIONLESS TEMPERATURES

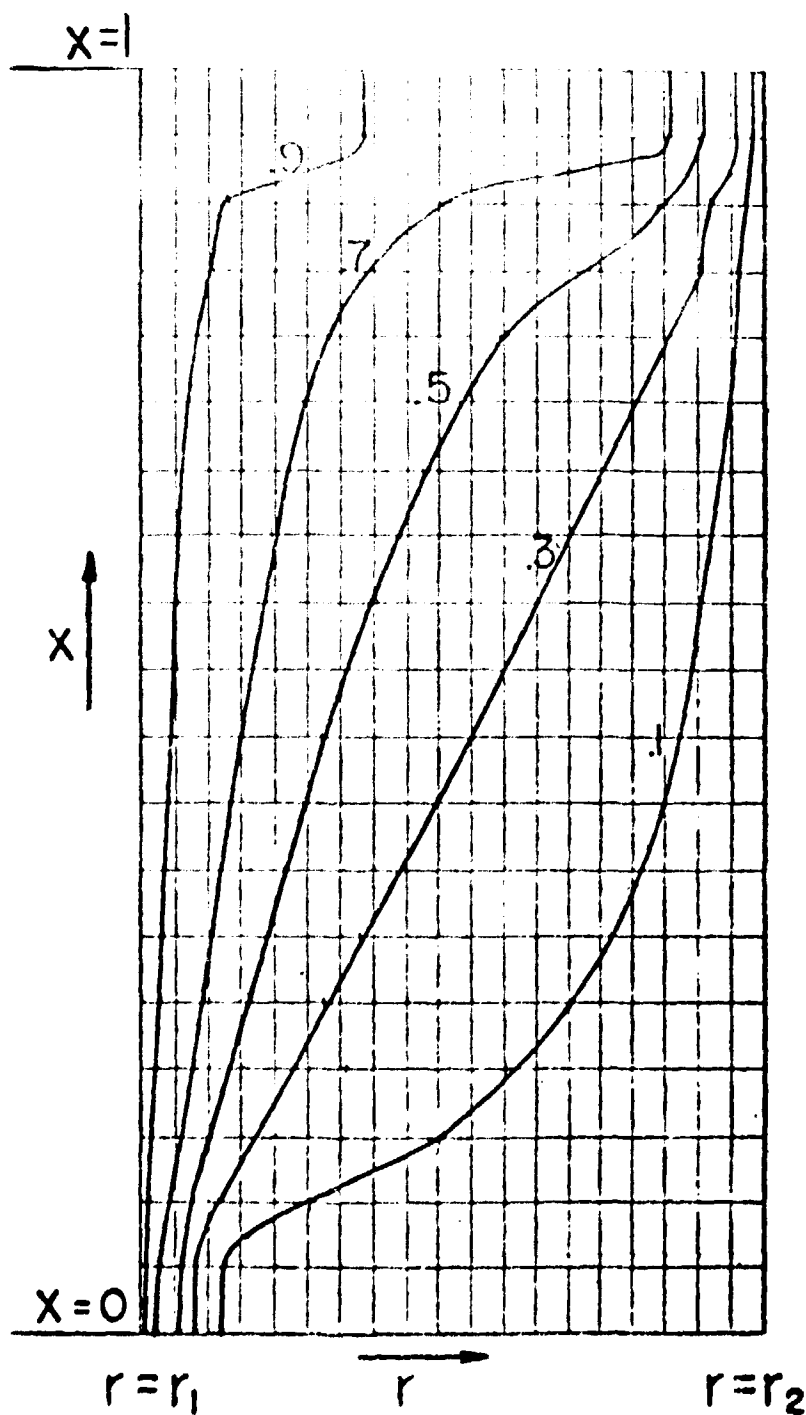


Figure 12-b - Vertical Cylinder, Dimensionless Temperatures $Rs = 204$, $Nu = 2.32$. The scale in the radial direction is 5 times the scale in the vertical direction.

The correlation of the overall Nusselt number with the parameter $\frac{Ra}{A}$ (Figure 10) results in the same equation (again within $\pm 10\%$) as for the horizontal cylinder which is

$$Nu = .48 \left(\frac{Ra}{A} \right)^{.5} \quad (66)$$

Figure 13 shows the variation of the cold face local Nusselt number Nu_j in the vertical direction. The ratio Nu_j/Nu was again obtained from Equation (63).

Permeable Hot Wall

When the hot wall is permeable, the velocity and temperature distribution in the insulation depends on the Rayleigh and Reynolds numbers. The Reynolds number is taken as positive when the hot gas in the inner pipe flows downward and as negative when it flows upward. The Reynolds number is defined by

$$Re = \frac{K_x (\partial p_o / \partial x - \rho_o g_x) d}{\mu_r \nu_o} \quad (67)$$

Where the subscript o refers to free stream and r to average insulation temperatures. The pressure gradient $\partial p_o / \partial x$ can be evaluated from friction factor correlations for flow inside pipes.

Figures 14-a and 14-b show the stream lines and dimensionless temperatures for $Ra = 51$ and $Re = 0$. The Nusselt number is 2.09. The outside gas enters the insulation layer in the upper part, flows downward along the cold wall leaves in the lower part, following a circulation pattern close to that observed in the case of a closed hot wall. Comparison of Figures 12-b and 14-b shows that temperature distributions are also similar. When $Ra = 5.1$ and $Re = 161$, flow and temperature fields characteristic of forced convection are seen (Figures 15-a and 15-b). They are different from what is encountered for natural circulation, but the flow is still downward along the cold wall. Figures 16-a and 16-b show a case where natural and forced convection ($Ra = 51$ and $Re = 29$) patterns are both present.

When Re is negative (upward gas flow in the pipe), forced and natural convection act in opposite directions. For $Ra = 5.1$ and $Re = -132$ (Figures 17-a and 17-b), the gas enters the insulation in the lower part, flows upward along the cold wall, and leaves in the upper part. The difference in the temperature profiles corresponding to positive and negative Reynolds numbers (with negligible natural circulation effects) can be observed comparing Figures 17-b and 15-b. When $Ra = 51$ and $Re = 88$, Figure 18-a shows a forced convection pattern on the hot wall side and a natural circulation pattern on the cold wall side, acting in opposite direction; which results in the almost parallel and equidistant isotherms, which $Nu = 1.17$ (Figure 18-b). For $Ra = 126$ and $Re = -128$ (Figures 19-a and 19-b), the beginning of forced convection permeation flows can be observed near the hot wall, while the rest of the flow field is dominated by natural circulation. Finally, for $Re = 51$ and $Re = -28$, the forced convection is not observable any more (Figures 20-a and 20-b), but it reduces the amount of natural convection which results in a Nu of 1.58 ($Nu = 2.09$ for $Ra = 50$ and $Ra = 0$).

The variation of local Nusselt number on the cold wall is shown on Figure 21. For $Re \geq 0$, Nu is maximum on top and decreases towards the bottom. For $Re < 0$, and $Ra \ll Re$, Nu is maximum on the bottom and decreases towards the top. For Re negative and Ra of comparable magnitude, variation of the local Nusselt number show a mixed forced-natural convection pattern.

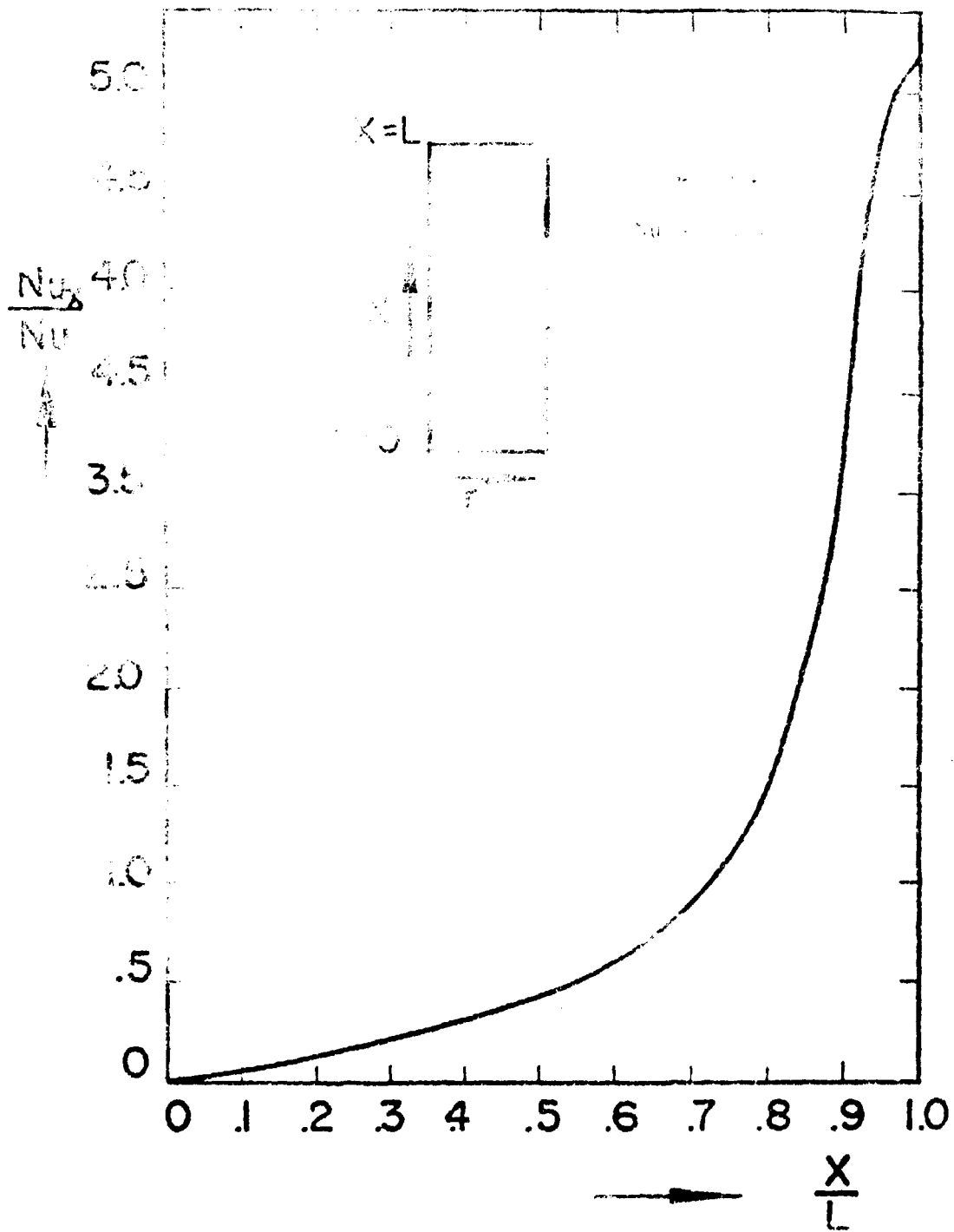


Figure 13 - Vertical Cylinder Without Permeation Flow, Variation of the Local Nusselt Number Along the Cold Wall

STREAM FUNCTION

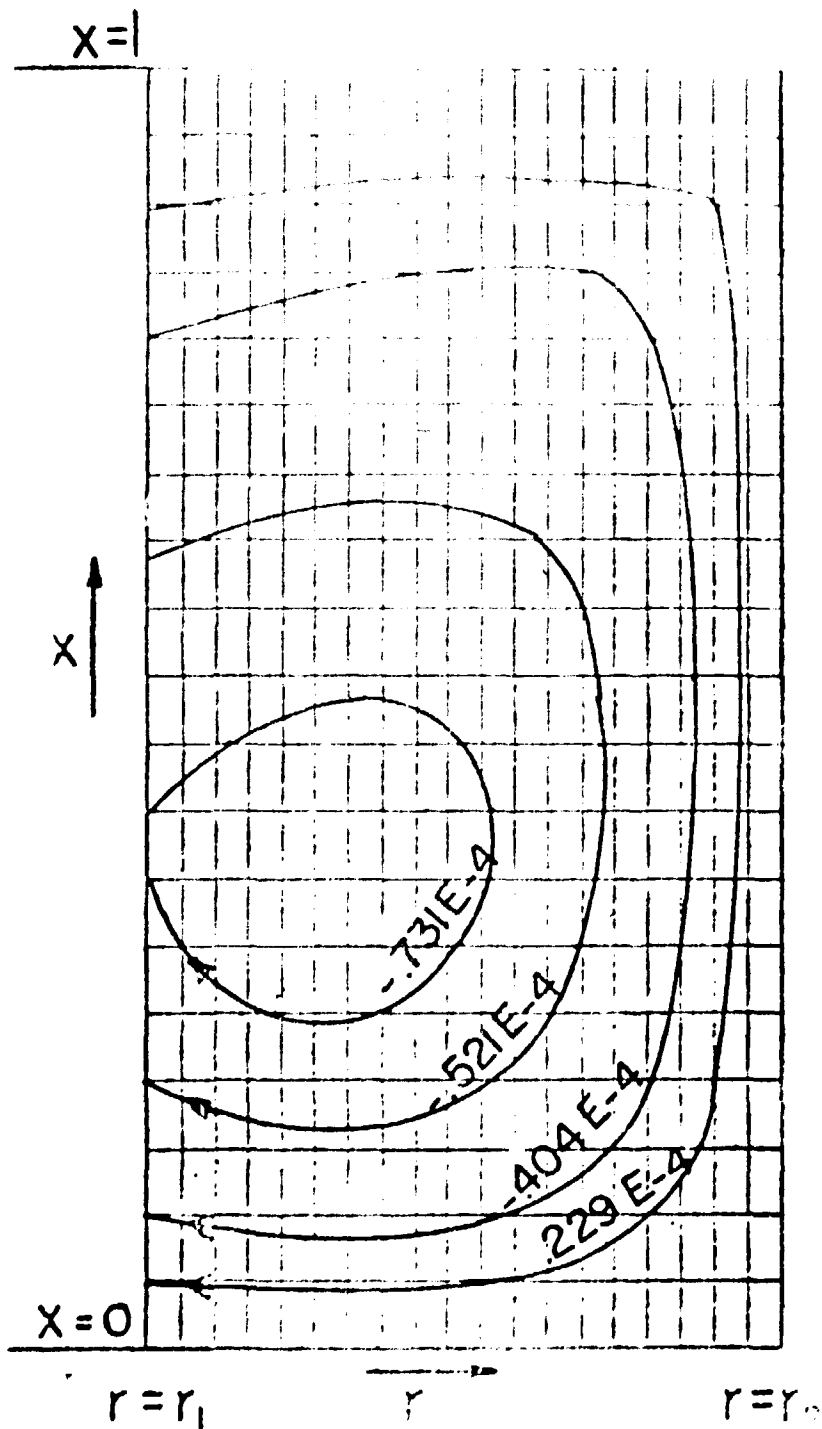


Figure 14-e — Vertical Cylinder. Stream Lines $Re = 51$, $Re = 0$, $Nu = 2.09$. The scale in the radial direction is 6 times the scale in the vertical direction.

DIMENSIONLESS TEMPERATURES

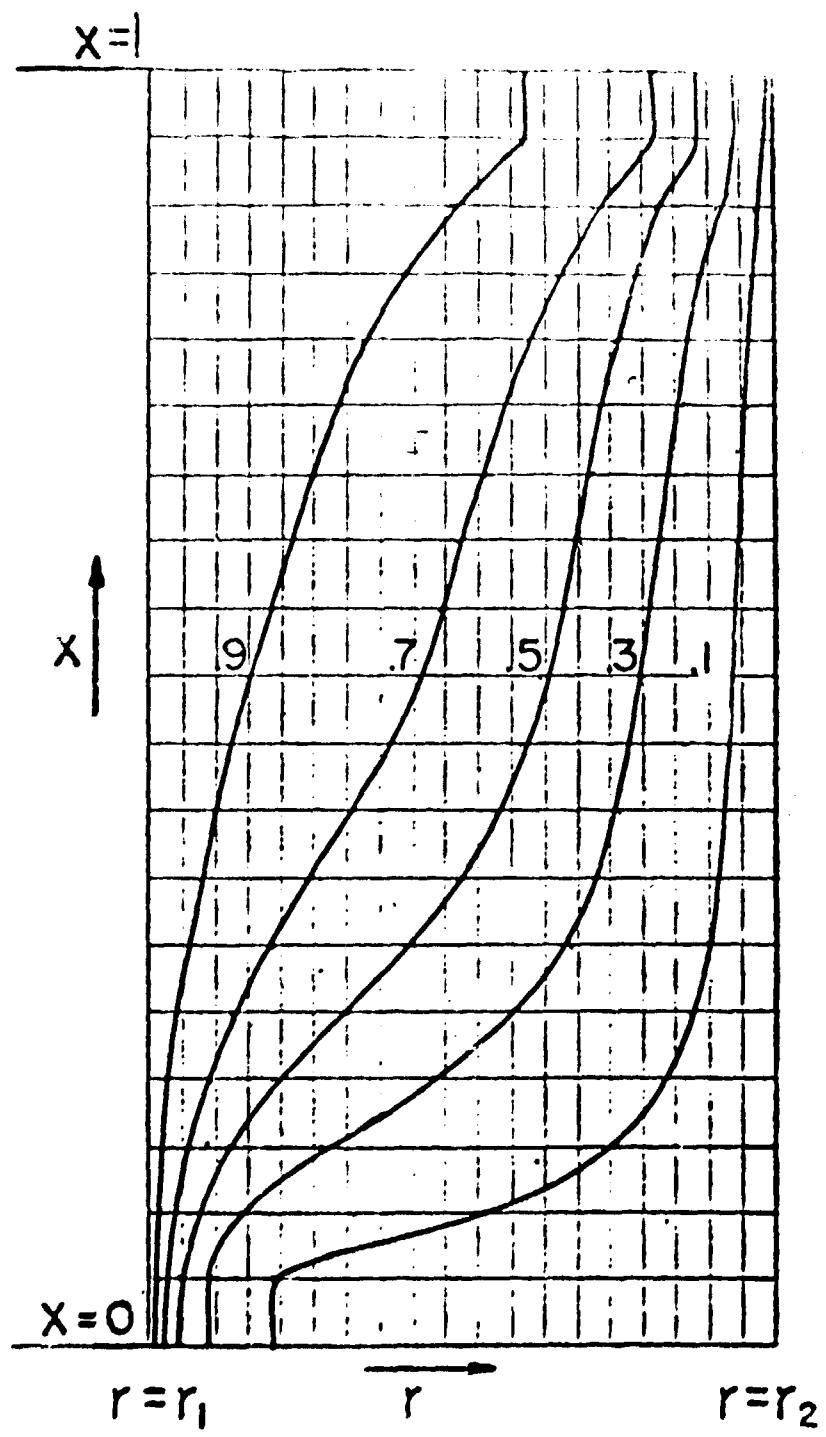


Figure 14-b - Vertical Cylinder, Dimensionless Temperatures $Ra = 51$, $Re = 0$, $Nu = 2.00$. The scale in the radial direction is 5 times the scale in the vertical direction.

STREAM FUNCTION

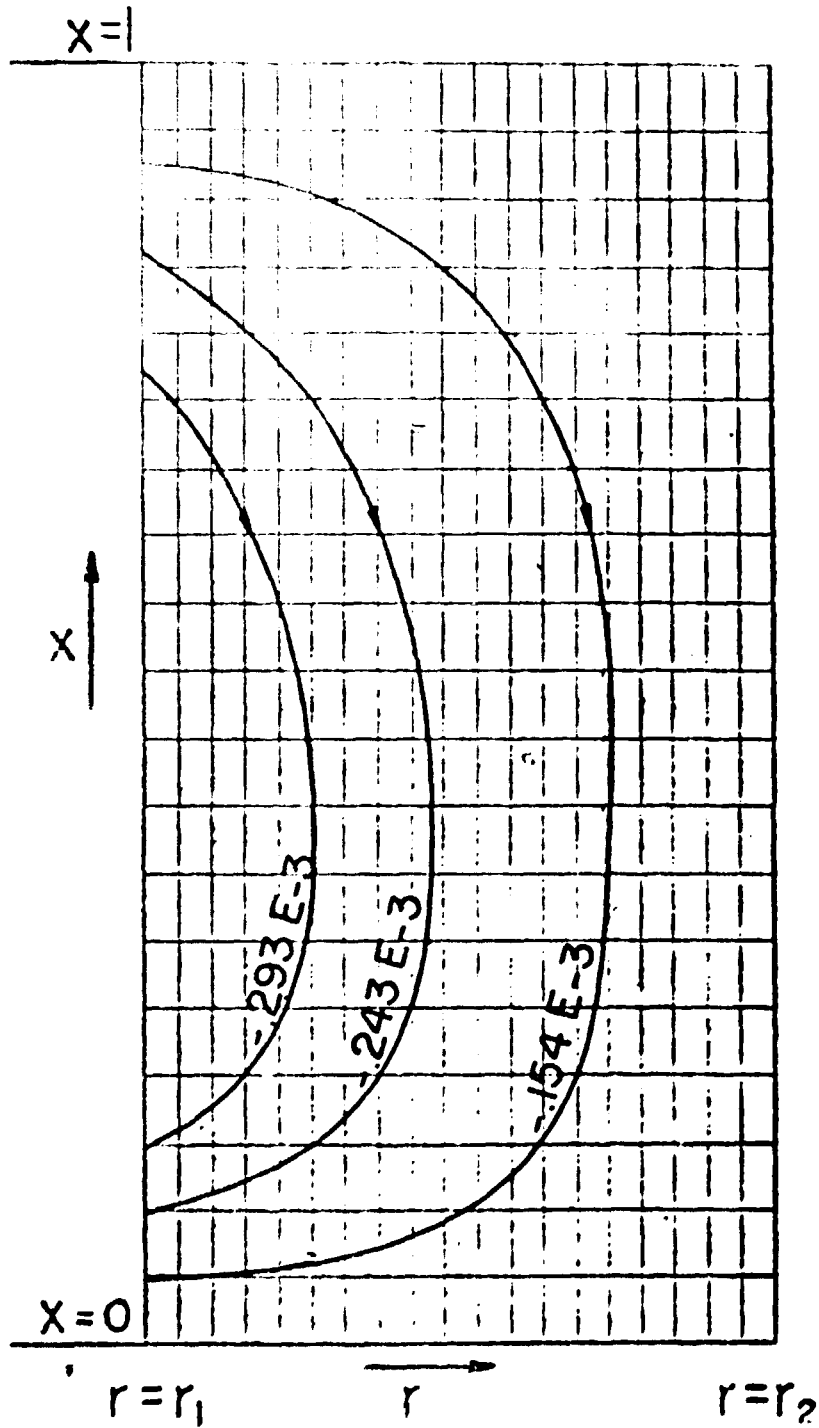


Figure 18-a - Vertical Cylinder, Stream Lines $Ra = 5.1$, $Re = 101$, $Nu = 4.67$. The scale in the radial direction is 5 times the scale in the vertical direction.

DIMENSIONLESS TEMPERATURES

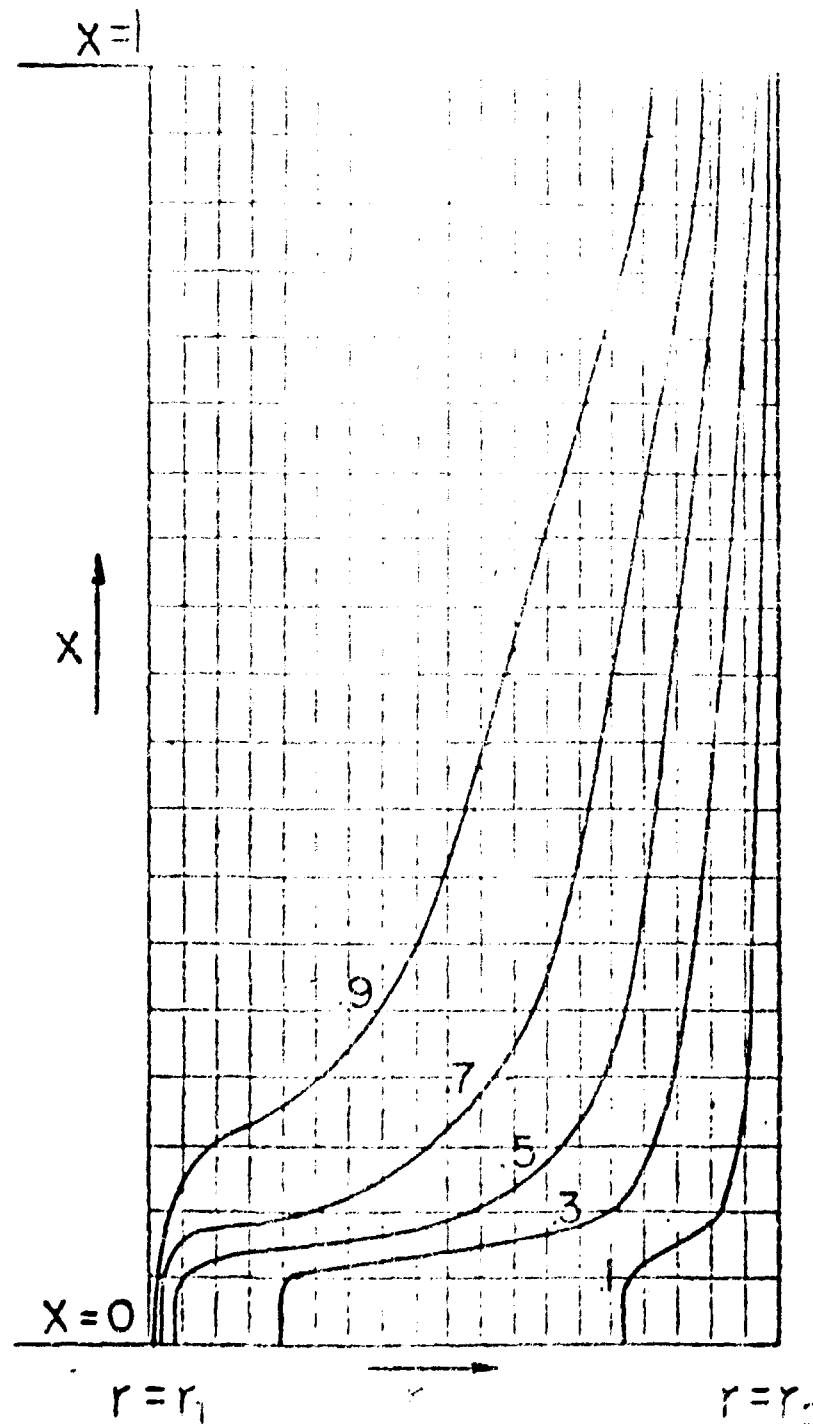


Figure 15-b - Vertical Cylinder, Dimensionless Temperatures $Ra = 5.1$, $Re = 161$, $Nu = 4.57$. The scale in the radial direction is 5 times the scale in the vertical direction.

STREAM FUNCTION

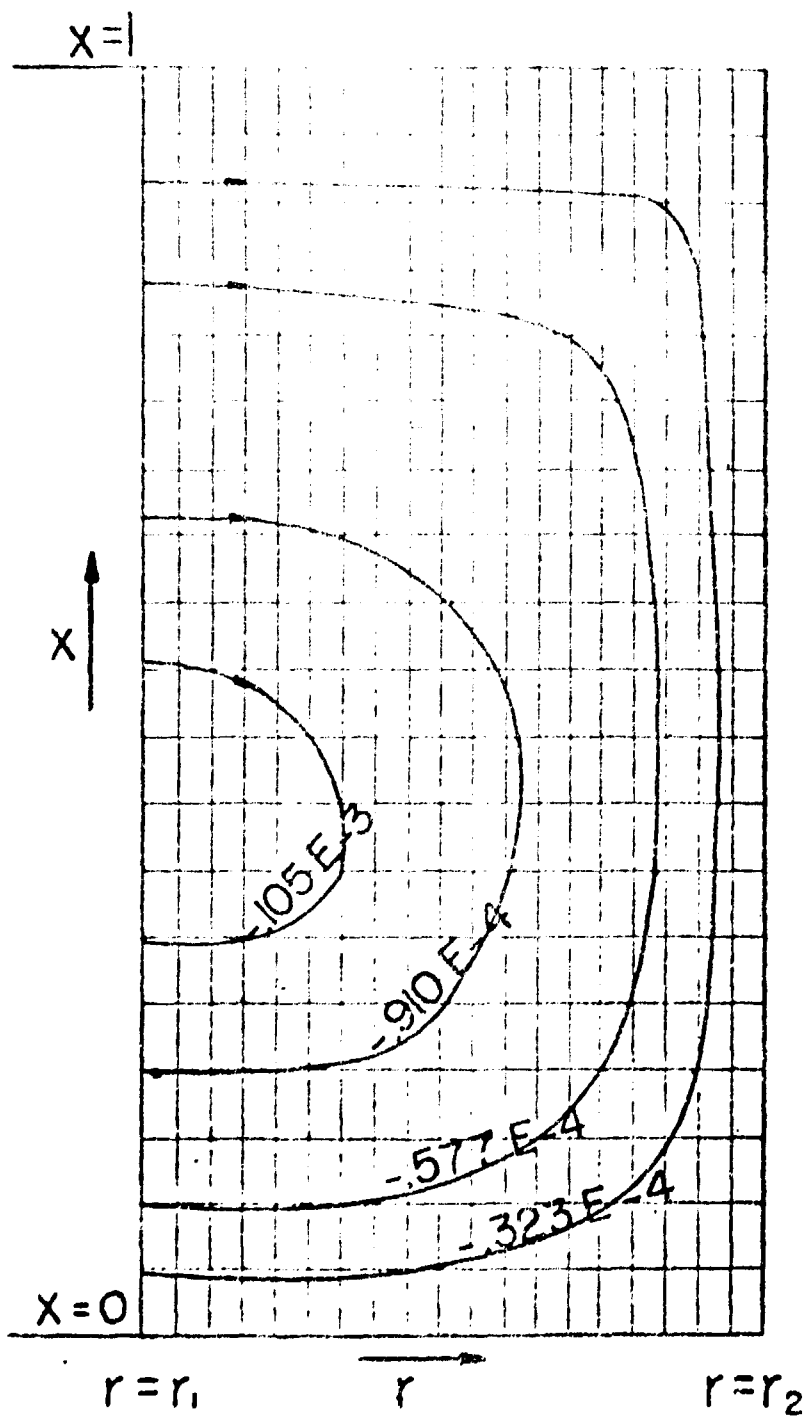


Figure 16-a - Vertical Cylinder, Stream Lines $Re = 61$, $Re = 29$, $Nu = 2.68$. The scale in the radial direction is 5 times the scale in the vertical direction.

DIMENSIONLESS TEMPERATURES

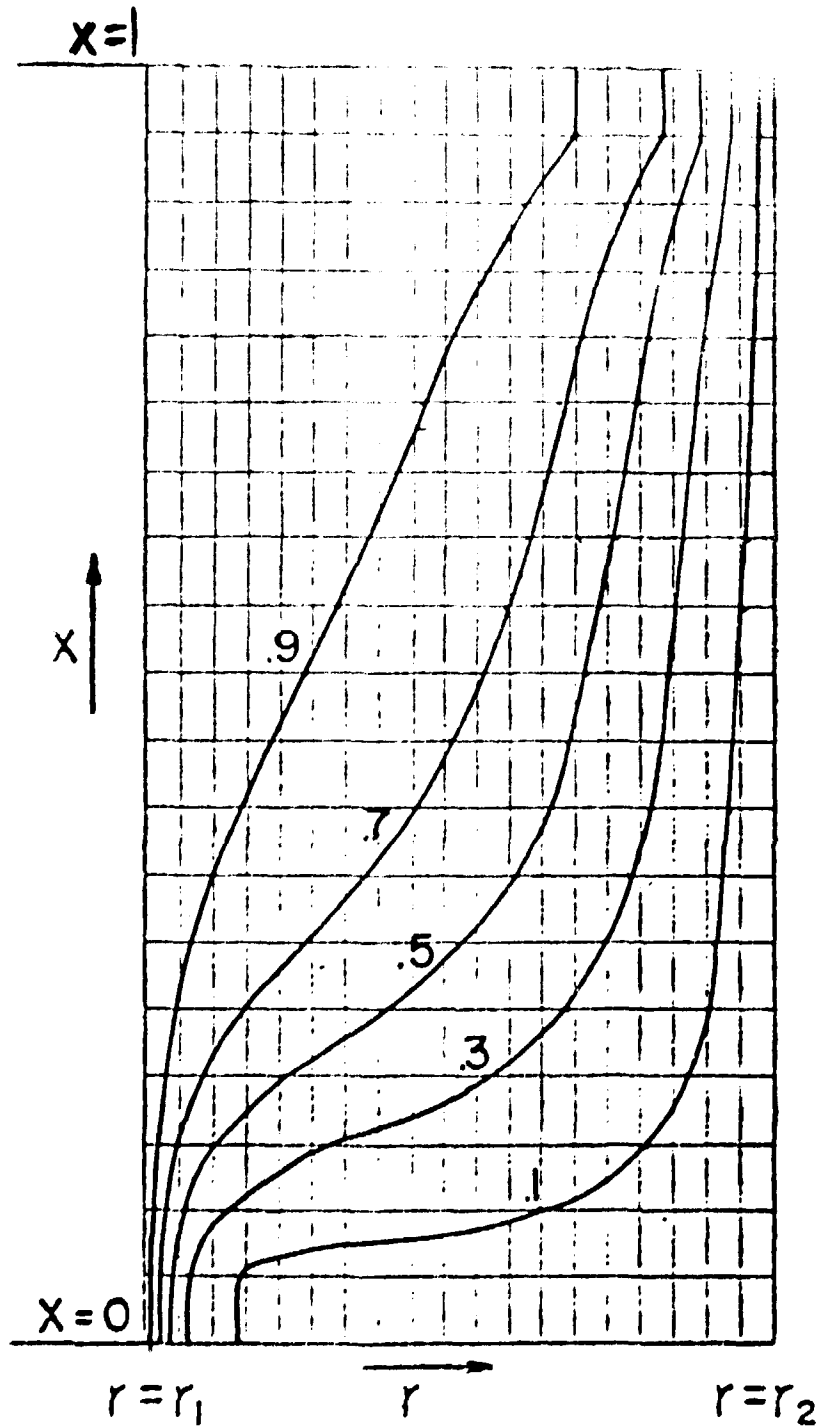


Figure 16-b - Vertical Cylinder, Dimensionless Temperatures $Ra = 51$, $Re = 29$, $Nu = 2.58$. The scale in the radial direction is 5 times the scale in the vertical direction

STREAM FUNCTION

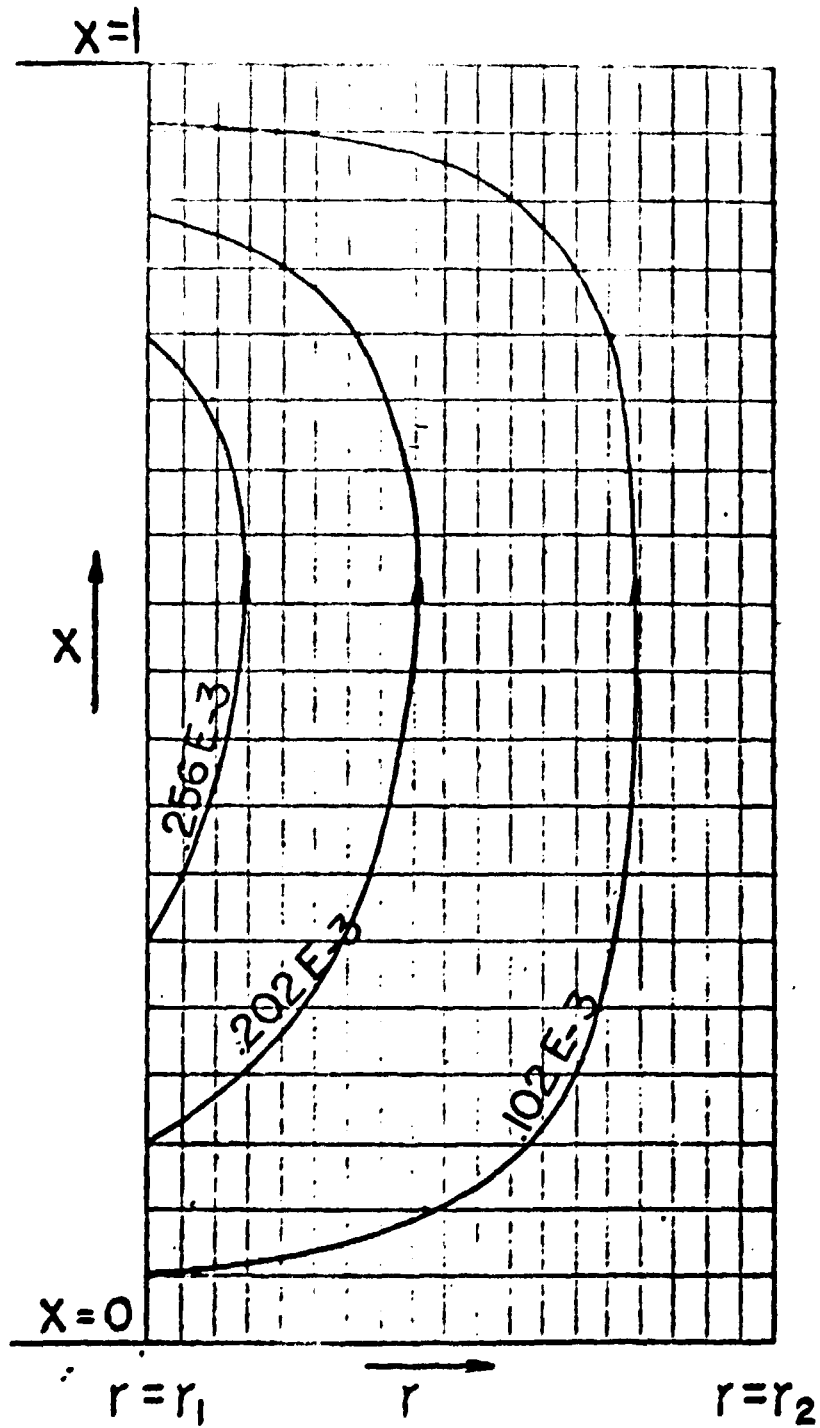


Figure 17a - Vertical Cylinder, Stream Lines $Re = 5.1$, $Re = -132$, $Nu = 3.89$. The scale in the radial direction is 5 times the scale in the vertical direction.

DIMENSIONLESS TEMPERATURES

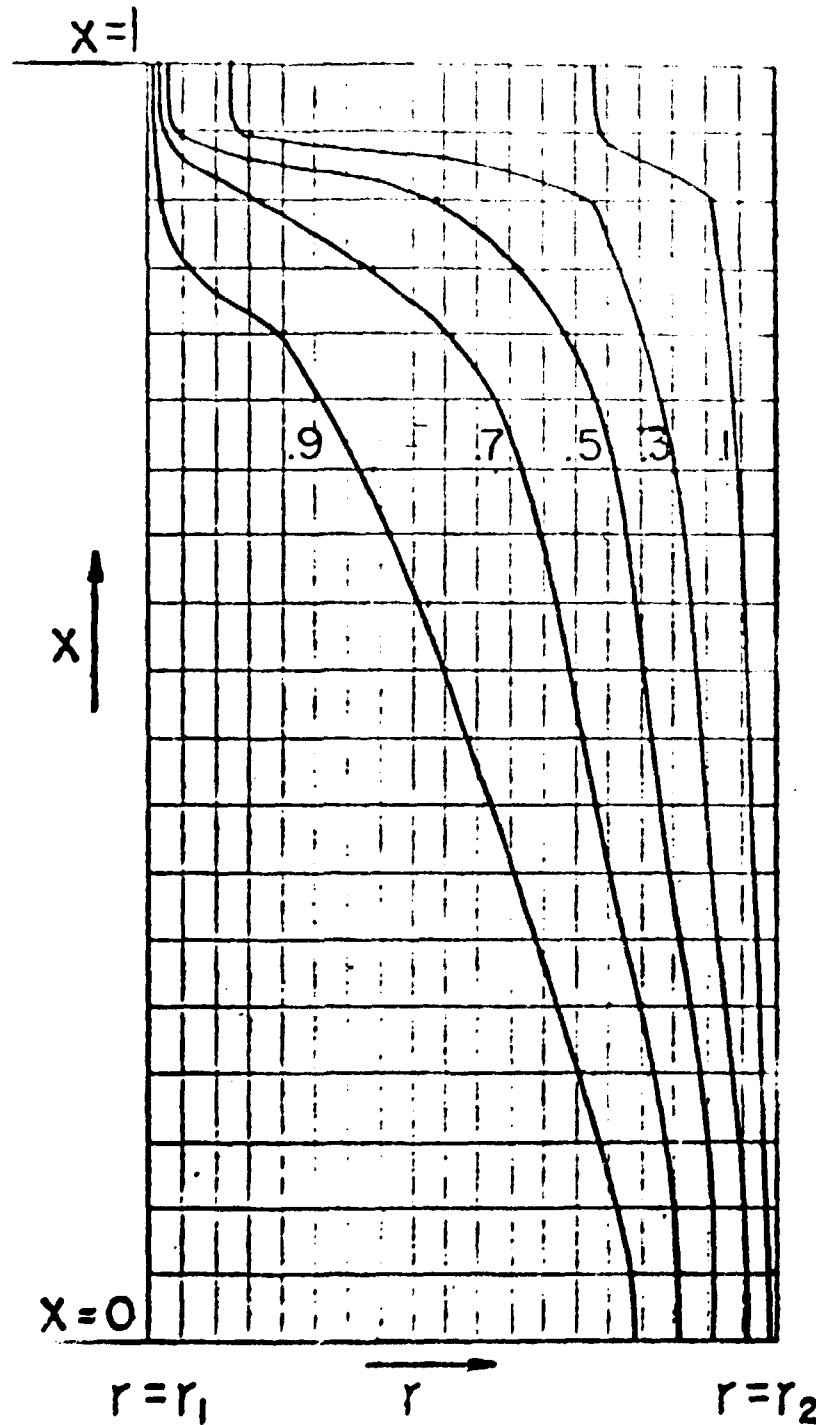


Figure 17-4 - Vertical Cylinder, Dimensionless Temperatures $Ra = 5.1$, $Re = -132$, $Nu = 3.88$. The scale in the radial direction is 5 times the scale in the vertical direction.

STREAM FUNCTION

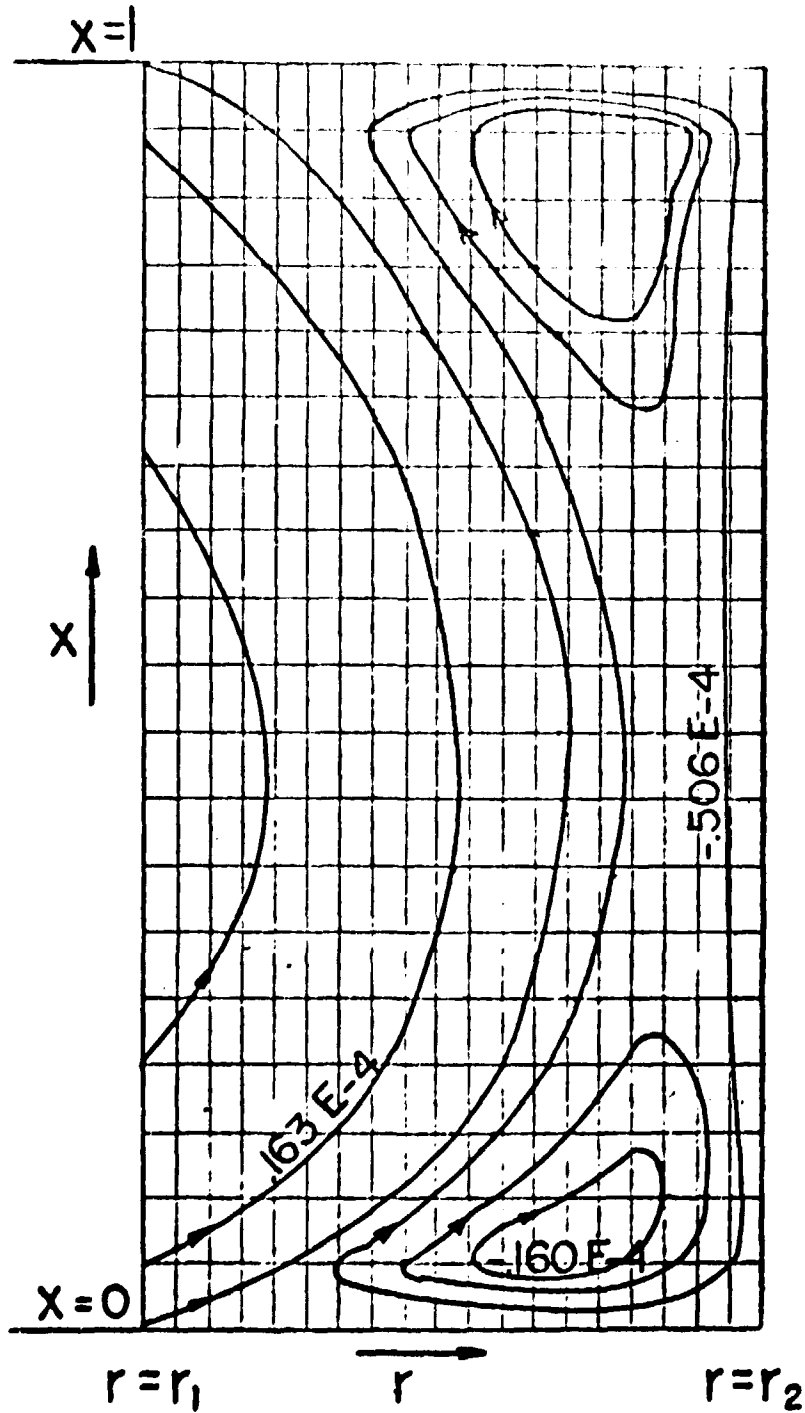


Figure 18-e - Vertical Cylinder, Stream Lines $Re = 51$, $Re = -85$, $Nu = 1,17$. The scale in the radial direction is 5 times the scale in the vertical direction.

DIMENSIONLESS TEMPERATURES

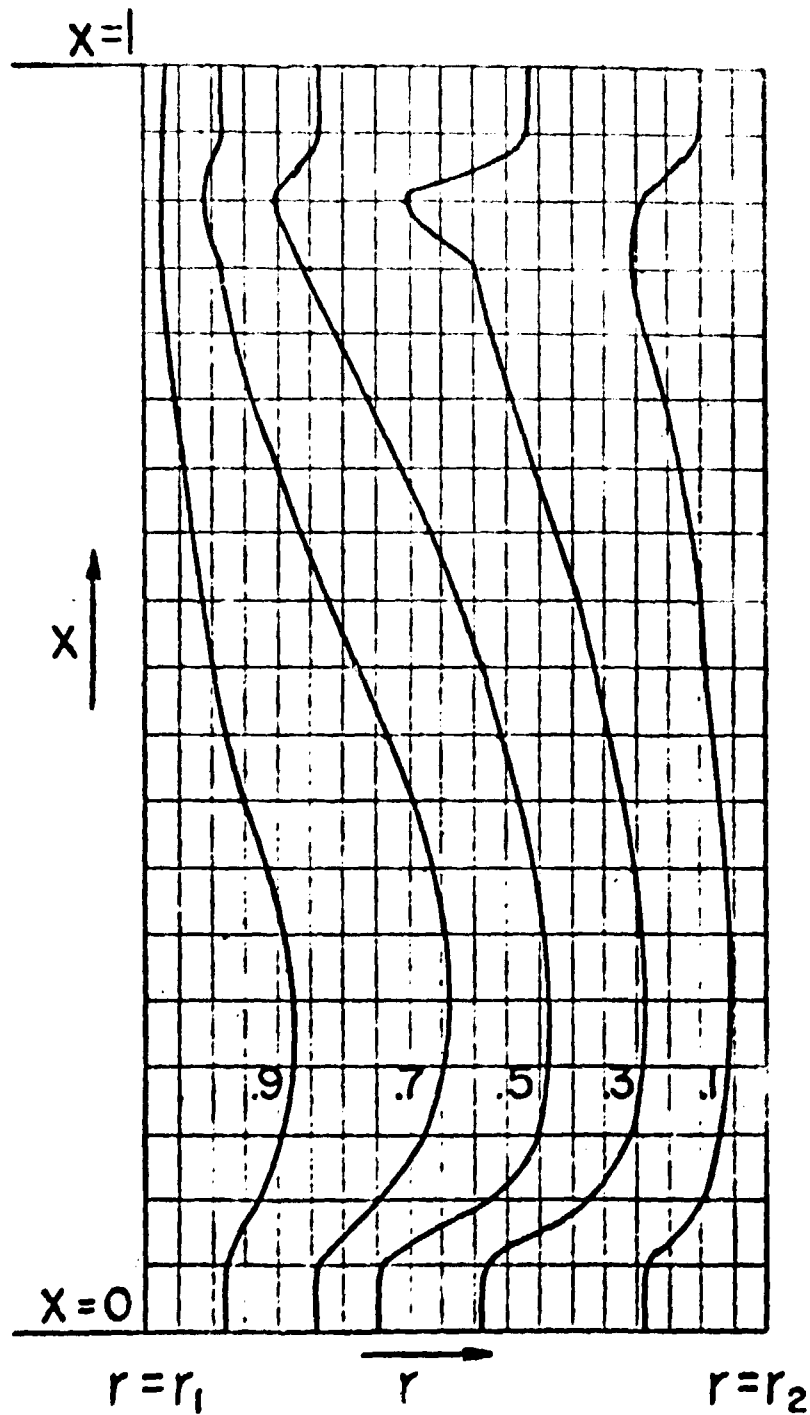


Figure 18-b - Vertical Cylinder, Dimensionless Temperatures $Ra = 51$, $Re = -85$, $Nu = 1.17$. The scale in the radial direction is 5 times the scale in the vertical direction.

STREAM FUNCTION

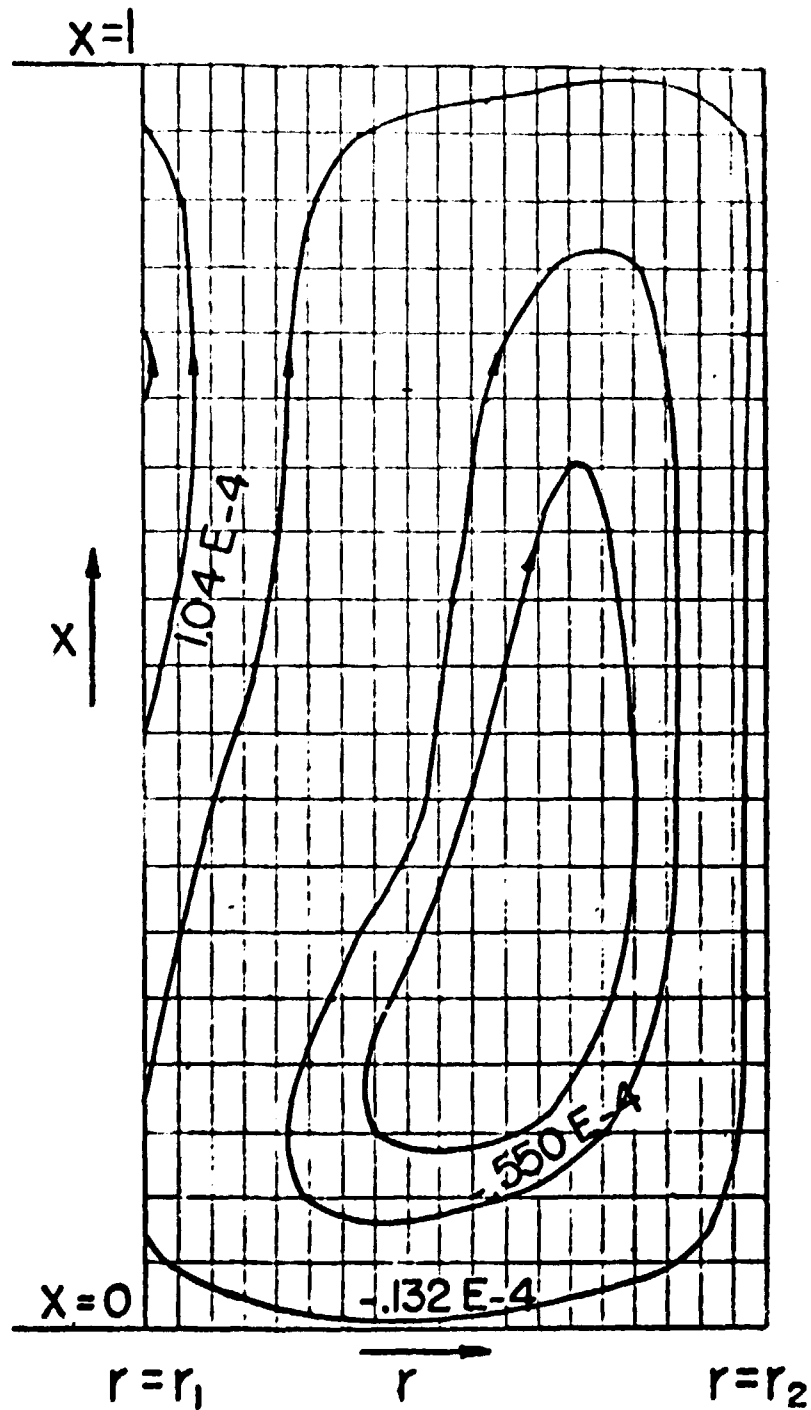


Figure 19-a - Vertical Cylinder, Stream Lines $Re = 126$, $Re = -126$, $Nu = 1.62$. The scale in the radial direction is 5 times the scale in the vertical direction.

DIMENSIONLESS TEMPERATURES

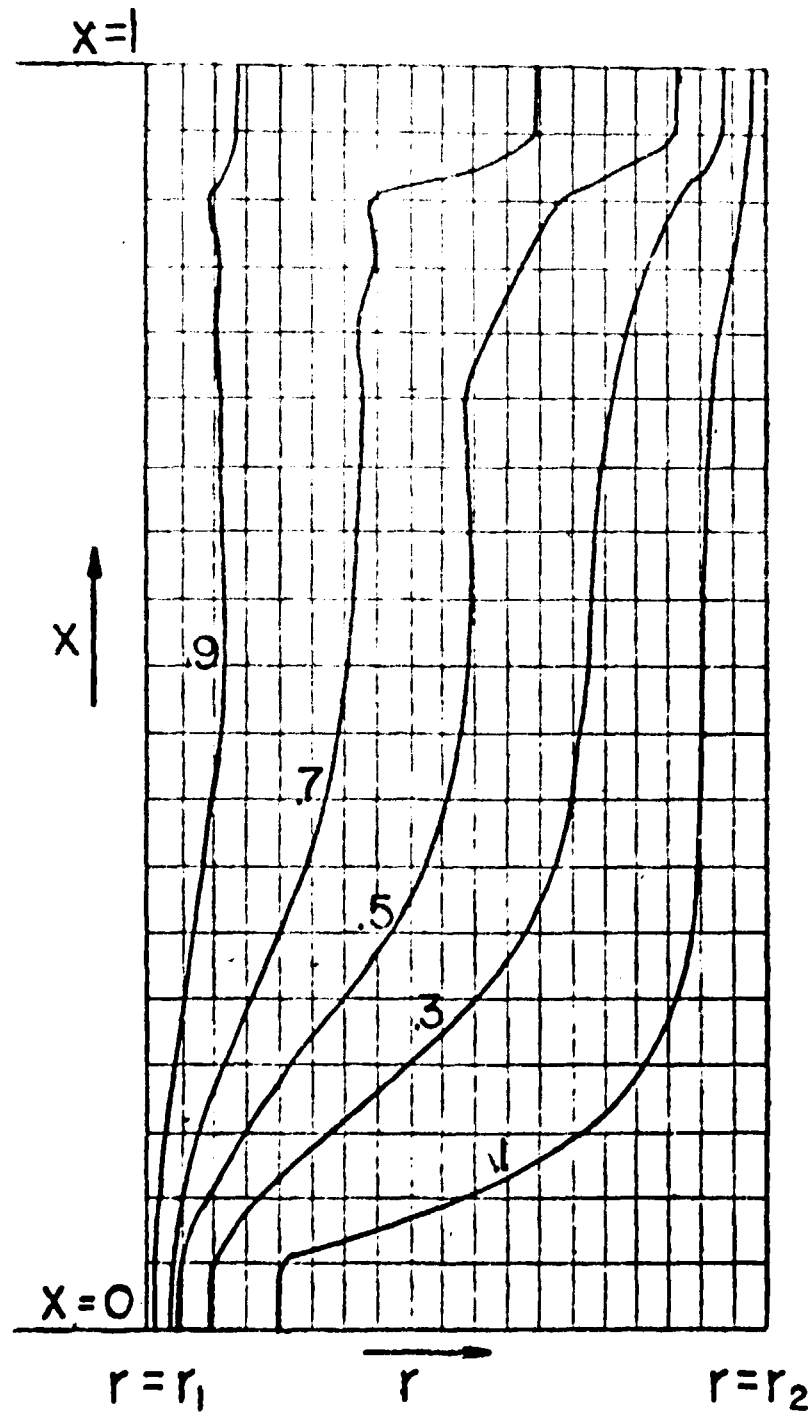


Figure 19-b - Vertical Cylinder, Dimensionless Temperatures $Ra = 120$, $Re = -128$, $Nu = 1.62$. The scale in the radial direction is 5 times the scale in the vertical direction.

STREAM FUNCTION

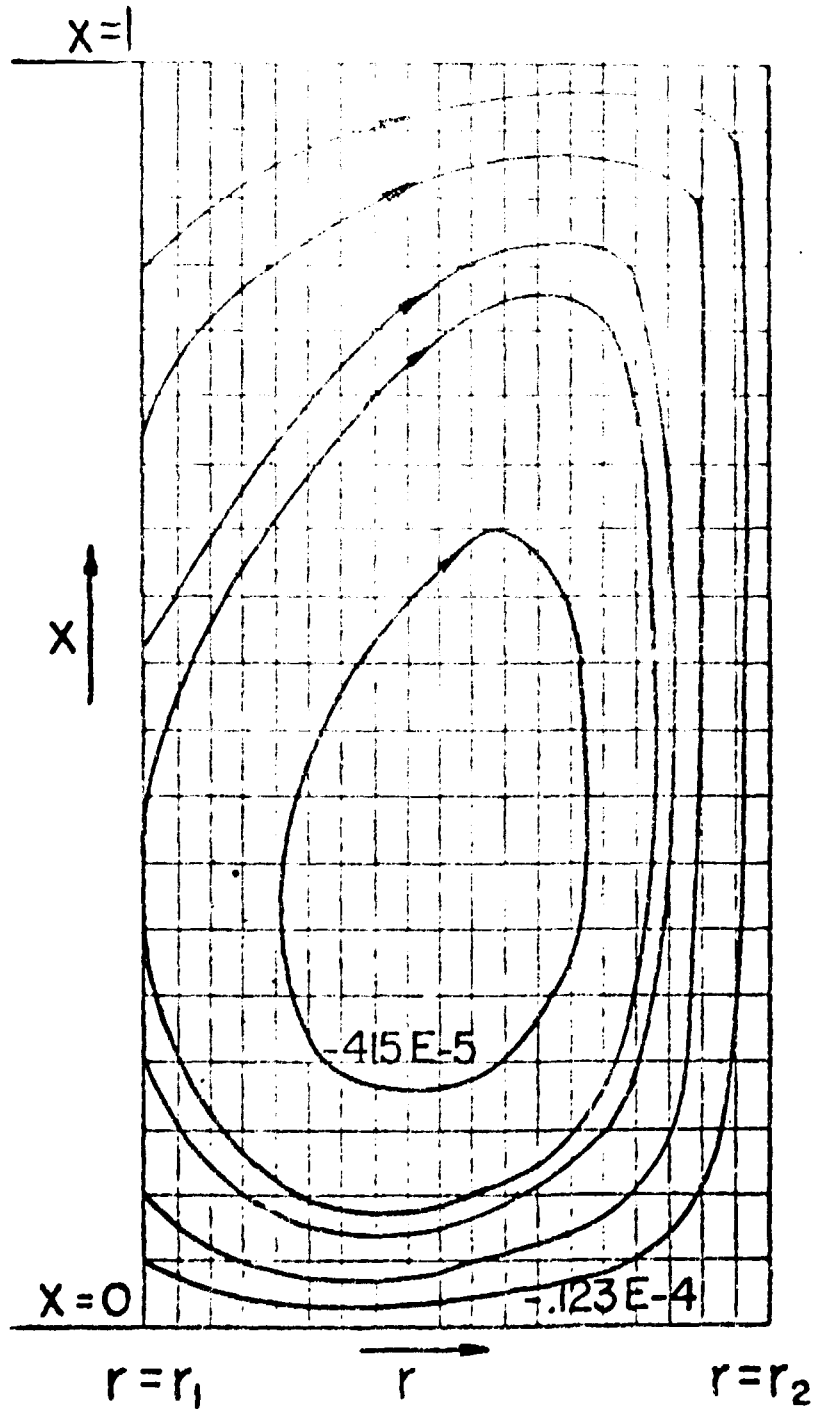


Figure 20-e - Vertical Cylinder, Stream Lines $Re = 51$, $Re = -28$, $Nu = 1.58$. The scale in the radial direction is 5 times the scale in the vertical direction.

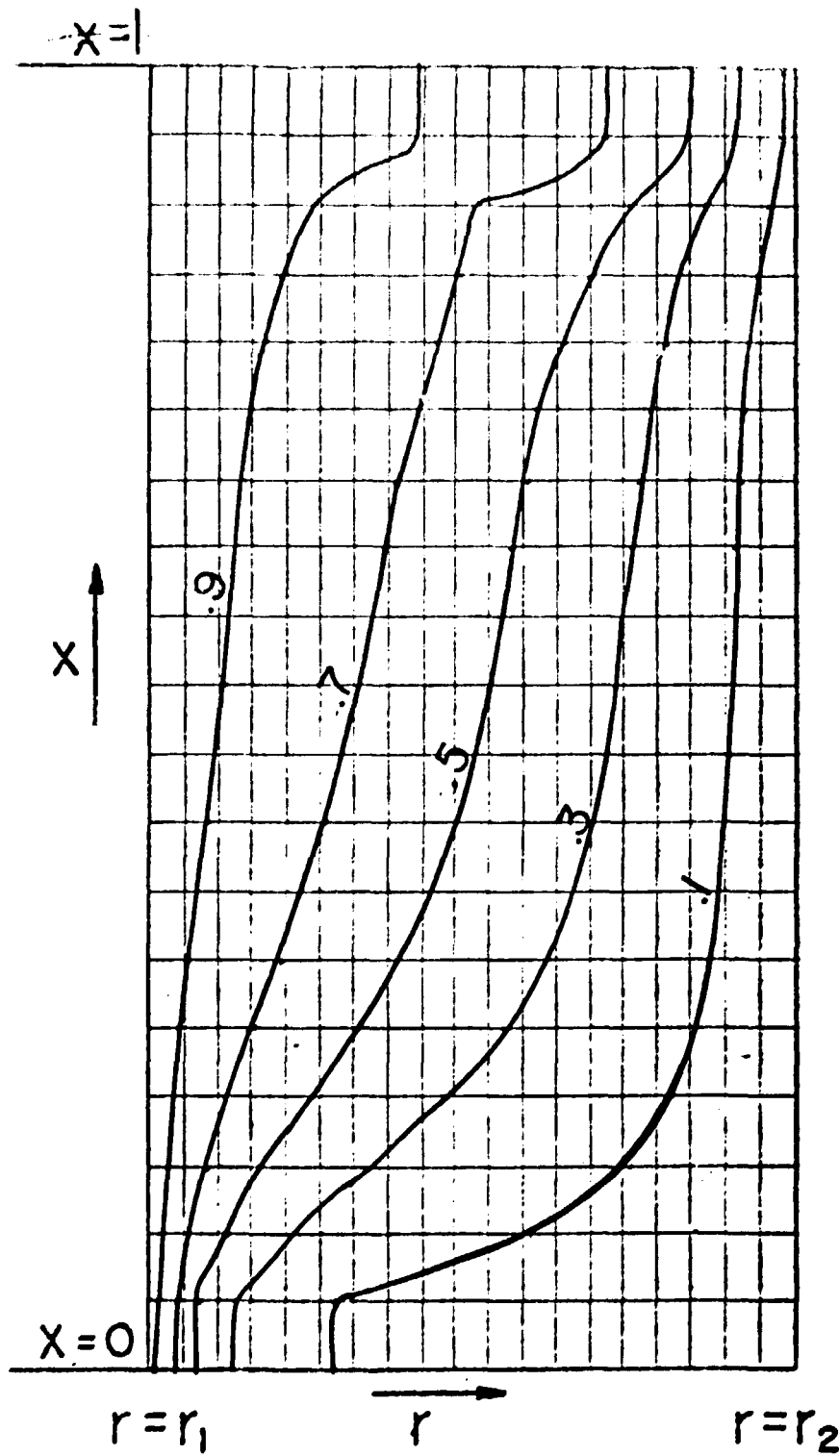


Figure 20-b - Vertical Cylinder, Dimensionless Temperatures $Ra = 51$, $Re = -28$, $Nu = 1.58$. The scale in the radial direction is 5 times the scale in the vertical direction.

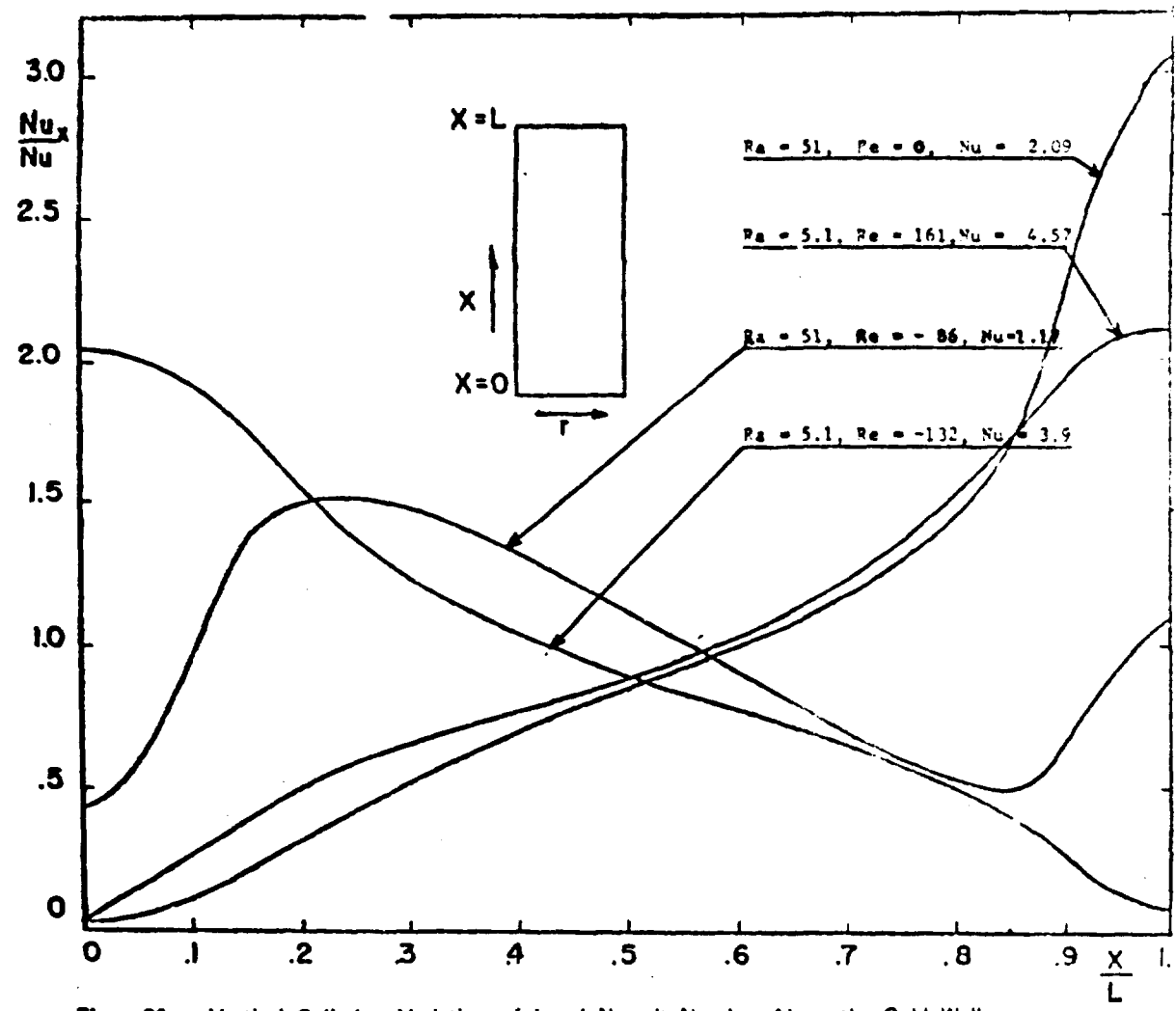


Figure 21 - Vertical Cylinder, Variation of Local Nusselt Number Along the Cold Wall

Correlation of Nu

For a rectangular cavity, Nu has been correlated by Jannot (6) by

$$Nu = .8 \left(\frac{Ra + 1.4 Pr Re}{A} \right)^{.55} \quad (68)$$

for $4 < .8 \left(\frac{Ra + 1.4 Pr Re}{A} \right)^{.55} < 10^3$, $A \geq 10$. Where Re can be positive or negative. To obtain Equation (68), the parameter $p = \frac{aK}{dK_c}$ has been set to 0 and Pr to 1. When $p = 0$, the resistance of the permeable wall to flow is much less than that of the insulation. Then the pressure gradient in the insulation will be equal to that outside the permeable hot wall. Thus the Reynolds number (Equation (67)) which uses the outside pressure gradient $\frac{\partial p_o}{\partial x}$ will give a good representation of permeation flow in the insulation. On the other hand, when the wall has a relatively high resistance to flow, only a fraction of the outside pressure gradient $\frac{\partial p_o}{\partial x}$ will exist in the insulation. To better approximate the pressure gradient $\frac{\partial p}{\partial x}$ inside the insulation, a combination of Re and p would be necessary. In the present study, Pr has been varied from .4 to .90, and p from 0 to 8. Figure 22 shows a plot of Nu vs $.8 \left(\frac{Ra + 1.4 Pr Re}{A} \right)^{.55}$. The spread of the points indicates the influence of parameters p and Pr. The upper values of Nu are from runs with p between 0 and .4, and lower values from runs with p between 4 and 8. On the other hand runs with constant p and varying Pr indicated that Nu increases about 10% as Pr increases from .4 to .9.

The parameter p can be combined with Re by defining a new Reynolds number, Re^* by

$$Re^* = \frac{Re}{1 + cp} \quad (69)$$

where c is a constant.

Thus, when the hot wall has no resistance to flow, $p = 0$, and $Re^* = Re$. The shielding effect of the high wall resistance would be taken into account by Re^* since Re^* would be smaller as p became larger. Also the influence of the Prandtl number could be better represented by raising it to an exponent. A more general correlation for Nu thus would be of the form

$$Nu = C_1 \left(\frac{Ra + m Pr^n Ra^*}{A} \right)^{.55} \quad (70)$$

where c_1 , m, and n are constants.

Due to time constraints, numerical values for the constants have not been evaluated. Instead, the Nusselt number have been correlated with

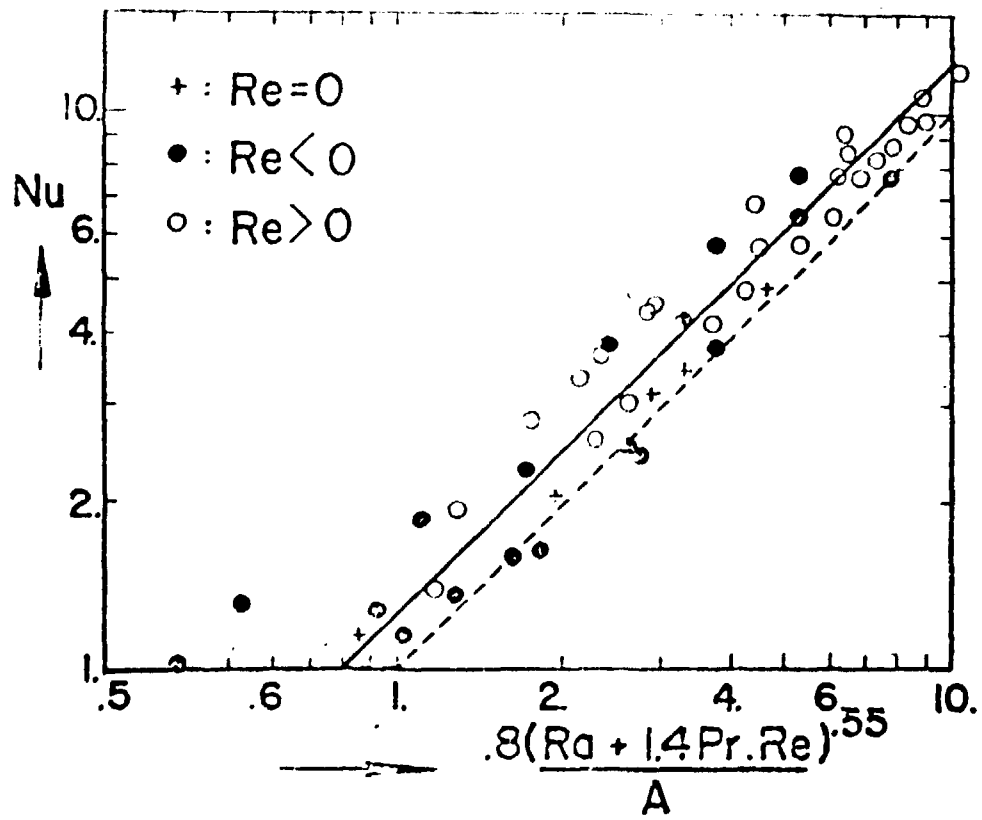


Figure 22 - Vertical Cylinder With Permeation Flow, Correlation of Nusselt Number

$$Nu = .98 \left(\frac{Ra + 1.4 Pr Re}{A} \right)^{.55} \quad (71)$$

$$\text{for } .98 \left(\frac{Ra + 1.4 Pr Re}{A} \right)^{.55} > 1 \text{ and } A \geq 10.$$

For values of $A < 10$, Equation (68) does not hold, and Nu becomes a function of A also (Equation (68) predicts higher values for $A < 10$). The validity of Equation (71) is assumed to be for $A \geq 10$ also. The computer points are for $20 > A \geq 10$.

The error band for the correlation was $\pm 10\%$. Some of the spread is due to the non-ideal behaviour of carbon dioxide. For Helium, the computed Nusselt numbers are expected to show a smaller spread.

4.3 – Rectangular Geometries

Rectangular geometries have been modeled to check the accuracy of the present numerical model. The range of parameters as well as the computation of the heat fluxes and Nu have been similar to what has been discussed earlier for the cylinders.

Closed Hot Wall

Figure 23 shows the temperature and flow fields for a vertical cavity with closed hot wall. The profiles are close to those for a vertical cylinder (Figure 12).

The correlation of Nu with Ra is shown on Figure 24. The computed points agree very well with the correlation of Jannot⁽⁵⁾, obtained from computed results, and also checked with experimental data. The correlation is

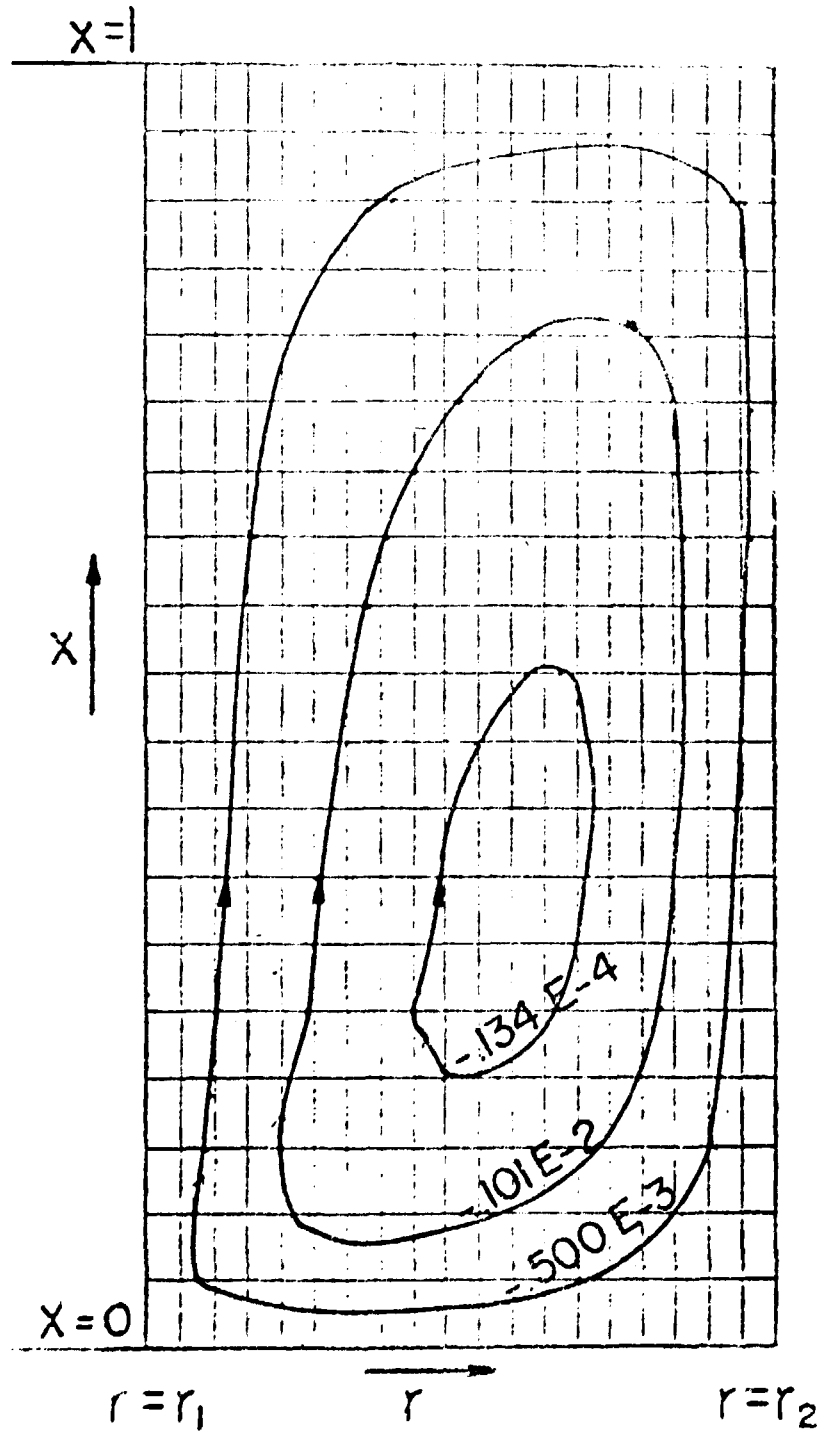
$$Nu = .55 \left(\frac{Ra}{A} \right)^{.5} \quad \text{for } A \geq 10 \quad (72)$$

Permeable Hot Wall

To compare the present model with that of Jannot, FINS has been run with the same p ($p \approx 0$) and same Pr ($Pr \approx 1$). The results of the flow temperature profiles are not shown since they are similar to those for a vertical cylinder. The case with horizontal walls has not been investigated, but the computer program FINS has been provided with that option.

The computed Nu has been plotted vs $.8 \left(\frac{Ra + 1.4 Pr Re}{A} \right)^{.55}$ (Figure 25). It is seen that there is very good agreement between the results of Jannot and those from the present work for a rectangular cavity.

STREAM FUNCTION



—Figure 23-a — Rectangular Geometry Without Permeation Flow, Stream Lines, $Re = 200$, $Nu = 2.45$

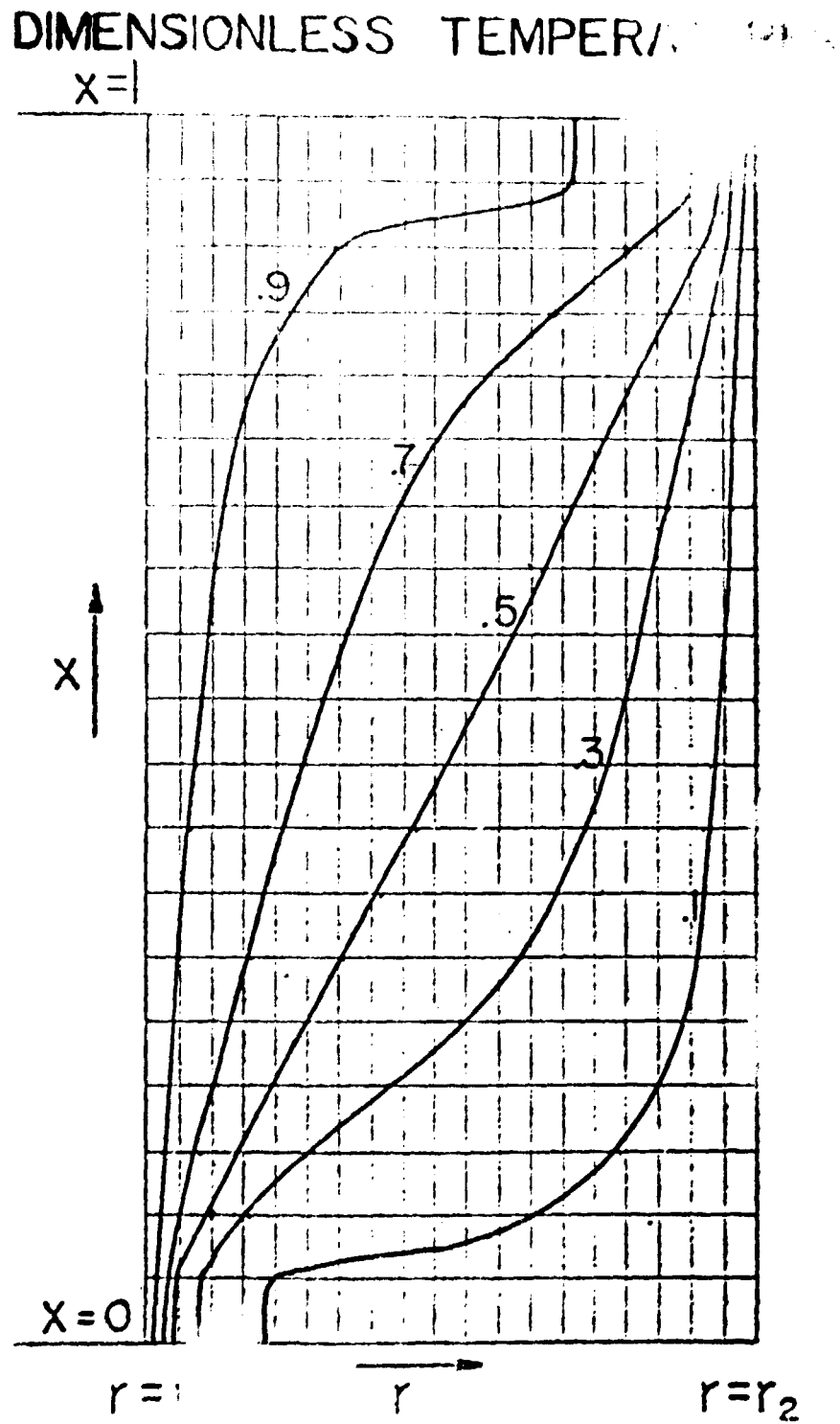


Figure 23-b - Rectangular Geometry Without Permeation Flow, Dimensionless Temperatures, $Ra = 209$, $Nu = 2.45$.

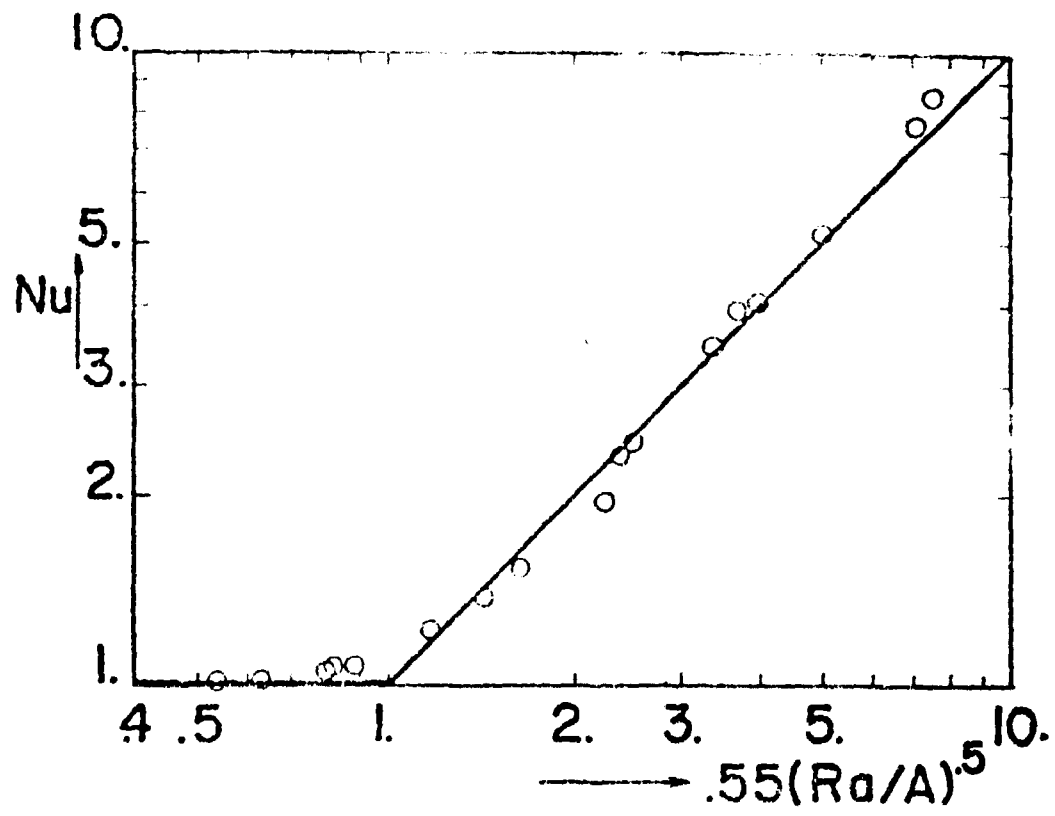


Figure 24 -- Correlation of Overall Nusselt Number, Rectangular Geometry Without Permeation Flow

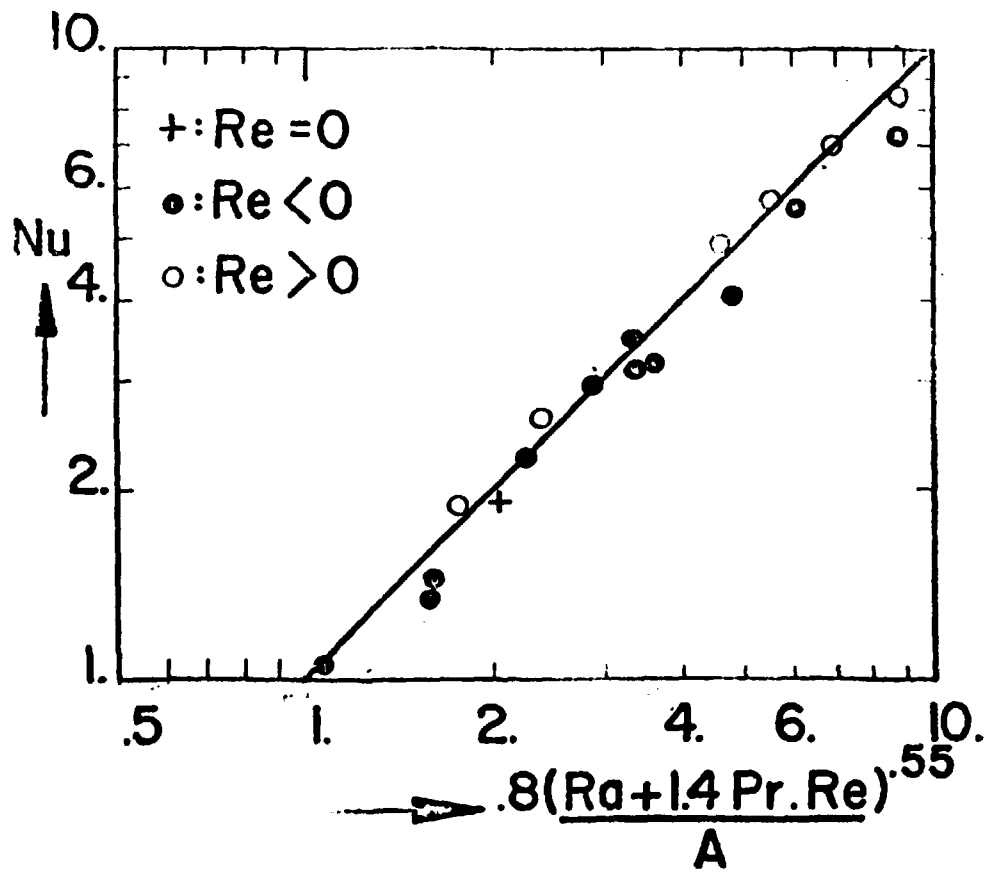


Figure 25 - Correlation of Overall Nusselt Number, Rectangular Geometry, Permeable Hot Wall

5 – CONCLUSION

5.1 – Summary of Results

A numerical solution of flow and heat transfer in fibrous insulation has been obtained by modelling the fibrous insulation as a porous medium. Cylindrical and rectangular geometries with a variety of boundary conditions (vertical or horizontal walls, closed or permeable hot wall, convective boundary conditions or given temperatures at the walls) have been considered. The method allows for variable properties and real gas behaviour through the use of property subroutines.

Excellent agreement has been obtained with published numerical and experimental work on insulation packed between rectangular walls. Comparisons for the cylindrical geometries have not been possible due to the lack of available literature results. However, since the same physical model and the same numerical method have been used for 2D geometries, the computations for cylindrical cavities are expected to be valid also.

The details of flow and heat transfer for natural, forced, and a combination of both have been obtained for a vertical cylinder. Forced convection flows are in the same direction (downward along the cold wall) as natural convection flows when the gas outside the hot wall flows downward and in the opposite direction (upward along the cold wall) when it flows upward. The local Nusselt number along the cold wall increases from bottom to top for natural convection, forced convection with downward gas flow, or both. When the gas flow upwards, and the natural convection is negligible, Nusselt number decreases from bottom to top.

Average effective Nusselt number for the insulation has been correlated with Rayleigh number only for closed hot wall, and with Rayleigh and Reynolds numbers for permeable hot wall. Forced convection flows can reach very high values when the outside pressure gradient is high. Both natural and forced circulation increase with increasing fiber permeability and gas density. The insulation must be protected from forced convection flows and fibers of small permeability (high packing density, small fibre diameter) must be chosen to obtain good insulation between the hot gas and the cold wall. The numerical model presented provides a useful tool to assess the performance of the fibrous insulation by accounting accurately for the natural and forced convection flows.

5.2 – Future Work

Future work concerns the application of the numerical method to experimental work on fibrous insulation to be carried out at IEA Helium Loop. The application involves first the incorporation in the computer program FINS of the test section geometry and boundary conditions, and second the determination of input parameters to the numerical method. The modification of FINS is not difficult, and no convergence problems should be encountered with the test section geometry. The input parameters which must be determined are the permeability and the effective thermal conductivity of the insulation in the absence of forced and natural convection.

NOMENCLATURE

A aspect ratio, defined by

$$A = \frac{1}{d} \text{ for a rectangular cavity}$$

$$A = \frac{1}{r_2 - r_1} \text{ for a vertical cylinder}$$

$$A = \frac{\pi (r_1 + r_2)}{2 (r_2 - r_1)} \text{ for a horizontal cylinder}$$

a thickness of the cover plate

b weight parameter for updating stream function

c weight parameter for updating temperatures

C_p specific heat of the fluid

d thickness of the porous medium for a rectangular cavity $d = r_2 - r_1$ for cylindrical geometries

h heat transfer coefficient

K permeability of porous medium

l, L length of the vertical cylinder or rectangular cavity

Nu Nusselt number = $\frac{hd}{\lambda_r}$

p pressure

Pr Prandtl number = $\frac{\mu_r c_p}{\lambda}$

q heat flux

r radius, radial cylindrical coordinate

Ra Rayleigh number, defined by

$$Ra = \frac{g(T_H - T_c) K \beta}{\nu_r \alpha_r}$$

Where $K = K_r$ for cylindrical and $K = K_v$ for rectangular geometries

Re Reynolds number, defined by

$$Re = \frac{K_x (\partial p_o / \partial x - \rho_o g_x) d}{\mu_r \nu_o}$$

T temperature

v velocity

x axial cylindrical coordinate or cartesian coordinate parallel to the cover plate

y cartesian coordinate perpendicular to cover plate

α thermal diffusivity = $\frac{\lambda}{\rho C_p}$

β thermal coefficient of volume expansion at constant pressure, defined by

$$\beta = -\frac{1}{\rho} \left(\frac{\partial \rho}{\partial T} \right)_p$$

θ angular cylindrical coordinate

λ thermal conductivity of the porous medium with stationary fluid

μ dynamic viscosity of the fluid

ν kinematic viscosity of the fluid

ρ density of the fluid

ψ stream function

Subscripts

C cover plate properties

n iteration number

o free stream quantities

quantities calculated at reference condition

θ, r, x, y properties in θ , r , x , and y directions

Superscripts

node index in r or y directions

node index in x direction

APPENDIX

COMPUTER PROGRAM FINS INPUT INFORMATION AND LISTING

Input Information

- 1) Permeability K (PER in FINS) depends on the fiber diameter, orientation of the fibers, and the insulation porosity. Furber and Davidson⁽³⁾ give their experimental results and an empirical correlation of K with the fiber diameter and insulation porosity. The correlation is not very accurate because it does not account for the fiber orientation. For more accurate determination of K, experimental data on the particular insulation to be modeled is required. The insulation must be packed in the apparatus for K measurement the same way as it is packed in the test section to be simulated by FINS.
- 2) Effective thermal conductivity λ (AKO in FINS) of the insulation with stagnant gas.

λ has been modeled as

$$\lambda = \epsilon \lambda_{\text{gas}} + \lambda_{\text{fiber}} \quad (\text{AKO} = \text{POR} \cdot \text{AK} + \text{AKO})$$

λ_{fiber} (AKO) must be given either from a correlation or from experimental data. It depends on the porosity, fiber diameter fiber orientation, and the thermal conductivity of the fiber material. Thermal conductivity experiments under vacuum yield λ_{fiber} . An approximation to λ_{fiber} has been obtained from available literature data as

$$\lambda_f = 6.0 \text{ E-5 } d \quad (\text{W/ms})$$

where d is the insulation density (Kg/m^3).

- 3) Permeability of the hot wall, K_c (PERCEP in FINS) information on the evaluation of K_c can be found in reference⁽⁵⁾.

- 4) Pressure drop in the gas stream $\frac{\partial p_o}{\partial x}$ (DPDX in FINS). This can be calculated from:

$$\frac{\partial p_o}{\partial x} = \frac{f \rho v^2}{D_H 2 g_c}$$

where f is the friction factors, ρ gas density, and v gas velocity.

COMPILER OPTIONS - NAME= MAIN,OPT=00,LINECNT=60,SIZE=0000K,
 SOURCE,FACDIO,NOLIST,NODECK,LOAD,NOMAP,NOEDIT,NOID,NOXREF
 C FINS (FIBROUS INSULATION) WRITTEN BY AYDIN KONUK,
 C INSTITUTO DE ENERGIA ATOMICA, SAO PAULO, BRAZIL
 C PRESENT VERSION COMPLETED IN MAY 15, 1976.

PURPOSE

C TO OBTAIN A NUMERICAL SOLUTION OF TEMPERATURE AND
 C VELOCITY DISTRIBUTION IN FIBROUS ISOLATION PACKED
 C IN ANNULAR OR RECTANGULAR CAVITIES.

DESCRIPTION INPUT PARAMETERS

GEOMETRY AND BOUNDARY CONDITIONS

C NG=1 -HORIZONTAL CYLINDER
 C NG=2 -VERTICAL CYLINDER
 C NG=3 -RECTANGULAR GEOMETRY, VERTICAL WALLS
 C NG=4 -RECTANGULAR GEOMETRY, HORIZONTAL WALLS,
 C WITH HOT WALL ON TOP
 C NG=5 -RECTANGULAR GEOMETRY, HORIZONTAL WALLS,
 C WITH COLD WALL ON TOP
 C MBC=0 -GIVEN COLD WALL TEMPERATURE T2
 C MBC=1 -GIVEN OUTSIDE TEMPERATURE TOUT AND HEAT
 C TRANSFER COEFFICIENT AGAC
 C NBC=0 -GIVEN HOT WALL TEMPERATURE T1
 C NRC=1 -GIVEN HOT GAS TEMPERATURE TGAS AND HEAT
 C TRANSFER COEFFICIENT AGAH.
 C NPF=0 -CLOSED HOT WALL
 C NPF=1 -PERMEABLE HOT WALL, WITH PERMEABILITY
 C PERCP AND THICKNESS TCP
 C R1 -INSIDE RADIUS. R1=0 FOR RECTANGULAR
 C GEOMETRIES. IN M.
 C R2 -OUTSIDE RADIUS. R2=INSULATION THICKNESS
 C FOR RECTANGULAR GEOMETRIES. IN M.
 C AL -LENTH OF THE INSULATION. FOR HORIZONTAL
 C CYLINDER, AL=3.1416. IN M.
 C PER -PERMEABILITY OF THE INSULATION, M**2
 C GR,GTH -FOR A RECTANGULAR GEOMETRY, COMPONENTS
 C OF THE ACCELERATION OF GRAVITY G.
 C GR=0, GTH=-G FOR VERTICAL WALLS
 C GR=1., GTH=.0 FOR HORIZONTAL WALLS,
 C WITH HOT WALL ON TOP,
 C GR=-1., GTH=.0 FOR HORIZONTAL WALLS,
 C WITH COLD WALL ON TOP.

INSULATION PROPERTIES

C AKM -CONDUCTION THROUGH FIBER
 C POR -POROSITY OF THE INSULATION
 C AKO -EFFECTIVE THERMAL CONDUCTIVITY OF THE
 C INSULATION WHEN THE FLUID IS STAGNANT.
 C DPOX -PRESSURE GRADIENT IN THE GAS STREAM,
 C N/M**3
 C PO -PRESSURE OF THE GAS, BAR
 C V -KINETIC VISCOSITY OF THE GAS, M**2/S
 C AK -THERMAL CONDUCTIVITY OF THE GAS, W/(M*K)
 C CP -HEAT CAPACITY OF THE GAS, W*(M/K)*K
 C RO -DENSITY OF THE GAS, KG/M**3

```

C
C DESCRIPTION OF OUTPUT PARAMETERS
C   J   -NODE INDEX ALONG THE LENGTH OF THE
C         INSULATION. J=1 AT THE BOTTOM
C         J=JJ ON TOP
C   I   -NODE INDEX ACROSS THE THICKNESS OF THE
C         INSULATION. I=1 AT THE HOT WALL
C         I=II AT THE COLD WALL
C
C   TR,TEND -TEMPERATURES FROM THE SOLUTION OF THE
C             ENERGY EQUATION,K
C   T,TLINE -TEMPERATURES USED TO LINEARIZE THE
C             EQUATIONS OF MOMENTUM AND ENERGY.
C   SR,SEND -STREAM FUNCTION FROM THE SOLUTION.
C             OF THE MOMENTUM EQUATION.
C   S,SLINE -STREAM FUNCTION USED TO LINEARIZE
C             THE ENERGY EQUATION.
C   TEMER   -TEND-T AT THE LAST ITERATION.
C   MVR     -RADIAL COMPONENT OF MASS VELOCITY,KG/(S*M**2)
C           MVR IS POSITIVE GOING FROM THE HOT TO THE
C           COLD WALL
C   MVTH    -COMPONENT OF MASS VELOCITY ALONG THE WALLS.
C           POSITIVE GOING UP FOR VERTICAL WALLS.
C   AKEF    -EFFECTIVE THERMAL CONDUCTIVITY AT EACH
C           J. AT THE WALLS, THE TEMPERATURE
C           GRADIENT IS APPROXIMATED WITH LINEAR
C           TEMPERATURE PROFILES.
C   AKEFH   -EFFECTIVE THERMAL CONDUCTIVITY USING
C           THE HEAT FLUX AT THE HOT WALL FROM
C           PARABOLIC APPROXIMATION OF THE
C           TEMPERATURE PROFILE.
C   AKEFC   -SIMILAT TO AKEFH, BUT AT THE COLD WALL.
C
C   AKHC    -ARITHMETIC AVERAGE OF AKEFH AND AKEFC.
C   RNUS    NUSSEL NUMBER, RNUS=AKHC/AKO.
C   RA      RAYLEIGH NUMBER FOR THE INSULATION.
C   RE      REYNOLDS NUMBER FOR THE INSULATION.
C   PR      PRANDTL NUMBER FOR THE INSULATION
C
C NUMERICAL SOLUTION PARAMETERS
C   II     -NUMBER OF NODES ACROSS THE THICKNESS
C           OF THE INSULATION
C   JJ     -NUMBER OF NODES ALONG THE WALLS
C   TW1,TW2 -WEIGHT FACTORS TO UPDATE TEMPERATURES
C   SW1,SW2 -WEIGHT FACTORS TO UPDATE STREAM FUNCTION
C   KMAX   -MAXIMUM NUBER OF ITERATIONS
C   EPT    -CONVERGENCE CRITERION, K
C
C SUBROUTINES AND FUNCTION SUBPROGRAMS USED ARE
C   SUBROUTINE GELB OF IBM SSP TO SOLVE THE SYSTEM OF
C   SIMULTANEOUS LINEAR EQUATIONS
C   FUNCTIONS TO CALCULATE GAS PROPERTIES
C   FUNCTION RGF(P,T) TO COMPUTE DENSITY
C   FUNCTION CPF(P,T) TO COMPUTE SPECIFIC HEAT
C   FUNCTION CLAMP(P,T) TO COMPUTE THERMAL CONDUCTIVITY
C   FUNCTION ETAF(P,T) TO COMPUTE DYNAMIC VISCOSITY
C
C DIMENSIONS

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ISN 0048      M=(I+J)
ISN 0049      N=1
ISN 0050      MOD=II
ISN 0051      EPS=10.**(-9)
ISN 0052      MLD=II
ISN 0053      ME=M*(1+MOD+MLD)- II*(1+II)
C             APD IS THE PRESSURE IN ATM
ISN 0054      APD=.98692*P0
C             CALCULATE INITIAL TEMPERATURE DISTRIBUTION
ISN 0055      DO 30 L=1,M
ISN 0056      TLINE(L)=(T1+T2)/2.
ISN 0057      30 CONTINUE
C
ISN 0058      KOUNT=0
ISN 0059      TB=(T1+T2)/2.
ISN 0060      120 CONTINUE
ISN 0061      SW1=.5
ISN 0062      TW1=.0
ISN 0063      SW2=1.-SW1
ISN 0064      TW2=1.-TW1
ISN 0065      DO 5 J=1,JJ
ISN 0066      DO 6 I=1,II
ISN 0067      TIJ=T(I,J)
ISN 0068      RO(I,J)=ROF(APD,TIJ)
ISN 0069      VI(I,J)=ETAFA(APD,TIJ)/ROF(APD,TIJ)
ISN 0070      AK(I,J)=POR*CLAMF(APD,TIJ)+AKM
ISN 0071      CP(I,J)=CPF(APD,TIJ)
ISN 0072      6 CONTINUE
ISN 0073      5 CONTINUE
ISN 0074      DO 15 L=1,ME
ISN 0075      15 A(L)=.0
C
C             MOMENTUM EQUATION
C
ISN 0076      IF(NG.EQ.2) GO TO 400
ISN 0078      IF(NG.LE.5 .AND. NG.GE.3) GO TO 410
C
C             HORIZONTAL CYLINDER
C
C             FILL IN COEFFICIENT MATRIX A AND
C             RIGHT HAND SIDE VECTOR SR
ISN 0080      K=-II
ISN 0081      L=-1
ISN 0082      DO 1 I=1,II
ISN 0083      L=L+1
ISN 0084      K=K+II+L
ISN 0085      H=II+1-(II-L)
ISN 0086      A(K+H)=1.
ISN 0087      SR(L+1)=.0
ISN 0088      1 CONTINUE
ISN 0089      K=K-1
ISN 0090      DO 10 J=2,JJ1
ISN 0091      L=L+1
ISN 0092      K=K+2*II+1
ISN 0093      H=II+1
ISN 0094      A(K+H)=1.
ISN 0095      SR(L+1)=.0
ISN 0096      DO 11 I=1,II

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ISN 0097      L=L+1
ISN 0098      K=K+2*I+1
ISN 0099      H=I+1
ISN 0100      A(K+H)=- (V(I,J-1)+2.*V(I,J)+V(I,J+1))/(2.*DTH2*R(I))
              1-(P(I+1)*V(I+1,J)+2.*R(I)*V(I,J)+R(I-1)*V(I-1,J))/(2.*DR2)
ISN 0101      H=1
ISN 0102      A(K+H)=(V(I,J-1)+V(I,J))/(2.*DTH2*R(I))
ISN 0103      H=2*I+1
ISN 0104      A(K+H)=(V(I,J)+V(I,J+1))/(2.*DTH2*R(I))
ISN 0105      H=I
ISN 0106      A(K+H)=(R(I)*V(I,J)+R(I-1)*V(I-1,J))/(2.*DR2)
ISN 0107      H=I+2
ISN 0108      A(K+H)=(R(I+1)*V(I+1,J)+R(I)*V(I,J))/(2.*DR2)
ISN 0109      DROTH=(R(I,J+1)-R(I,J-1))/(2.*DTH)
ISN 0110      DROR=(R(I+1,J)-R(I-1,J))/(2.*DR)
ISN 0111      SR(L+1)=G*PER*(COS(TH(J))*DROTH+R(I)*SIN(TH(J))*DROR)
ISN 0112      11 CONTINUE
ISN 0113      L=L+1
ISN 0114      K=K+2*I+1
ISN 0115      H=I+1
ISN 0116      A(K+H)=1.
ISN 0117      SR(L+1)=.0
ISN 0118      10 CONTINUE
ISN 0119      DO 2 I=1,I1
ISN 0120      L=L+1
ISN 0121      K=K+M+I+1-L
ISN 0122      H=I+1
ISN 0123      A(K+H)=1.
ISN 0124      SR(L+1)=.0
ISN 0125      2 CONTINUE
ISN 0126      CALL GELBISR,A,M,N,MUD,MLD,EPS,IER)
ISN 0127      KOUNT=KOUNT+1
C             PRINT STREAM FUNCTION
C
ISN 0128      WRITE(6,200) KOUNT
ISN 0129      WRITE(6,218) IER
ISN 0130      WRITE(6,201) SR
C             UPDATE STREAM FUNCTION
C
ISN 0131      IF(KOUNT.EQ.1) GO TO 71
ISN 0132      DO 70 L= 1,M
ISN 0133      70 SLINE(L)=SW1*SLINE(L)+SW2*SR(L)
ISN 0134      GO TO 73
ISN 0135      71 DO 72 L=1,M
ISN 0136      72 SLINE(L)=SR(L)
ISN 0137      73 CONTINUE
C             ENERGY EQUATION
ISN 0138      DO 16 L=1,ME
ISN 0139      16 A(L)=.0
C
C             FILL IN COEFFICIENT MATRIX A AND
C             RIGHT HAND SIDE VECTOR TR
ISN 0140      IF(INBC.EQ.1)GOTO 42
ISN 0141      A(1)=1.
ISN 0142      TR(1)=T1
ISN 0143      GO TO 43
ISN 0144      42 CONTINUE
ISN 0145      A(1)=3.*AK(1,1)/(2.*DR)+AGAM
ISN 0146
ISN 0147

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ISN 0165      A(I,1)=AG*(1,1)/DR
ISN 0166      A(I,1)=A(I,1)/(2.*DR)
ISN 0167      A(I,1)=AG*TA5
ISN 0168      TR(I,1)=0
ISN 0169      H=I
ISN 0170      DO 110 J=2, I1
ISN 0171      H=I+1
ISN 0172      K=K+1+1
ISN 0173      A(K+H)=A(I,1)
ISN 0174      A(K+H)=1.
ISN 0175      TR(I+1)=0
ISN 0176      A(K+H)=A(I,1)/(2.*DR)
ISN 0177      TR(I+1)=AG*TA5
ISN 0178      45 CONTINUE
ISN 0179      K=K+1
ISN 0180      DO 110 J=2, J1
ISN 0181      H=I+1
ISN 0182      K=K+2+1+1
ISN 0183      IF (TR(I,1)) GOTO 52
ISN 0184      A=I+1
ISN 0185      A(K+H)=1.
ISN 0186      TR(I+1)=1
ISN 0187      GO TO 53
ISN 0188      51 CONTINUE
ISN 0189      H=I+1
ISN 0190      A(K+H)=3.*A(I,J)/(2.*DR)+AG*H
ISN 0191      H=I+2
ISN 0192      A(K+H)=-2.*A(I,J)/DR
ISN 0193      H=I+3
ISN 0194      A(K+H)=A(I,J)/(2.*DR)
ISN 0195      TR(I+1)=AG*H*TA5
ISN 0196      54 CONTINUE
ISN 0197      DO 111 I=2, I1
ISN 0198      L=L+1
ISN 0199      K=K+2+1+1
ISN 0200      H=I+1
ISN 0201      A(K+H)=(R(I+1)*A(I+1,J)+2.*R(I)*A(I,J)+R(I-1)*A(I-1,J))/
ISN 0202      (2.*DR2)+(A(I,J+1)+2.*A(I,J)+A(I,J-1))/(2.*R(I))*DTH2
ISN 0203      H=I
ISN 0204

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ISN 0205      A(K+H)=(S(I+1,J)-S(I-1,J))*CP(I,J-1)/(4.*DTH*DR)
              1-(AK(I,J)+AK(I,J-1))/(2.*R(I)*DTH2)
              H=2*I+1
ISN 0206      A(K+H)=-(S(I+1,J)-S(I-1,J))*CP(I,J+1)/(4.*DTH*DR)
ISN 0207      1-(AK(I,J+1)+AK(I,J))/(2.*R(I)*DTH2)
              H=I+1
ISN 0208      A(K+H)=-(S(I,J+1)-S(I,J-1))*CP(I-1,J)/(4.*DTH*DR)
ISN 0209      1-(R(I)*AK(I,J)+R(I-1)*AK(I-1,J))/(2.*DR2)
              H=I+2
ISN 0210      A(K+H)=(S(I,J+1)-S(I,J-1))*CP(I+1,J)/(4.*DTH*DR)
ISN 0211      1-(R(I+1)*AK(I+1,J)+R(I)*AK(I,J))/(2.*DR2)
ISN 0212      TR(L+1)=.0
ISN 0213      111 CONTINUE

C
ISN 0214      L=L+1
ISN 0215      K=K+2*I+1
ISN 0216      IF(MBC.EQ.1)GOTO 55
ISN 0218      H=I+1
ISN 0219      A(K+H)=1.
ISN 0220      TR(L+1)=I2
ISN 0221      GOTO 56
ISN 0222      55 CONTINUE
ISN 0223      H=I+1
ISN 0224      A(K+H)=AGAC+3.*AK(I,J)/(2.*DR)
ISN 0225      H=I
ISN 0226      A(K+H)=-2.*AK(I,J)/DR
ISN 0227      H=I-1
ISN 0228      A(K+H)=AK(I,J)/(2.*DR)
ISN 0229      TR(L+1)=AGAC*TOUT
ISN 0230      56 CONTINUE
ISN 0231      110 CONTINUE

C
ISN 0232      L=L+1
ISN 0233      K=K+2*I+1
ISN 0234      IF(NBC.EQ.1)GOTO 62
ISN 0236      H=I+1
ISN 0237      A(K+H)=1.
ISN 0238      TR(L+1)=I1
ISN 0239      GO TO 63
ISN 0240      62 CONTINUE
ISN 0241      H=I+1
ISN 0242      A(K+H)=3.*AK(I,JJ)/(2.*DR)+AGAH
ISN 0243      H=I+2
ISN 0244      A(K+H)=-2.*AK(I,JJ)/DR
ISN 0245      H=I+3
ISN 0246      A(K+H)=AK(I,JJ)/(2.*DR)
ISN 0247      TR(L+1)=AGAH*TGAS
ISN 0248      63 CONTINUE

C
ISN 0249      DO 102 I=2,I1
ISN 0250      L=L+1
ISN 0251      K=K+M+I+1-L
ISN 0252      H=I+1
ISN 0253      A(K+H)=1.
ISN 0254      H=1
ISN 0255      A(K+H)=-1.
ISN 0256      TR(L+1)=.0
ISN 0257      102 CONTINUE

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C
ISN 0258      I=I+1
ISN 0259      K=K+M+I+1-L
ISN 0260      IF(MHC.EQ.1)GOTO 65
ISN 0262      H=I+1
ISN 0263      A(K+H)=1.
ISN 0264      TR(L+1)=T2
ISN 0265      GOTO 66
ISN 0266      65 CONTINUE
ISN 0267      H=I+1
ISN 0268      A(K+H)=AGAC+3.*AK(II,JJ)/(2.*DR)
ISN 0269      H=II
ISN 0270      A(K+H)=-2.*AK(II,JJ)/DR
ISN 0271      H=II-1
ISN 0272      A(K+H)=AK(II,JJ)/(2.*DR)
ISN 0273      TR(L+1)=AGAC*TOUT
ISN 0274      66 CONTINUE
C
C      MATRIX A IS FILLED IN
C
ISN 0275      CALL GFLB(TR,A,M,N,MUD,MID,EPS,ITER)
C
C      WRITE TEMPERATURES
C
ISN 0276      WRITE(6,210)
ISN 0277      WRITE(6,211) TR
C
C
C      CHECK THE TEMPERATURES FOR CONVERGENCE
C
ISN 0278      DO 1300 L=1,M
ISN 0279      IF(DABS(TR(L)-TLINE(L)).GT.EPT) GO TO 116
ISN 0281      1300 CONTINUE
ISN 0282      GO TO 115
C
C      CHECK IF KOUNT IS LESS THAN KMAX
C
ISN 0283      116 IF(KOUNT.EQ.KMAX) GO TO 115
C
C      UPDATE TEMPERATURES
ISN 0285      DO 80 L=1,M
ISN 0286      80 TLINE(L)=TW1*TLINE(L)+TW2*TR(L)
ISN 0287      GO TO 120
ISN 0288      115 CONTINUE
C
ISN 0289      T1=.0
ISN 0290      T2=.0
ISN 0291      DO 97 J=1,JJ
ISN 0292      T1=TEND(1,J)/JJ+T1
ISN 0293      T2=TEND(II,J)/JJ+T2
ISN 0294      97 CONTINUE
ISN 0295      TB=(T1+T2)/2.
C
C      CALCULATE VELOCITIES
C      MVR MASS VELOCITY IN THE R DIRECTION
C      MVTH MASS VELOCITY IN THE TH DIRECTION
ISN 0296      DO 82 J=1,JJ1
ISN 0297      DO 81 I=1,II
ISN 0298      SEND(I,J)=S(I,J)
ISN 0299      81 MVR(I,J)=(SEND(I,J+1)-SEND(I,J))/(R(I)*DTH)

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ISN 0300      92 CONTINUE
ISN 0301      DO 85 I=1,III
ISN 0302      DO 86 J=1,JJ
ISN 0303      86 MVT(I,J)=-TEND(I+1,J)-SEND(I,J)/DR
ISN 0304      85 CONTINUE
C             CALCULATE THE HEAT FLUX AND K EFFECTIVE
ISN 0305      QT(I)=.0
ISN 0306      QPHT=.0
C
ISN 0307      DO 90 J=1,JJ1
ISN 0308      T(I,J)=TEND(I,J)
ISN 0309      QPH(J)=-(-T(3,J)-T(3,J+1)+4.*(T(2,J)+T(2,J+1))-3.*(T(1,J)+
1T(1,J+1)))*(AK(1,J)+AK(1,J+1))/(8.*DR)*R(I)*DTH
C
ISN 0310      CONV(I,J)=.0
ISN 0311      CONO(I,J)=(T(1,J)+T(1,J+1)-T(2,J)-T(2,J+1))*(AK(1,J)+AK(1,J+1))
1/(4.*DR)*R(I)*DTH
C
ISN 0312      Q(I,J)=CONV(I,J)+CONO(I,J)
ISN 0313      QPHT=QPHT+QPH(J)
ISN 0314      90 QT(I)=QT(I)+Q(I,J)
C
ISN 0315      DO 91 I=2,III
ISN 0316      QT(I)=.0
ISN 0317      DO 92 J=1,JJ1
ISN 0318      T(I,J)=TEND(I,J)
ISN 0319      CONO(I,J)=(T(I-1,J)+T(I-1,J+1)-T(I+1,J)-T(I+1,J+1))
1*(AK(I,J)+AK(I,J+1))/(8.*DR)*R(I)*DTH
ISN 0320      CONV(I,J)= MVR(I,J)*(T(I,J)+T(I,J+1))*(CP(I,J)+CP(I,J+1))/4.
1*R(I)*DTH
ISN 0321      Q(I,J)=CONV(I,J)+CONO(I,J)
ISN 0322      92 QT(I)=QT(I)+Q(I,J)
ISN 0323      91 CONTINUE
C
ISN 0324      QT(III)=.0
ISN 0325      QPCT=.0
ISN 0326      DO 93 J=1,JJ1
ISN 0327      T(II,J)=TEND(II,J)
ISN 0328      QPC(J)=-(-T(II-2,J)-T(II-2,J+1)+4.*(T(II-1,J)+T(II-1,J+1))
1-3.*(T(II,J)+T(II,J+1)))*(AK(II,J)+AK(II,J+1))/(8.*DR)*R(II)*DTH
C
ISN 0329      CONV(II,J)=.0
ISN 0330      CONO(II,J)=(T(II-1,J)+T(II-1,J+1)-T(II,J)-T(II,J+1))
1*(AK(II,J)+AK(II,J+1))/(4.*DR)*R(II)*DTH
ISN 0331      Q(II,J)=CONV(II,J)+CONO(II,J)
ISN 0332      QPCT=QPCT+QPC(J)
ISN 0333      93 QT(II)=QT(II)+Q(II,J)
C
ISN 0334      QTAV=.0
ISN 0335      DO 96 I=1,II
ISN 0336      QTAV=QTAV+QT(I)/II
ISN 0337      96 CONTINUE
ISN 0338      AKAV=QTAV*ALOG(R2/R1)/((T1-T2)*3.1416)
ISN 0339      AKFC=QPCT*ALOG(R2/R1)/((T1-T2)*3.1416)
ISN 0340      AKFHM=QPHT*ALOG(R2/R1)/((T1-T2)*3.1416)
ISN 0341      GO TO 450
ISN 0342      400 CONTINUE
C

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C      VERTICAL CYLINDER
C
C      FILL IN COEFFICIENT MATRIX A AND
C      RIGHT HAND SIDE VECTOR SR
ISN 0343      K=-1
ISN 0344      L=1
ISN 0345      DO 601 I=1,II
ISN 0346      I=I+1
ISN 0347      K=K+I+1
ISN 0348      H=I+1-(II-L)
ISN 0349      A(K+H)=1.
ISN 0350      SR(L+1)=.0
ISN 0351      601 CONTINUE
ISN 0352      K=K-1
ISN 0353      DO 610 J=2,JJ1
ISN 0354      L=L+1
ISN 0355      K=K+2*I+1
ISN 0356      IF(NPE,EQ.1) GO TO 612
ISN 0358      H=I+1
ISN 0359      A(K+H)=1.
ISN 0360      SR(L+1)=.0
ISN 0361      GO TO 611
ISN 0362      612 CONTINUE
ISN 0363      H=1
ISN 0364      A(K+H)=C1*(V(I,J)+V(I,J-1))/(V(I,J))
ISN 0365      H=2*I+1
ISN 0366      A(K+H)=C1*(V(I,J)+V(I,J+1))/(V(I,J))
ISN 0367      H=I+1
ISN 0368      A(K+H)=-1./DR-C1*(V(I,J-1)+2.*V(I,J)+V(I,J+1))/(V(I,J))
ISN 0369      H=I+2
ISN 0370      A(K+H)=1./DR
ISN 0371      SR(L+1)=R(I)*PER/V(I,J)*(G*RO(I,J)+DPDX)
ISN 0372      511 CONTINUE
ISN 0373      DO 611 I=2,III
ISN 0374      L=L+1
ISN 0375      K=K+2*I+1
ISN 0376      H=I+1
ISN 0377      A(K+H)=-1*(V(I,J-1)+2.*V(I,J)+V(I,J+1))/(2.*DTH2*R(I))
      1-(V(I+1,J)/R(I+1)+2.*V(I,J)/R(I)+V(I-1,J)/R(I-1))/(2.*DR2)
ISN 0378      H=1
ISN 0379      A(K+H)=(V(I,J-1)+V(I,J))/(2.*DTH2*R(I))
ISN 0380      H=2*I+1
ISN 0381      A(K+H)=(V(I,J)+V(I,J+1))/(2.*DTH2*R(I))
ISN 0382      H=I
ISN 0383      A(K+H)=(V(I,J)/R(I)-V(I-1,J)/R(I-1))/(2.*DR2)
ISN 0384      H=I+2
ISN 0385      A(K+H)=(V(I+1,J)/R(I+1)+V(I,J)/R(I))/(2.*DR2)
ISN 0386      DRDTH=(RO(I,J+1)-RO(I,J-1))/(2.*DTH)
ISN 0387      DRDR=(RO(I+1,J)-RO(I-1,J))/(2.*DR)
ISN 0388      SR(L+1)=G*PER*DRDR
ISN 0389      611 CONTINUE
ISN 0390      L=L+1
ISN 0391      K=K+2*I+1
ISN 0392      H=I+1
ISN 0393      A(K+H)=1.
ISN 0394      SR(L+1)=.0
ISN 0395      610 CONTINUE
ISN 0396      DO 602 I=1,II

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ISN 0397      L=I+L
ISN 0398      K=K+M+II+1-L
ISN 0399      H=II+1
ISN 0400      A(K+H)=1.
ISN 0401      SR(L+1)=.0
ISN 0402      602 CONTINUE
ISN 0403      CALL GELB(SR,A,M,N,MUD,MLD,EPS,IER)
ISN 0404      KOUNT=KOUNT+1
              PRINT STREAM FUNCTION
C
C
ISN 0405      WRITE(6,200) KOUNT
ISN 0406      WRITE(6,218) IER
ISN 0407      WRITE(6,201) SR
              UPDATE STREAM FUNCTION
C
C
ISN 0408      IF(KOUNT.EQ.1) GO TO 671
ISN 0410      DO 670 L=1,M
ISN 0411      670 SLINE(L)=SW1*SLINE(L)+SW2*SR(L)
ISN 0412      GO TO 673
ISN 0413      671 DO 672 L=1,M
ISN 0414      672 SLINE(L)=SR(L)
ISN 0415      673 CONTINUE
              ENERGY EQUATION
C
ISN 0416      DO 616 L=1,ME
ISN 0417      616 A(L)=.0
C
C
              FILL IN COEFFICIENT MATRIX A AND
              RIGHT HAND SIDE VECTOR TR
C
ISN 0418      IF(NBC.EQ.1) GO TO 642
ISN 0420      A(1)=1.
ISN 0421      TR(1)=T1
ISN 0422      GO TO 643
ISN 0423      642 CONTINUE
ISN 0424      A(1)=3.*AK(1,1)/(2.*DR)+AGAH
ISN 0425      A(2)=-2.*AK(1,1)/DR
ISN 0426      A(3)=AK(1,1)/(2.*DR)
ISN 0427      TR(1)=AGAH*TGAS
ISN 0428      643 CONTINUE
ISN 0429      L=0
ISN 0430      K=0
ISN 0431      DO 6101 I=2,III
ISN 0432      L=L+1
ISN 0433      K=K+II+L
ISN 0434      H=II+1-(II-L)
ISN 0435      A(K+H)=1.
ISN 0436      H=2*II+1-(II-L)
ISN 0437      A(K+H)=-1.
ISN 0438      TR(L+1)=.0
ISN 0439      6101 CONTINUE
C
ISN 0440      L=L+1
ISN 0441      K=K+II+L
ISN 0442      IF(MBC.EQ.1) GO TO 645
ISN 0444      H=II+1-(II-L)
ISN 0445      A(K+H)=1.
ISN 0446      TR(L+1)=T2
ISN 0447      GO TO 646
ISN 0448      645 CONTINUE

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ISN 0449      H=I+1-(I+1)
ISN 0450      A(K+H)=AGAC+3.*AK(I+1,1)/(2.*DR)
ISN 0451      H=I-(I+1)
ISN 0452      A(K+H)=-2.*AK(I+1,1)/DR
ISN 0453      H=I-1-(I+1)
ISN 0454      A(K+H)=AK(I+1,1)/(2.*DR)
ISN 0455      TR(L+1)=AGAC*THUT
ISN 0456      646 CONTINUE
ISN 0457      K=K+1
ISN 0458      DO 6110 J=2, J1
ISN 0459      L=L+1
ISN 0460      K=K+2*I+1
ISN 0461      IF (NRC.EQ.1) GO TO 652
ISN 0463      H=I+1
ISN 0464      A(K+H)=1.
ISN 0465      TR(L+1)=I

C
ISN 0466      G) T) 653
ISN 0467      652 CONTINUE
ISN 0468      H=I+1
ISN 0469      A(K+H)=3.*AK(I, J)/(2.*DR)+AGAH
ISN 0470      H=I+2
ISN 0471      A(K+H)=-2.*AK(I, J)/DR
ISN 0472      H=I+3
ISN 0473      A(K+H)=AK(I, J)/(2.*DR)
ISN 0474      TR(L+1)=AGAH*TGAS
ISN 0475      653 CONTINUE
ISN 0476      DO 6111 I=2, I1
ISN 0477      L=L+1
ISN 0478      K=K+2*I+1
ISN 0479      H=I+1
ISN 0480      A(K+H)=(R(I+1)*AK(I+1, J)+7.*R(I)*AK(I, J)+R(I-1)*AK(I-1, J))/
1(2.*DR2)+(AK(I, J+1)+2.*AK(I, J)+AK(I, J-1))*R(I)/(2.*DTH2)
ISN 0481      H=I
ISN 0482      A(K+H)=(S(I+1, J)-S(I-1, J))*CP(I, J-1)/(4.*DTH*DR)
1-(AK(I, J+1)+AK(I, J-1))*R(I)/(2.*DTH2)
ISN 0483      H=2*I+1
ISN 0484      A(K+H)=-(S(I+1, J)-S(I-1, J))*CP(I, J+1)/(4.*DTH*DR)
1-(AK(I, J+1)+AK(I, J-1))*R(I)/(2.*DTH2)
ISN 0485      H=I
ISN 0486      A(K+H)=-(S(I, J+1)-S(I, J-1))*CP(I-1, J)/(4.*DTH*DR)
1-(R(I)*AK(I, J)+R(I-1)*AK(I-1, J))/(2.*DR2)
ISN 0487      H=I+2
ISN 0488      A(K+H)=(S(I, J+1)-S(I, J-1))*CP(I+1, J)/(4.*DTH*DR)
1-(R(I+1)*AK(I+1, J)+R(I)*AK(I, J))/(2.*DR2)
ISN 0489      TR(L+1)=0
ISN 0490      6111 CONTINUE

C
ISN 0491      L=L+1
ISN 0492      K=K+2*I+1
ISN 0493      IF (NRC.EQ.1) GO TO 655
ISN 0495      H=I+1
ISN 0496      A(K+H)=1.
ISN 0497      TR(L+1)=T2
ISN 0498      G) T) 656
ISN 0499      655 CONTINUE
ISN 0500      H=I+1
ISN 0501      A(K+H)=AGAL+3.*AK(I+1, J)/(2.*DR)

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ISN 0502      H=II
ISN 0503      A(K+H)=-2.*AK(II,J)/DR
ISN 0504      H=II-1
ISN 0505      A(K+H)=AK(II,J)/(2.*DR)
ISN 0506      TR(L+1)=AGAC*TOUT
ISN 0507      656 CONTINUE
ISN 0508      6110 CONTINUE
C
ISN 0509      L=L+1
ISN 0510      K=K+2*II+1
ISN 0511      IF(NBC.EQ.1) GO TO 662
ISN 0513      H=II+1
ISN 0514      A(K+H)=1.
ISN 0515      TR(L+1)=T1
ISN 0516      GO TO 663
ISN 0517      662 CONTINUE
ISN 0518      H=II+1
ISN 0519      A(K+H)=3.*AK(1,JJ)/(2.*DR)+AGAH
ISN 0520      H=II+2
ISN 0521      A(K+H)=-2.*AK(1,JJ)/DR
ISN 0522      H=II+3
ISN 0523      A(K+H)=AK(1,JJ)/(2.*DR)
ISN 0524      TR(L+1)=AGAH*TGAS
ISN 0525      663 CONTINUE
C
ISN 0526      DO 6102 I=2,III
ISN 0527      L=L+1
ISN 0528      K=K+M+II+1-L
ISN 0529      H=II+1
ISN 0530      A(K+H)=1.
ISN 0531      H=1
ISN 0532      A(K+H)=-1.
ISN 0533      TR(L+1)=.0
ISN 0534      6102 CONTINUE
C
ISN 0535      L=L+1
ISN 0536      K=K+1+II+1-L
ISN 0537      IF(MBC.EQ.1) GO TO 665
ISN 0539      H=II+1
ISN 0540      A(K+H)=1.
ISN 0541      TR(L+1)=T2
ISN 0542      GO TO 666
ISN 0543      665 CONTINUE
ISN 0544      H=II+1
ISN 0545      A(K+H)=AGAC+3.*AK(II,JJ)/(2.*DR)
ISN 0546      H=II
ISN 0547      A(K+H)=-2.*AK(II,JJ)/DR
ISN 0548      H=II-1
ISN 0549      A(K+H)=AK(II,JJ)/(2.*DR)
ISN 0550      TR(L+1)=AGAC*TOUT
ISN 0551      666 CONTINUE
C
C
C
C
ISN 0552      CALL GELB(TR,A,M,N,MUO,MLD,EPS,IER)
C
C
C
C
ISN 0553      WRITE(6,210)

```

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ISN 0554      WPI(I,6,211) TR
C
C
C      CHECK THE TEMPERATURES FOR CONVERGENCE
C
ISN 0555      DO 630 I=1,M
ISN 0556      IF (DABS(T(I)-T1NE(I)).GT.EPT) GO TO6116
ISN 0558      630 CONTINUE
ISN 0559      GO TO 6115
C
C      CHECK IF COUNT IS LESS THAN KMAX
C
ISN 0560      6116 IF (COUNT.EQ.KMAX) GO TO 6115
C
C      UPDATE TEMPERATURES
ISN 0562      DO 640 I=1,M
ISN 0563      640 T1NE(I)=TW1*T1NE(I)+TW2*TR(I)
ISN 0564      GO TO 120
ISN 0565      6115 CONTINUE
C
ISN 0566      T1=.0
ISN 0567      T2=.0
ISN 0568      DO 697 J=1,JJ
ISN 0569      T1=TND(I,J)/DT+T1
ISN 0570      T2=TND(I,J)/DJ+T2
ISN 0571      697 CONTINUE
ISN 0572      TR=(T1+T2)/2.
C
C      CALCULATE VELOCITIES
C      MVR - MASS VELOCITY IN THE R DIRECTION
C      MVT - MASS VELOCITY IN THE TH DIRECTION
ISN 0573      DO 682 J=1,JJ
ISN 0574      DO 681 I=1,II
ISN 0575      S(I,J)=S(I,J)
ISN 0576      681 MVR(I,J)=(S(I,J+1)-S(I,J))/R(I)*DTH
ISN 0577      682 CONTINUE
ISN 0578      DO 685 I=1,III
ISN 0579      DO 686 J=1,JJ
ISN 0580      686 MVT(I,J)=-(S(I+1,J)-S(I,J))/L.5*DR*(R(I)+R(I+1))
ISN 0581      685 CONTINUE
C
C      CALCULATE THE HEAT FLUX AND K EFFECTIVE
ISN 0582      QTEI=.0
ISN 0583      QPHI=.0
C
ISN 0584      DO 670 J=1,JJ
ISN 0585      T(I,J)=TND(I,J)
C
ISN 0586      COND(I,J)=(T(I,J)+T(I,J+1)-T(2,J)-T(2,J+1))*(AK(I,J)+AK(I,J+1))
ISN 0587      CONV(I,J)=MVR(I,J)*(T(I,J)+T(I,J+1))*(CP(I,J)+CP(I,J+1))/4.
ISN 0588      IF (MVR(I,J).LT. .0) CONV(I,J)=MVR(I,J)*(T(I,J)+T(2,J)
ISN 0590      +T(I,J+1)+T(2,J+1))*(CP(I,J)+CP(I,J+1))/8.*DTH*2.*3.1416*R(I)
ISN 0591      QPHI(J)=-(-T(3,J)-T(3,J+1)+4.*T(2,J)+T(2,J+1))-3.*(T(I,J)+T(I,J+1)
ISN 0592      )*(AK(I,J)+AK(I,J+1))/8.*DR)*DTH*2.*3.1416*R(I)
ISN 0591      QPHI(J)=QPHI(J)+CONV(I,J)
C
ISN 0592      J(T,1)=CONV(I,J)+COND(I,J)

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ISN 0593      QPHT=QPHT+QPH(J)
ISN 0594      690 QT(I)=QT(I)+Q(I,J)
C
ISN 0595      DO 691 I=2, I11
ISN 0596      JT(I)=.0
ISN 0597      691 J=1, J11
ISN 0598      T(I,J)=TEND(I,J)
ISN 0599      CONDI(J)=(T(I-1,J)+T(I-1,J+1)-T(I+1,J)-T(I+1,J+1))
ISN 0600      1*(AK(I,J)+AK(I,J+1))/(4.*DP)*DTH*2.*3.1416*RI
ISN 0601      CONVI(J)=MVR(I,J)*(T(I,J)+T(I,J+1))*(CP(I,J)+CP(I,J+1))/4.
ISN 0602      1*DTH*2.*3.1416*RI
ISN 0603      692 QT(I)=QT(I)+Q(I,J)
ISN 0604      691 CONTINUE
ISN 0605      DO 693 J=1, J11
ISN 0606      T(I,J)=TEND(I,J)
ISN 0607      QPC(J)=(-T(I-2,J)-T(I-2,J+1)+4.*(T(I-1,J)+T(I-1,J+1))-3.*(T(I-
ISN 0608      1,J)+T(I,J+1))*(AK(I,J)+AK(I,J+1))/(4.*DP)*DTH*2.*3.1416*RI
C
ISN 0609      CONVI(J)=.0
ISN 0610      CONDI(J)=(T(I-1,J)+T(I-1,J+1)-T(I,J)-T(I,J+1))
ISN 0611      1*(AK(I,J)+AK(I,J+1))/(4.*DP)*DTH*2.*3.1416*RI
ISN 0612      Q(I,J)=CONVI(J)+CONDI(J)
C
ISN 0613      JPCT=QPCT+QPC(J)
ISN 0614      693 QT(I)=QT(I)+Q(I,J)
ISN 0615      DO 694 I=1, I1
ISN 0616      QTAV=QTAV+QT(I)/I1
ISN 0617      694 CONTINUE
ISN 0618      AKAV=QTAV*ALOG(ER2/RI)/(I1-1)*3.1416*2.*AL
ISN 0619      AKEFH=QPHT*ALOG(ER2/RI)/(I1-1)*3.1416*2.*AL
ISN 0620      AKFC=QPCT*ALOG(ER2/RI)/(I1-1)*3.1416*2.*AL
C
ISN 0621      GO TO 450
ISN 0622      410 CONTINUE
C
ISN 0623      RECTANGULAR GEOMETRIES
C
ISN 0624      FILL IN COEFFICIENT MATRIX A AND
ISN 0625      RIGHT HAND SIDE VECTOR S
C
ISN 0626      K=-11
ISN 0627      L=-1
ISN 0628      DO 701 I=1, I1
ISN 0629      L=L+1
ISN 0630      K=K+1+L
ISN 0631      H=I+1-(I-L)
ISN 0632      A(K+H)=1.
ISN 0633      S(I+1)=.0
ISN 0634      701 CONTINUE
ISN 0635      K=K-1
ISN 0636      DO 710 J=2, J11
ISN 0637      L=L+1
ISN 0638      K=K+2+1+L
ISN 0639      IF(NPF.EQ.1) GO TO 712
ISN 0640      H=I+1
ISN 0641      A(K+H)=1.
ISN 0642      S(I+1)=.0

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ISN 0638          DO 713
ISN 0639          712 CONTINUE
ISN 0640          H=I
ISN 0641          A(K+H)=C1*(V(I,J)+V(I,J-1))/(V(I,J))
ISN 0642          H=2*I+1
ISN 0643          A(K+H)=C1*(V(I,J)+V(I,J+1))/(V(I,J))
ISN 0644          H=I+1
ISN 0645          A(K+H)=-1./DR-C1*(V(I,J-1)+2.*V(I,J)+V(I,J+1))/(V(I,J))
ISN 0646          H=I+2
ISN 0647          A(K+H)=1./DR
ISN 0648          SK(L+1)=PER/V(I,J)*(1-GTH*RO(I,J)+DPOX+TCP*GR)
ISN 0649          713 CONTINUE
ISN 0650          DO 711 I=2,I11
ISN 0651          L=L+1
ISN 0652          K=K+2*I+1
ISN 0653          H=I+1
ISN 0654          A(K+H)=- (V(I,J-1)+2.*V(I,J)+V(I,J+1))/(2.*DTH2)
          I-(V(I+1,J)+2.*V(I,J)+V(I-1,J))/(2.*DR2)
          H=I
          A(K+H)=(V(I,J-1)+V(I,J))/(2.*DTH2)
          H=2*I+1
          A(K+H)=(V(I,J)+V(I,J+1))/(2.*DTH2)
          H=I
          A(K+H)=(V(I,J)+V(I-1,J))/(2.*DR2)
          H=I+2
          A(K+H)=(V(I+1,J)+V(I,J))/(2.*DR2)
          DRDTH=(RO(I,J+1)-RO(I,J-1))/(2.*DTH)
          DRDR=(RO(I+1,J)-RO(I-1,J))/(2.*DR)
          SK(L+1)=PER*(GR*DRDTH-GTH*DRDR)
          711 CONTINUE
          L=L+1
          K=K+2*I+1
          H=I+1
          A(K+H)=1.
          SR(L+1)=.0
          710 CONTINUE
          DO 702 I=1,I1
          L=L+1
          K=K+M*I+1-L
          H=I+1
          A(K+H)=1.
          SR(L+1)=.0
          702 CONTINUE
          CALL GELB(SP,A,M,N,MUD,MLD,FPS,IER)
          KOUNT=KOUNT+1
          PRINT STREAM FUNCTION
C
C
          WRITE(6,200) KOUNT
          WRITE(6,218) IER
          WRITE(6,201) SR
C
C
          CALCULATE VELOCITIES
          UPDATE STREAM FUNCTION
C
          IF(KOUNT.EQ.1) GO TO 771
          DO 770 L=1,M
          770 SLINE(L)=SW1*SLINE(L)+SW2*SR(L)
          GO TO 773
          771 DO 772 L=1,M

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ISN 0691      772 SRIL(L)=SRIL)
ISN 0692      773 CONTINUE
C
ISN 0693      ENERGY EQUATION
ISN 0694      DO 716 I=1,ME
716 A(I)=.0
C
C      FILL IN COEFFICIENT MATRIX A AND
C      RIGHT HAND SIDE VECTOR TR
C
ISN 0695      IF(NBC.EQ.1) GO TO 74
ISN 0697      A(I)=1.
ISN 0698      TR(I)=T1
ISN 0699      GO TO 743
ISN 0700      742 CONTINUE
ISN 0701      A(I)=3.*AK(I,1)/(2.*DR)+AGAH
ISN 0702      A(I)=-2.*AK(I,1)/DR
ISN 0703      A(I)=AK(I,1)/(2.*DR)
ISN 0704      TR(I)=AGAH*IGAS
ISN 0705      743 CONTINUE
ISN 0706      L=0
ISN 0707      K=0
ISN 0708      DO 7101 I=2,III
ISN 0709      L=L+1
ISN 0710      K=K+I+L
ISN 0711      H=I+1-(I-L)
ISN 0712      A(K+H)=1.
ISN 0713      H=2*I+1-(I-L)
ISN 0714      A(K+H)=-1.
ISN 0715      TR(L+1)=.0
ISN 0716      7101 CONTINUE
C
ISN 0717      L=L+1
ISN 0718      K=K+I+L
ISN 0719      IF(NBC.EQ.1) GO TO 745
ISN 0721      H=I+1-(I-L)
ISN 0722      A(K+H)=1.
ISN 0723      TR(L+1)=T2
ISN 0724      GO TO 746
ISN 0725      745 CONTINUE
ISN 0726      H=I+1-(I-L)
ISN 0727      A(K+H)=AGAC+3.*AK(I,1)/(2.*DR)
ISN 0728      H=I-(I-L)
ISN 0729      A(K+H)=-2.*AK(I,1)/DR
ISN 0730      H=I-1-(I-L)
ISN 0731      A(K+H)=AK(I,1)/(2.*DR)
ISN 0732      TR(L+1)=AGAC*TOUT
ISN 0733      746 CONTINUE
ISN 0734      K=K-1
ISN 0735      DO 7110 J=2,JJ1
ISN 0736      L=L+1
ISN 0737      K=K+2*I+1
ISN 0738      IF(NBC.EQ.1) GO TO 752
ISN 0740      H=I+1
ISN 0741      A(K+H)=1.
ISN 0742      TR(L+1)=T1
ISN 0743      GO TO 753
ISN 0744      752 CONTINUE
ISN 0745      H=I+1

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ISN 0746      A(K+H)=3.*AK(I,J)/(2.*DR)+AGAH
ISN 0747      H=II+2
ISN 0748      A(K+H)=-2.*AK(I,J)/DR
ISN 0749      H=II+3
ISN 0750      A(K+H)=AK(I,J)/(2.*DR)
ISN 0751      TR(I+1)=AGAH*IGAS
ISN 0752      753 CONTINUE

ISN 0753      DO 7111 I=2,III
ISN 0754      L=L+1
ISN 0755      K=K+2*II+1
ISN 0756      H=II+1
ISN 0757      A(K+H)=(AK(I+1,J)+2.*AK(I,J)+AK(I-1,J))/(2.*DR2)
              1+(AK(I,J+1)+2.*AK(I,J)+AK(I,J-1))/(2.*DR2)
ISN 0758      H=1
ISN 0759      A(K+H)=(S(I+1,J)-S(I-1,J))*CP(I,J-1)/(4.*DTH*DR)
              1-(AK(I,J-1)+AK(I,J))/(2.*DR2)
ISN 0760      H=2*II+1
ISN 0761      A(K+H)=- (S(I+1,J)-S(I-1,J))*CP(I,J+1)/(4.*DTH*DR)
              1-(AK(I,J+1)+AK(I,J))/(2.*DR2)
ISN 0762      H=II
ISN 0763      A(K+H)=- (S(I,J+1)-S(I,J-1))*CP(I-1,J)/(4.*DTH*DR)
              1-(AK(I-1,J)+AK(I,J))/(2.*DR2)
ISN 0764      H=II+2
ISN 0765      A(K+H)=(S(I,J+1)-S(I,J-1))*CP(I+1,J)/(4.*DTH*DR)
              1-(AK(I+1,J)+AK(I,J))/(2.*DR2)
ISN 0766      TR(L+1)=.0
ISN 0767      7111 CONTINUE

C
ISN 0768      L=L+1
ISN 0769      K=K+2*II+1
ISN 0770      IF(MBC.EQ.1) GO TO 755
ISN 0771      H=II+1
ISN 0772      A(K+H)=1.
ISN 0773      TR(L+1)=T2
ISN 0774      GO TO 756
ISN 0775      755 CONTINUE
ISN 0776      H=II+1
ISN 0777      A(K+H)=AGAC+3.*AK(II,J)/(2.*DR)
ISN 0778      H=II
ISN 0779      A(K+H)=-2.*AK(II,J)/DR
ISN 0780      H=II-1
ISN 0781      A(K+H)=AK(II,J)/(2.*DR)
ISN 0782      TR(L+1)=AGAC*TOUT
ISN 0783      756 CONTINUE
ISN 0784      7110 CONTINUE

C
ISN 0786      L=L+1
ISN 0787      K=K+2*II+1
ISN 0788      IF(NBC.EQ.1) GO TO 762
ISN 0789      H=II+1
ISN 0790      A(K+H)=1.
ISN 0791      TR(L+1)=T1
ISN 0792      GO TO 763
ISN 0793      762 CONTINUE
ISN 0794      H=II+1
ISN 0795      A(K+H)=3.*AK(II,J)/(2.*DR)+AGAH
ISN 0796      H=II+2

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ISN 0798      A(K+H)=-2.*AK(1, JJ)/DR
ISN 0799      H=II+3
ISN 0800      A(K+H)=AK(1, JJ)/(2.*DR)
ISN 0801      TR(L+1)=AGAH*TGAS
ISN 0802      763 CONTINUE
C
ISN 0803      DO 7102 I=2, III
ISN 0804      L=I+1
ISN 0805      K=K+M+II+1-L
ISN 0806      H=II+1
ISN 0807      A(K+H)=1.
ISN 0808      H=1
ISN 0809      A(K+H)=-1.
ISN 0810      TR(L+1)=.0
ISN 0811      7102 CONTINUE
C
ISN 0812      L=L+1
ISN 0813      K=K+M+II+1-L
ISN 0814      IF(MHC.EQ.1) GO TO 765
ISN 0816      H=II+1
ISN 0817      A(K+H)=1.
ISN 0818      TR(L+1)=T2
ISN 0819      GO TO 766
ISN 0820      765 CONTINUE
ISN 0821      H=II+1
ISN 0822      A(K+H)=AGAC+3.*AK(II, JJ)/(2.*DR)
ISN 0823      H=II
ISN 0824      A(K+H)=-2.*AK(II, JJ)/DR
ISN 0825      H=II-1
ISN 0826      A(K+H)=AK(II, JJ)/(2.*DR)
ISN 0827      TR(L+1)=AGAC*TOUT
ISN 0828      766 CONTINUE
C
C      MATRIX A IS FILLED IN
C
ISN 0829      CALL GELB(TR, A, M, N, MUD, MLD, EPS, IER)
C
C      WRITE TEMPERATURES
C
ISN 0830      WRITE(6, 210)
ISN 0831      WRITE(6, 211) TR
C
C
C      CHECK THE TEMPERATURES FOR CONVERGENCE
C
ISN 0832      DO 7300 L=1, M
ISN 0833      IF(DABS(TR(L)-TLINE(L)).GT.EPT) GO TO 7116
ISN 0835      7300 CONTINUE
ISN 0836      GO TO 7115
C
C      CHECK IF KOUNT IS LESS THAN KMAX
C
ISN 0837      7116 IF(KOUNT.EQ.KMAX) GO TO 7115
C
C      UPDATE TEMPERATURES
ISN 0839      DO 780 L=1, M
ISN 0840      780 TLINE(L)=TW1*TLINE(L)+TW2*TR(L)
ISN 0841      GO TO 120
ISN 0842      7115 CONTINUE

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C
ISN 0843 T1=.0
ISN 0844 T2=.0
ISN 0845 DO 777 J=1,JJ
ISN 0846 T1=TEND(I,J)/JJ+T1
ISN 0847 T2=TEND(I,J)/JJ+T
ISN 0848 777 CONTINUE
ISN 0849 TB=(T1+T2)/2.
C MVR MASS VELOCITY IN THE R DIRECTION
C MVTH MASS VELOCITY IN THE TH DIRECTION
ISN 0850 DO 782 J=1,JJ1
ISN 0851 DO 781 I=1,II
ISN 0852 SEND(I,J)=S(I,J)
ISN 0853 781 MVR(I,J)=(SEND(I,J+1)-SEND(I,J))/DTH
ISN 0854 782 CONTINUE
ISN 0855 DO 785 I=1,III
ISN 0856 DO 786 J=1,JJ
ISN 0857 786 MVTH(I,J)=-(SEND(I+1,J)-SEND(I,J))/DR
ISN 0858 785 CONTINUE
C CALCULATE THE HEAT FLUX AND K EFFECTIVE
ISN 0859 QT(I)=.0
ISN 0860 QPHT=.0
C
ISN 0861 DO 790 J=1,JJ1
ISN 0862 T(I,J)=TEND(I,J)
ISN 0863 QPH(J)=-(-T(3,J)-T(3,J+1)+4.*(T(2,J)+T(2,J+1))-3.*(T(1,J)+
T(1,J+1)))*(AK(1,J)+AK(1,J+1))/(8.*DR)*DTH
C
ISN 0864 CONV(I,J)=MVR(I,J)*(T(1,J)+T(1,J+1))*(CP(1,J)+CP(1,J+1))/4.*DTH
ISN 0865 IF(MVR(I,J).LT..0) CONV(I,J)=MVR(I,J)*(T(1,J)+T(2,J)
1+T(1,J+1)+T(2,J+1))*(CP(1,J)+CP(1,J+1))/8.*DTH
ISN 0867 QPH(J)=QPH(J)+CONV(I,J)
ISN 0868 COND(I,J)=(T(1,J)+T(1,J+1)-T(2,J)-T(2,J+1))*(AK(1,J)+AK(1,J+1)
1/(4.*DR)*DTH
C
ISN 0869 Q(I,J)=CONV(I,J)+COND(I,J)
ISN 0870 QPHT=QPHT+QPH(J)
ISN 0871 790 QT(I)=QT(I)+Q(I,J)
C
ISN 0872 DO 791 I=2,III
ISN 0873 QT(I)=.0
ISN 0874 DO 792 J=1,JJ1
ISN 0875 T(I,J)=TEND(I,J)
ISN 0876 COND(I,J)=(T(I-1,J)+T(I-1,J+1)-T(I+1,J)-T(I+1,J+1)
1*(AK(I,J)+AK(I,J+1)))/(8.*DR)*DTH
ISN 0877 CONV(I,J)= MVR(I,J)*(T(I,J)+T(I,J+1))*(CP(I,J)+CP(I,J+1))/4.*DTH
ISN 0878 Q(I,J)=CONV(I,J)+COND(I,J)
ISN 0879 792 QT(I)=QT(I)+Q(I,J)
ISN 0880 791 CONTINUE
C
ISN 0881 QT(1)=.0
ISN 0882 QPCT=.0
ISN 0883 DO 793 J=1,JJ1
ISN 0884 T(I,J)=TEND(I,J)
ISN 0885 QPC(J)=-(-T(11-2,J)-T(11-2,J+1)+4.*(T(11-1,J)+T(11-1,J+1)
1-3.*(T(11,J)+T(11,J+1)))*(AK(11,J)+AK(11,J+1))/(8.*DR)*DTH
C
ISN 0886 CONV(11,J)=.0

```

```

ISN 0897      COND(I1,J1)=T(I1-1,J1)+T(I1-1,J+1)-T(I1,J)-T(I1,J+1)
              Y*(AK(I1,J)+AK(I1,J+1))/(4.*DR)*DTH
ISN 0898      Q(I1,J)=CONV(I1,J)+COND(I1,J)
ISN 0899      JPC(I)=JPC(I)+Q(I,J)
ISN 0900      799 QT(I1)=QT(I1)+Q(I1,J)
C
ISN 0901      QJAV=Q
ISN 0902      DO 796 I=1,II
ISN 0903      QJAV=QJAV+QT(I1)/I1
ISN 0904      796 CONTINUE
ISN 0905      AKAV=QJAV/((T1-T2)*AL)*(R2-R1)
ISN 0906      AKFE=QPC/I/((T1-T2)*AL)      *(R2-R1)
ISN 0907      AKFPH=QPH/I/((T1-T2)*AL)      *(R2-R1)
ISN 0908      450 CONTINUE
C
ISN 0909      CALCULATE RALEIGH NUMBER RA
C
ISN 0910      ANU=ETA/(APD,TH)/(ROF(APD,TH)
ISN 0911      ANU=ETA/(APD,T1)/ROF(APD,T1)
ISN 0912      AKM=PR*CLAME(APD,TH)*AKM
ISN 0913      ALF=AKM/(ROF(APD,TH)*CPE(APD,TH)
C
ISN 0914      TH=273.15/TB
ISN 0915      DR=1-7.18E-2*APD*TH*TH*(1.+0.804*APD*TH**5)
ISN 0916      BETA=-DR/DI/ROF(APD,TH)
ISN 0917      RA=6*PEP*(T1-T2)*(R2-R1)*BETA/(ALF*ANU)
ISN 0918      RE=PEP*(DPRX-GTH*ROF(APD,T1))*(R2-R1)/(ETA/(APD,TH)*ANU)
ISN 0919      RNUS=AKHC/AKD
ISN 0920      AKHC=(AKFEC+AKFPH)/2.
C
ISN 0921      WRITE(6,260)
ISN 0922      WRITE(6,261)
ISN 0923      DO 150 I=1,II
ISN 0924      DO 151 J=1,JJ
ISN 0925      TEMER=TEND(I,J)-T(I,J)
ISN 0926      WRITE(6,262) I,J,SEND(I,J),TEND(I,J),TTOIM(I,J),TEMER
ISN 0927      151 CONTINUE
ISN 0928      150 CONTINUE
ISN 0929      WRITE(6,272)
ISN 0930      DO 152 I=1,III
ISN 0931      DO 153 J=1,JJ1
ISN 0932      WRITE(6,273) I,J,MV(I,J),Q(I,J)
ISN 0933      153 CONTINUE
ISN 0934      152 CONTINUE
ISN 0935      WRITE(6,281)
ISN 0936      WRITE(6,282)
ISN 0937      DO 154 I=1,III
ISN 0938      DO 155 J=1,JJ
ISN 0939      WRITE(6,283) I,J,MVTH(I,J)
ISN 0940      155 CONTINUE
ISN 0941      154 CONTINUE
ISN 0942      WRITE(6,285)
ISN 0943      DO 157 J=1,JJ1
ISN 0944      WRITE(6,286) J,QPH(J),QPC(J)
ISN 0945      157 CONTINUE
ISN 0946      WRITE(6,287)
ISN 0947      WRITE(6,288)
ISN 0948      DO 158 I=1,II
ISN 0949      WRITE(6,289) I,QT(I)

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158 CONTINUE
159 WRITE(6,290) QPHT,QPCT,QTAV
160 PR=ANU/ALF
161 RNUS=AKHC/AKO
162 WRITE(6,291) ALF,ANU,AKO,PR,RA,RE ,RNUS
163 WRITE(6,292)
164 WRITE(6,293) AKAV,AKFFH,AKFFC,AKHC
165 200 FORMAT(10,20X,'SR',5X,'KOUNT=',1X,I3,///)
166 201 FORMAT(0,5X,1CF11.3)
167 210 FORMAT(0,20X,'TR TEMPERATURES',///)
168 211 FORMAT(0,5X,1CF11.3)
169 218 FORMAT(0,5X,14)
170 260 FORMAT(1,10X,'RESULTS FROM THE LAST ITERATION',/)
171 261 FORMAT(0,1 I J STREAM FUNCTION TEMPERATURES DIMENSIONAL
172 155 TEMPERATURES,K TEMER')
173 262 FORMAT(0,12,2X,12,5X,F10.3,4X,F8.3,15X,F10.3,10X,F5.2)
174 272 FORMAT(1,1 I J MVR (KG/(SQM*2) Q (W))
175 273 FORMAT(0,12,2X,12,8X,F10.4,6X,F10.3)
176 281 FORMAT(1,1 MVT,COMPONENT OF MASS VELOCITY PERPENDICULAR TO MVR')
177 282 FORMAT(0,1 I J MVT(KG/(SQM*2))
178 283 FORMAT(0,12,2X,12,8X,F10.3)
179 285 FORMAT(1,1 J QPH(W)
180 286 FORMAT(0,12,4X,F10.3,4X,E10.3)
181 287 FORMAT(1,1 HEAT FLOW THROUGH THE INSULATION,QT,QPHT,QPCT ,QTAV')
182 288 FORMAT(0,1 I QT(W))
183 289 FORMAT(0,12,4X,E10.3)
184 290 FORMAT(0,1 QPHT=',E10.3,'QPCT=',F10.3,'QTAV=',E10.3,///)
185 291 FORMAT(0,1 ALF=',E11.4,'ANU=',E11.4,'AKO=',E11.4,'PR=',E11.4,'RA=
186 1',E11.4,'RE=',F11.4,5X,'RNUS=',E11.4,///)
187 292 FORMAT(0,1 EFFECTIVE THERMAL CONDUCTIVITY, W/(KM),///)
188 293 FORMAT(0,1 AKAV=',E11.4,'AKFFH=',F11.4,'AKFFC=',E11.4,'AKHC='
189 1,E11.4)
190 300 FORMAT(5X,715,5X,F5.2)
191 301 FORMAT(5X,3F10.5)
192 302 FORMAT(5X,3F10.3)
193 303 FORMAT(5X,F5.3,2E10.3)
194 304 FORMAT(5X,3E10.3)
195 305 FORMAT(5X,2F10.4)
196 STOP
197 END

```

OPTIONS IN EFFECT NAME= MAIN,OPT=00,LINFCNT=60,SIZE=0000K,

OPTIONS IN EFFECT SOURCE,EBCDIC,NOLIST,NODFCK,LOAD,NOMAF,NOEDIT,NOID,NIXREF

STATISTICS SOURCE STATEMENTS = 974 ,PROGRAM SIZE = 203798

STATISTICS NO DIAGNOSTICS GENERATED

***** END OF COMPILATION *****

LEVEL 21.7 (JAN 73)

DS7160 FORTRAN H

COMPILER OPTIONS - NAME= MAIN,OPT=00,LINECNT=60,SIZE=0000K,
SOURCE,FBC,DEC,NOLIST,NODECK,LOAD,NOMAP,NOEDIT,NOID,NIXREF

C		GELB 10
C	GELB 20
C		GELB 30
C	SUBROUTINE GELB	GELB 40
C		GELB 50
C	PURPOSE	GELB 60
C	TO SOLVE A SYSTEM OF SIMULTANEOUS LINEAR EQUATIONS WITH A	GELB 70
C	COEFFICIENT MATRIX OF BAND STRUCTURE.	GELB 80
C		GELB 90
C	USAGE	GELB 100
C	CALL GELB(R,A,M,N,MUD,MLD,EPS,IFR)	GELB 110
C	DESCRIPTION OF PARAMETERS	GELB 130
C	R - M BY N RIGHT HAND SIDE MATRIX (DESTROYED).	GELB 140
C	ON RETURN R CONTAINS THE SOLUTION OF THE EQUATIONS.	GELB 150
C	A - M BY M COEFFICIENT MATRIX WITH BAND STRUCTURE	GELB 160
C	(DESTROYED).	GELB 170
C	M - THE NUMBER OF EQUATIONS IN THE SYSTEM.	GELB 180
C	N - THE NUMBER OF RIGHT HAND SIDE VECTORS.	GELB 190
C	MUD - THE NUMBER OF UPPER CUDIAGONALS (THAT MEANS	GELB 200
C	CUDIAGONALS ABOVE MAIN DIAGONAL).	GELB 210
C	MLD - THE NUMBER OF LOWER CUDIAGONALS (THAT MEANS	GELB 220
C	CUDIAGONALS BELOW MAIN DIAGONAL).	GELB 230
C	EPS - AN INPUT CONSTANT WHICH IS USED AS RELATIVE	GELB 240
C	TOLERANCE FOR TEST ON LOSS OF SIGNIFICANCE.	GELB 250
C	IFR - RESULTING ERROR PARAMETER CODED AS FOLLOWS	GELB 260
C	IFR=0 - NO ERROR.	GELB 270
C	IFR=-1 - NO RESULT BECAUSE OF WRONG INPUT PARAME-	GELB 280
C	TERS M,MUD,MLD OR BECAUSE OF PIVOT ELEMENT	GELB 290
C	AT ANY ELIMINATION STEP EQUAL TO 0.	GELB 300
C	IFR=K - WARNING DUE TO POSSIBLE LOSS OF SIGNIFI-	GELB 310
C	CANCE INDICATED AT ELIMINATION STEP K+1,	GELB 320
C	WHERE PIVOT ELEMENT WAS LESS THAN OR	GELB 330
C	EQUAL TO THE INTERNAL TOLERANCE EPS TIMES	GELB 340
C	ABSOLUTELY GREATEST ELEMENT OF MATRIX A.	GELB 350
C		GELB 360
C	REMARKS	GELB 370
C	BAND MATRIX A IS ASSUMED TO BE STORED ROWWISE IN THE FIRST	GELB 380
C	ME SUCCESSIVE STORAGE LOCATIONS OF TOTALLY NEEDED MA	GELB 390
C	STORAGE LOCATIONS, WHERE	GELB 400
C	MA=M*MC-ML*(ML+1)/2 AND ME=MA-MU*(MU+1)/2 WITH	GELB 410
C	MC=M*(M,1+MUD+MLD), ML=MC-1-MLD, MU=MC-1-MUD.	GELB 420
C	RIGHT HAND SIDE MATRIX R IS ASSUMED TO BE STORED COLUMNWISE	GELB 430
C	IN M*MC SUCCESSIVE STORAGE LOCATIONS. ON RETURN SOLUTION	GELB 440
C	MATRIX R IS STORED COLUMNWISE TOO.	GELB 450
C	INPUT PARAMETERS M, MUD, MLD SHOULD SATISFY THE FOLLOWING	GELB 460
C	RESTRICTIONS MUD NOT LESS THAN ZERO	GELB 470
C	MLD NOT LESS THAN ZERO	GELB 480
C	MUD+MLD NOT GREATER THAN 2*M-2.	GELB 490
C	NO ACTION BESIDES ERROR MESSAGE IFR=-1 TAKES PLACE IF THESE	GELB 500
C	RESTRICTIONS ARE NOT SATISFIED.	GELB 510
C	THE PROCEDURE GIVES RESULTS IF THE RESTRICTIONS ON INPUT	GELB 520
C	PARAMETERS ARE SATISFIED AND IF PIVOT ELEMENTS AT ALL	GELB 530
C	ELIMINATION STEPS ARE DIFFERENT FROM 0. HOWEVER WARNING	GELB 540
C	IFR=K - IF GIVEN - INDICATES POSSIBLE LOSS OF SIGNIFICANCE.	GELB 550
C	IN CASE OF A WELL SCALED MATRIX A AND APPROPRIATE TOLERANCE	GELB 560
C	EPS, IFR=K MAY BE INTERPRETED THAT MATRIX A HAS THE RANK K.	GELB 570


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C      INSERT ZEROS IN FIRST MU ROWS (NOT NECESSARY IN CASE MZ=0)
ISN 0031      IFMZ(14,14,10)
ISN 0032      10 JJ=1
ISN 0033      J=1+M2
ISN 0034      IC=1+MUJ
ISN 0035      DO 13 I=1,MU
ISN 0036      DO 12 K=1,MC
ISN 0037      A(I,J)=0.
ISN 0038      IFK=IC(11,11,12)
ISN 0039      11 A(I,J)=A(I,J)
ISN 0040      J=J+1
ISN 0041      12 JJ=JJ+1
ISN 0042      13 IC=IC+1
C
C      GENERATE TEST VALUE FOR SINGULARITY
ISN 0043      14 TOL=EPS*PIV
C
C      START DECOMPOSITION LOOP
ISN 0044      RST=1
ISN 0045      IDST=MC
ISN 0046      IC=MC-1
ISN 0047      DO 38 K=1,M
ISN 0048      IFK=MR(1,16,16,15)
ISN 0049      15 IDST=IDST-1
ISN 0050      16 ID=ID+1
ISN 0051      ILR=IC+D
ISN 0052      IF(I=1,18,18,17)
ISN 0053      17 ILC=IC
ISN 0054      18 ILC=IC
C
C      PIVOT SEARCH IN FIRST COLUMN (ROW INDEXES FROM I+1 TO I+ILR)
ISN 0055      I=I+1
ISN 0056      DO 21 J=K+1,ILR
ISN 0057      A=ABS(A(I,J))
ISN 0058      IF (A > PIV) 20,20,19
ISN 0059      19 P=PIV
ISN 0060      J=J+1
ISN 0061      A=0.
ISN 0062      20 ILC=MR(22,22,21)
ISN 0063      21 ILC=ID-1
ISN 0064      22 ILC=I+10
C
C      TEST ON SINGULARITY
ISN 0065      23 ILC=PIV(47,47,23)
ISN 0066      24 ILC=IFR(26,24,26)
ISN 0067      24 ILC=PIV-TOL(25,25,26)
ISN 0068      25 ILC=K-1
ISN 0069      26 ILC=1/A(I,J)
C
C      PIVOT ROW REDUCTION AND ROW INTERCHANGE IN RIGHT HAND SIDE R
ISN 0070      I=J-K
ISN 0071      DO 27 I=K+1,M
ISN 0072      I=I+10
ISN 0073      TR=PIV*(I)
ISN 0074      R(I)=R(I)
ISN 0075      27 R(I)=TR
C
GELB1150
GELB1160
GELB1170
GELB1180
GELB1190
GELB1200
GELB1210
GELB1220
GELB1230
GELB1240
GELB1250
GELB1260
GELB1270
GELB1280
GELB1290
GELB1300
GELB1310
GELB1320
GELB1330
GELB1340
GELB1350
GELB1360
GELB1370
GELB1380
GELB1390
GELB1400
GELB1410
GELB1420
GELB1430
GELB1440
GELB1450
GELB1460
GELB1470
GELB1480
GELB1500
GELB1510
GELB1520
GELB1530
GELB1540
GELB1550
GELB1560
GELB1570
GELB1580
GELB1590
GELB1600
GELB1610
GELB1620
GELB1630
GELB1640
GELB1650
GELB1660
GELB1670
GELB1680
GELB1690
GELB1700
GELB1710
GELB1720

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		PIVOT ROW REDUCTION AND ROW INTERCHANGE IN COEFFICIENT MATRIX A	GEL01730
ISN 0076		II=KST	GEL01740
ISN 0077		J=JJ+IC	GEL01750
ISN 0078		DO 28 I=JJ,J	GEL01760
ISN 0079		TH=PIV*AI(I)	GEL01770
ISN 0080		AI(I)=AI(I)	GEL01780
ISN 0081		AI(I)=TH	GEL01790
ISN 0082	28	II=II+1	GEL01800
	C		GEL01810
	C	ELEMENT REDUCTION	GEL01820
ISN 0083		IF (R=ILR) 23, 34, 34	GEL01830
ISN 0084	29	II=KST	GEL01840
ISN 0085		II=K+1	GEL01850
ISN 0086		MU=KST+1	GEL01860
ISN 0087		MZ=KST+IC	GEL01870
ISN 0088		DO 33 I=II,ILR	GEL01880
	C		GEL01890
	C	IN MATRIX A	GEL01900
ISN 0089		ID=ID+MC	GEL01910
ISN 0090		JJ=I-MH-1	GEL01920
ISN 0091		IF (JJ) 31, 31, 30	GEL01930
ISN 0092	30	ID=ID-JJ	GEL01940
ISN 0093	31	PIV=-AI(ID)	GEL01950
ISN 0094		J=ID+1	GEL01960
ISN 0095		DO 32 JJ=MU,MZ	GEL01970
ISN 0096		AI(J-1)=AI(J)+PIV*AI(JJ)	GEL01980
ISN 0097	32	J=J+1	GEL01990
ISN 0098		AI(J-1)=G.	GEL02000
	C		GEL02010
	C	IN MATRIX H	GEL02020
ISN 0099		J=K	GEL02030
ISN 0100		DO 33 JJ=I,MM,M	GEL02040
ISN 0101		H(IJJ)=H(IJJ)+PIV*P(IJ)	GEL02050
ISN 0102	33	J=J+M	GEL02060
ISN 0103	34	KST=KST+MC	GEL02070
ISN 0104		IF (ILR=MH) 34, 35, 35	GEL02080
ISN 0105	35	IC=IC-1	GEL02090
ISN 0106	36	ID=K-MH	GEL02100
ISN 0107		IF (ID) 36, 36, 37	GEL02110
ISN 0108	37	KST=KST-ID	GEL02120
ISN 0109	38	CONTINUE	GEL02130
	C	END OF DECOMPOSITION LOOP	GEL02140
	C		GEL02150
	C		GEL02160
	C	BACK SUBSTITUTION	GEL02170
ISN 0110		IF (MC=I) 46, 46, 45	GEL02180
ISN 0111	39	IC=2	GEL02190
ISN 0112		KST=MA+ML-MC+2	GEL02200
ISN 0113		II=M	GEL02210
ISN 0114		DO 45 I=2,M	GEL02220
ISN 0115		KST=KST-MC	GEL02230
ISN 0116		II=II-1	GEL02240
ISN 0117		J=II-MH	GEL02250
ISN 0118		IF (J) 41, 41, 40	GEL02260
ISN 0119	40	KST=KST+J	GEL02270
ISN 0120	41	DO 43 J=II,MM,M	GEL02280
ISN 0121		TH=H(IJ)	GEL02290
ISN 0122		MZ=KST+IC-2	GEL02300

ISN 0123	ID=J	GFLB2310
ISN 0124	(M) 47 JJ=KST,MZ	GFLB2320
ISN 0125	ID=ID+1	GELB2330
ISN 0126	42 TB=TB-AIJJ)ORII0I	GELB2340
ISN 0127	43 MCJJ=TB	GELB2350
ISN 0128	IF TIC=MC)44,45,45	GELB2360
ISN 0129	44 IC=IC+1	GELB2370
ISN 0130	45 CONTINUE	GELB2380
ISN 0131	46 RETURN	GELB2390
	C	GELB2400
	C	GELB2410
	C	GELB2420
ISN 0132	ENDIN RETURN	GFLB2430
ISN 0133	47 IER=-1	GELB244
ISN 0134	RETURN	GELB2450
	END	

OPTIONS IN EFFECT NAME= MAIN,OPT=00,LINECNT=60,SIZE=0000K.

OPTIONS IN EFFECT SOURCE,ERGOIC,NOLIST,NOOFCR,LOAD,NOMAP,NOEDIT,NOID,NOXREF

STATISTICS SOURCE STATEMENTS = 133 ,PROGRAM SIZE = 2786

STATISTICS NO DIAGNOSTICS GENERATED

***** END OF COMPILATION *****

187K BYTES OF CORE NOT USED

LEVEL 21.7 (JAN 71)

015/160 FORTRAN M

COMPILER OPTIONS - NAME = MAIN,OPT=00,LENGTH=60,SIZE=00000,
SOURCE,FREDEC,NOLIST,NODECK,LOAD,NOWAP,NOEDIT,NOID,NOXREF

```

TSN 0002      FUNCTION FTAF(P,T)
C              DYNAMIC VISCOSITY OF C12 (P=ATM, T=K, FTA=ERG/CM*SEC)
TSN 0003      P=1.
TSN 0004      T=273.16
TSN 0005      PF=P/P0
TSN 0006      TF=T/T0
TSN 0007      FTAN=(1.54E-70*SQT (T))/T*(1+(224./T))
TSN 0008      FTAF=FTAN*(1+.0478E-30*(PF-1.)*(TF**3))      *0.4068
TSN 0009      RETURN
TSN 0010      END

```

OPTIONS IN EFFECT NAME= MAIN,OPT=00,LENGTH=60,SIZE=00000.

OPTIONS IN EFFECT SOURCE,FREDEC,NOLIST,NODECK,LOAD,NOWAP,NOEDIT,NOID,NOXREF

STATISTICS SOURCE STATEMENTS = 4 ,PROGRAM SIZE = 424

STATISTICS NO DIAGNOSTICS GENERATED

***** END OF COMPILATION *****

205K BYTES OF CORE NOT USED

LEVEL 21.7 (JAN 73)

D57160 FORTRAN M

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      COMPILER OPTIONS - NAME = MAIN,OPT=00,LINFCNT=60,SIZE=0000K,
                        SOURCE,ERCDDIC,NOLIST,NODECK,LOAD,NOMAP,NOEDIT,NOID,NOXREF
ISN 0002      FUNCTION CLAMP(P,T)
      C      THERMAL CONDUCTIVITY OF CO2 (P=ATM, T=K, L=M/HR)
ISN 0003      P0=1.
ISN 0004      T0=273.16
ISN 0005      PF=P/P0
ISN 0006      TF=T/T0
ISN 0007      IF(PF-1.) 4,4,5
ISN 0008      4 FCL = PF-1.
ISN 0009      GO TO 6
ISN 0010      5 FCL = (PF -1.)**1.25
ISN 0011      6 IF(T-T0-725.) 1,1,2
ISN 0012      1 CA=3.6943E2
ISN 0013      CB=1.6768E5
ISN 0014      CC=2.7331E7
ISN 0015      GO TO 3
ISN 0016      2 CA=4.0476E2
ISN 0017      CC=-1.9204E7
ISN 0018      CB=1.5704E5
ISN 0019      3 CLAMP=(SQRT (T))/(CA+(CB/T)+(CC/(T*T)))
ISN 0020      CLAMP=CLAMP*(1.+2.14E-3*FCL*(TF**2.36))      01.1623
ISN 0021      RETURN
ISN 0022      END

*OPTIONS IN EFFECT*      NAME= MAIN,OPT=00,LINFCNT=60,SIZE=0000K,
*OPTIONS IN EFFECT*      SOURCE,ERCDDIC,NOLIST,NODECK,LOAD,NOMAP,NOEDIT,NOID,NOXREF
*STATISTICS*      SOURCE STATEMENTS =      21 ,PROGRAM SIZE =      674
*STATISTICS* NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****
*STATISTICS* NO DIAGNOSTICS THIS STEP
205K BYTES OF CORE NOT USED

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RESUMO

Uma solução numérica foi obtida para as distribuições de fluxo e de temperatura em isolamentos térmicos do tipo fibra. Foram consideradas geometrias retangular e cilíndrica. As condições de contorno utilizadas também incluíram os casos de parede quente permeável. Para cavidades retangulares foi obtida uma boa concordância com resultados publicados, tanto numéricos como experimentais. As distribuições de velocidade e de temperatura calculadas propiciaram um melhor entendimento dos fenômenos de escoamento e de transferência de calor em isolamentos tipo fibra. Os números de Nusselt, locais e médio, obtidos para a parede fria constituem informações úteis para o projeto de isolações do tipo de fibras para vasos e tubulações de usinas nucleares resfriadas a gás. O número de Nusselt médio foi correlacionado, com o Número de Rayleigh nos casos apenas de convecção natural, e com os números de Rayleigh e Reynolds nos casos de uma combinação de convecção natural com convecção forçada.

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